

Коллективные явления в холодных непрямых экситонах

Collective phenomena in cold indirect excitons

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Sen Yang, A.T. Hammack, A. V. Mintsev, A.G. Winbow, E.E. Novitskaya, A.A. High, M. Remeika, J.C. Graves, G. Grosso, A.K. Thomas, Y.Y. Kuznetsova, J.R. Leonard, P. Andreakou, S.V. Poltavtsev, E.V. Calman, M.W. Hasling, C.J. Dorow, L.H. Fowler-Gerace, D.J. Choksy (*UCSD*)

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S.V. Lobanov, N.A. Gippius (*Skoltech*)

J. Wilkes, A.L. Ivanov (*Cardiff*)

D.E. Nikonov, I.A. Young (Intel)

B.D. Simons (*Cambridge*), L.S. Levitov (*MIT*)

K.L. Campman, M. Hanson, A.C. Gossard (*UCSB*)

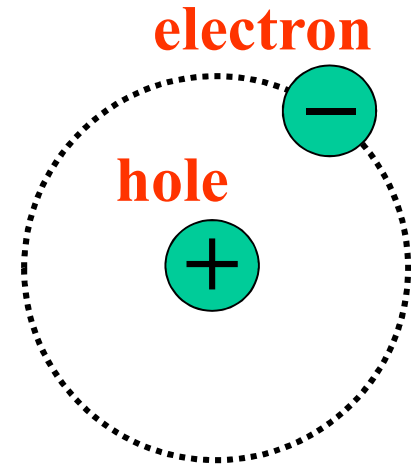
L.N. Pfeiffer, K.W. West (*Princeton*)

S. Hu, A. Mishchenko, A.K. Geim, K.S. Novoselov (*Manchester*)

- Indirect excitons (IXs) aka interlayer excitons
- Spontaneous coherence and condensation of IXs
- Phenomena in IX condensate
 - Density wave, commensurability effect
 - Spin textures
 - Pancharatnam-Berry phase, coherent spin transport
 - Phase singularities, interference dislocations
- IXs in van der Waals heterostructures
 - Opportunity to realize high-T IX condensation
 - IXs at room temperature
 - Indirect trions

exciton – bound pair of electron and hole

$$m_{exciton} = m_{electron} + m_{hole} \ll m_{atom}$$



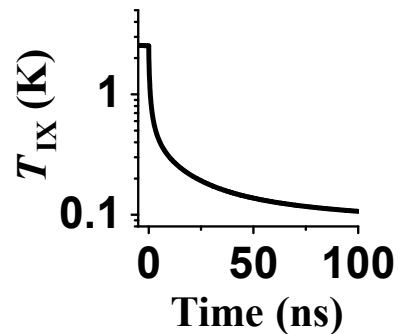
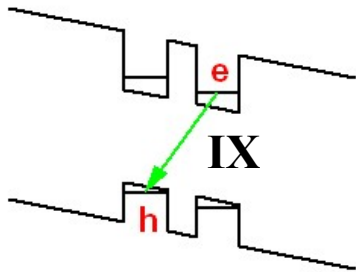
exciton – light bosonic particle in semiconductor

Indirect excitons (IXs) aka interlayer excitons

Degenerate Bose gas of excitons: thermal de Broglie wavelength $\lambda_{dB} = \left(\frac{2\pi\hbar^2}{mk_B T} \right)^{1/2}$
 \sim separation between excitons

Temperature of quantum degeneracy $T_0 = \frac{2\pi\hbar^2}{m_x} n \sim 3 \text{ K}$
 \uparrow
 excitons in GaAs QW
 $n = 10^{10} \text{ cm}^{-2}, m_x = 0.2 m_e$

IXs in CQW



IXs cool to 100 mK within $\sim 100 \text{ ns}$ lifetime
PRL 86, 5608 (2001)

$T_{IX} \sim 100 \text{ mK} \ll T_0$
 is realized for IXs



UCSD

Spontaneous coherence and condensation of IXs

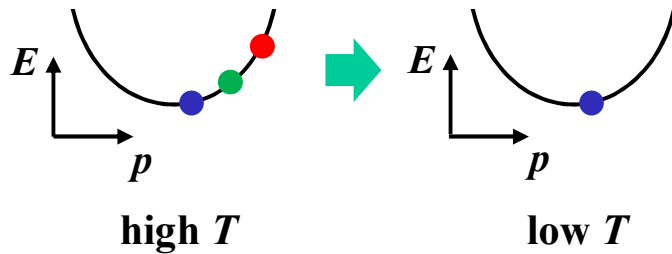
Sen Yang, A.T. Hammack, M.M. Fogler, L.V. Butov, A.C. Gossard,
PRL 97, 187402 (2006)

A.A. High, J.R. Leonard, A.T. Hammack, M.M. Fogler, L.V. Butov,
A.V. Kavokin, K.L. Campman, A.C. Gossard, *Nature* 483, 584 (2012)

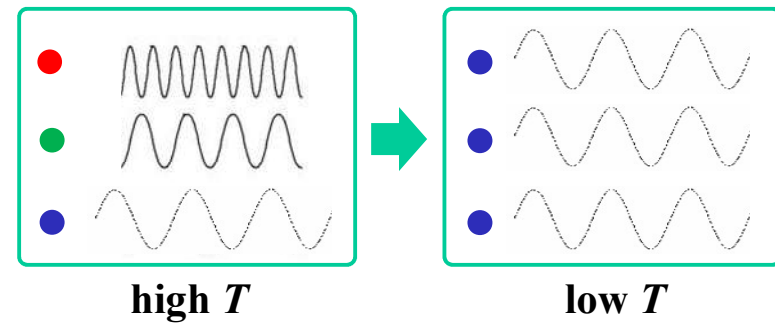
Below the temperature of quantum degeneracy
bosonic particles can form a **coherent state** \leftrightarrow **BEC**

Condensation in momentum space = Spontaneous coherence of matter waves

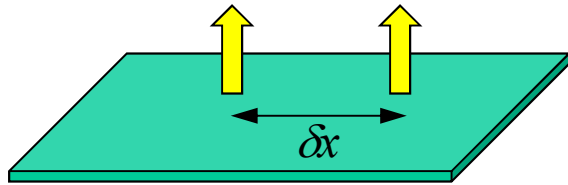
particles (in momentum space)



matter waves (in real space) $\lambda = h/p$

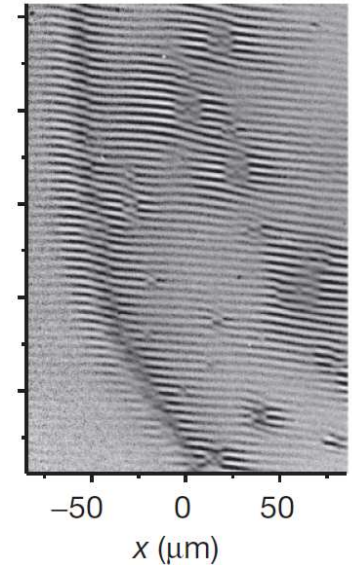
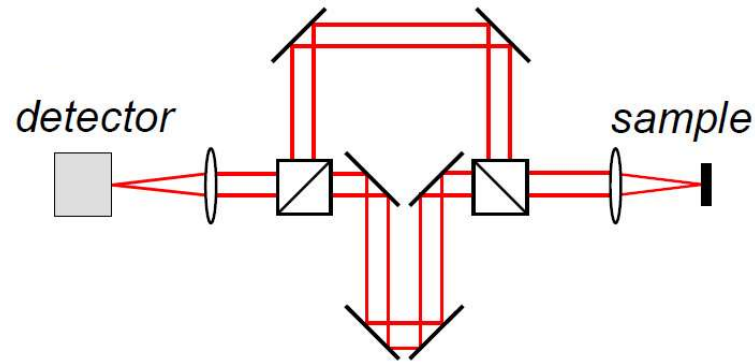


Direct measurement of spontaneous coherence and condensation

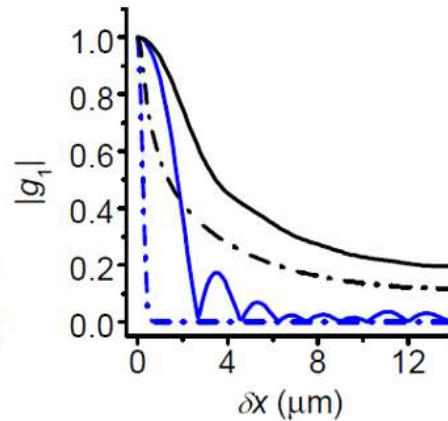
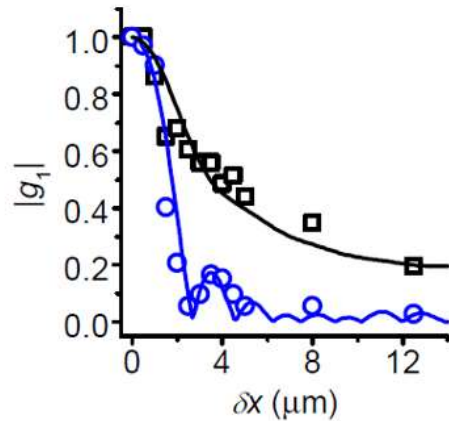


exciton coherence
is imprinted on coherence
of their light emission

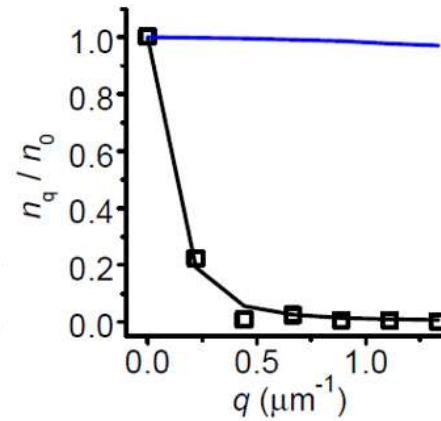
measured by shift-interferometry



First order coherence function $g_1(\Delta x)$



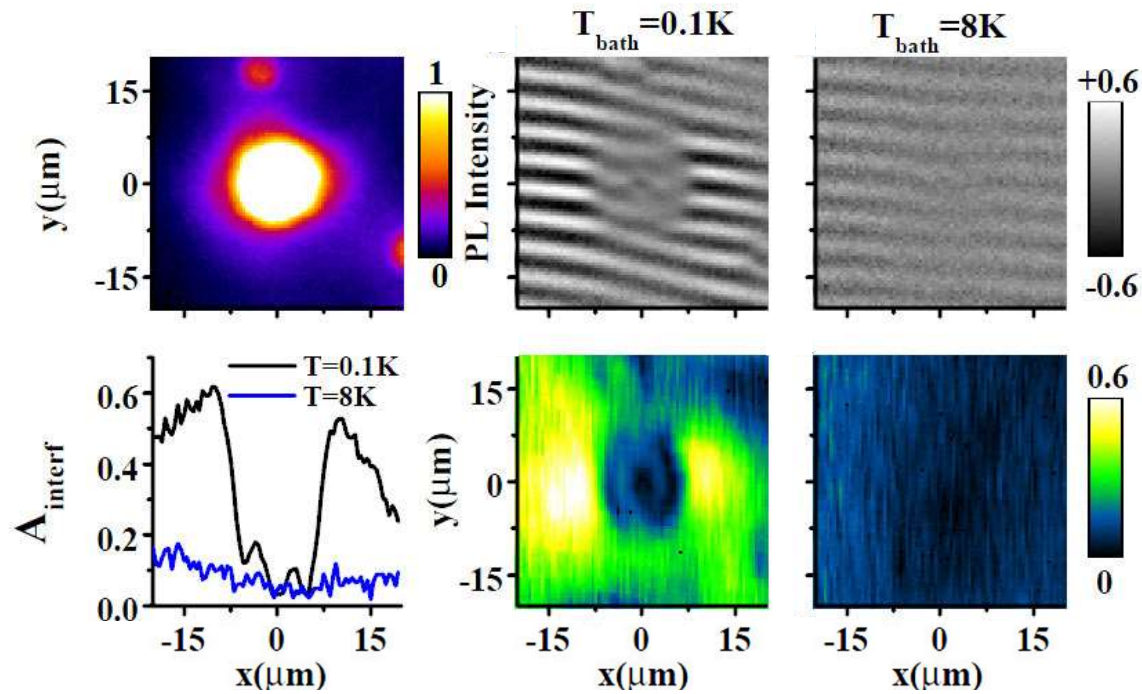
Distribution in q -space n_q



spontaneous
coherence of exciton
matter waves =
exciton condensation
in momentum space

$$g_1(r) \sim \int d^2q e^{iqr} n_q$$

IX spontaneous coherence



emergence of
spontaneous coherence
around source of IXs
at low T at $r > r_{\text{coh}}$

Phenomena in IX condensate

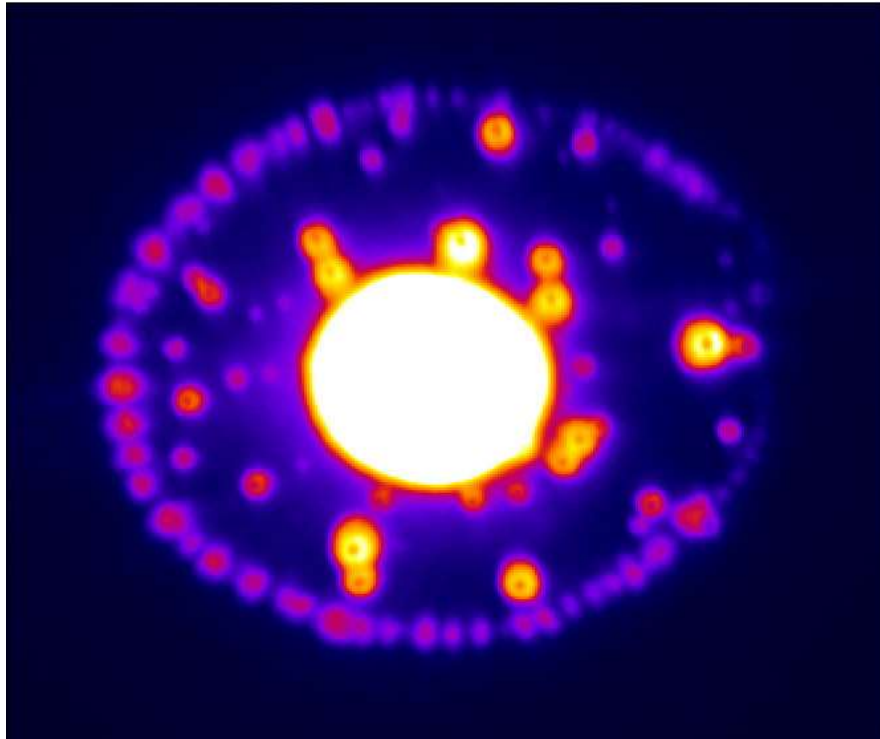
**Density wave
and commensurability effect
in IX condensate**

L.V. Butov, A.C. Gossard, D.S. Chemla, *Nature* 418, 751 (2002)

L.S. Levitov, B.D. Simons, L.V. Butov, *PRL* 94, 176404 (2005)

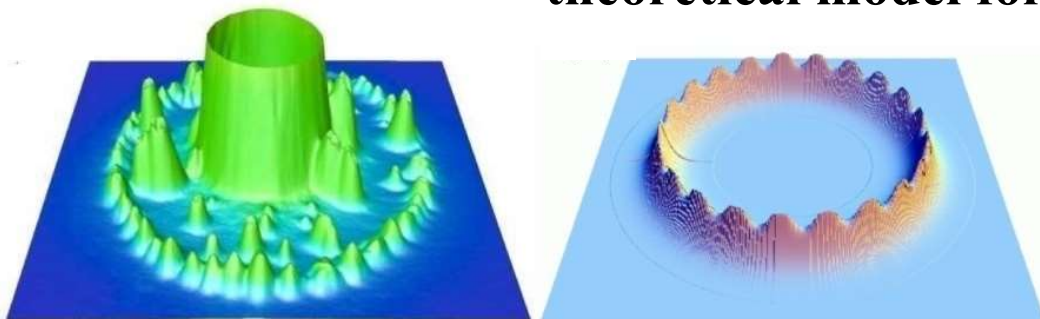
Sen Yang, L.V. Butov, B.D. Simons, K.L. Campman, A.C. Gossard,
PRB 91, 245302 (2015)

macroscopically ordered exciton state (MOES) or exciton density wave



L.V. Butov, A.C. Gossard, D.S. Chemla,
Nature 418, 751 (2002)

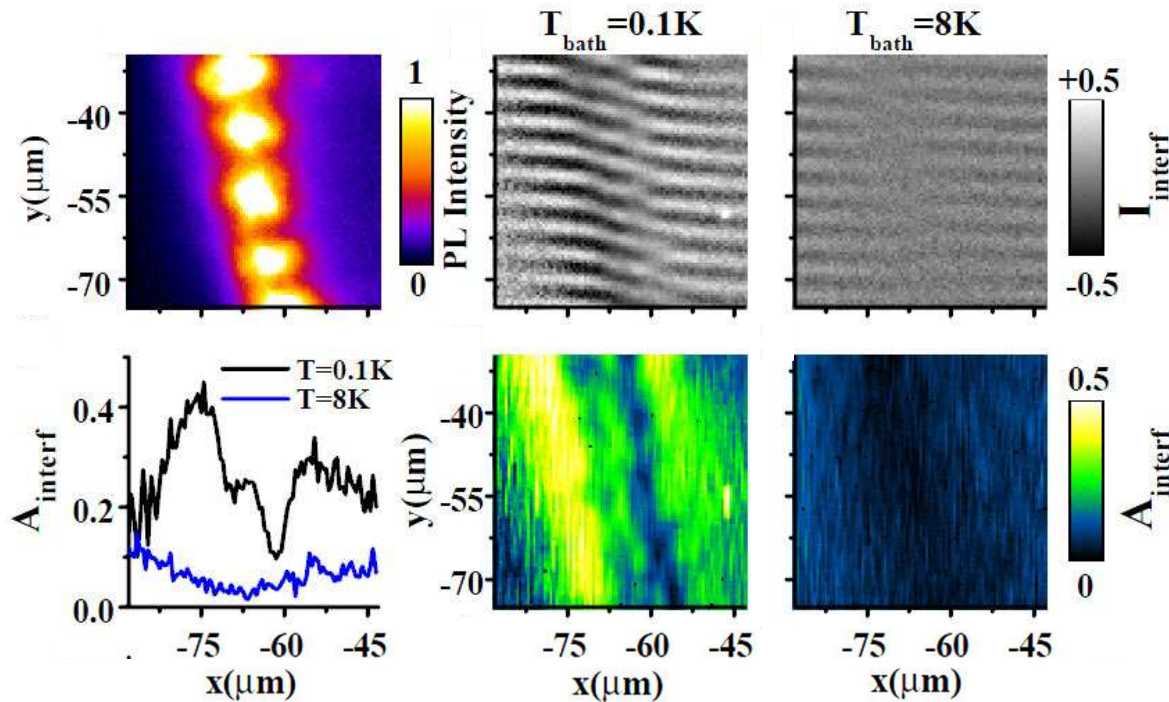
theoretical model for MOES



instability results from
quantum degeneracy due to
stimulated kinetics of exciton formation

L.S. Levitov, B.D. Simons, L.V. Butov,
PRL 94, 176404 (2005)

IX spontaneous coherence in MOES



emergence of
spontaneous coherence
in MOES
(in IX density wave)
at low T at $r > r_{\text{coh}}^*$

MOES is **condensate in k-space**
with **macroscopic spatial order**

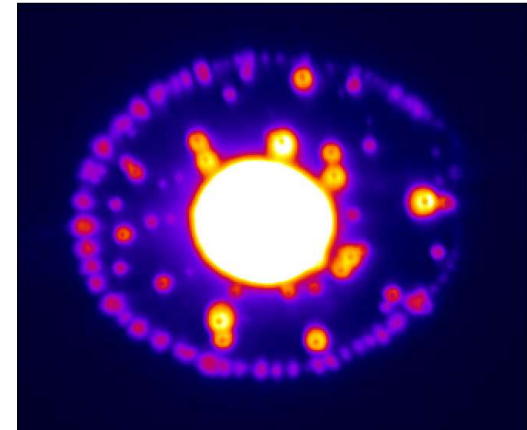
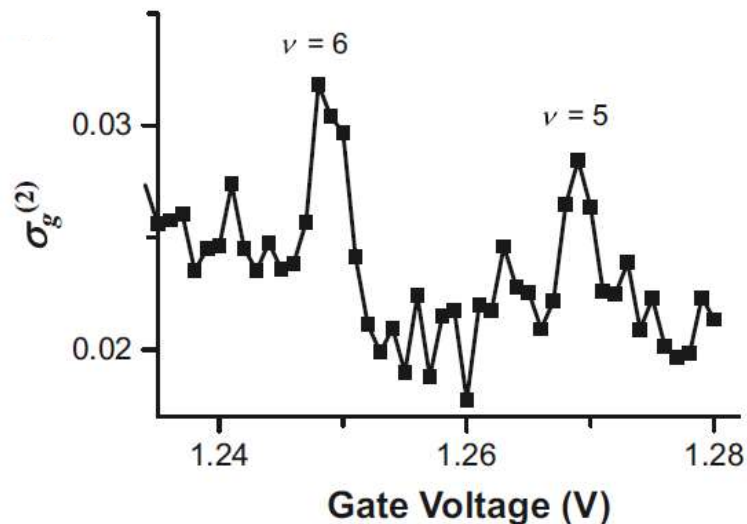
Sen Yang, A.T. Hammack, M.M. Fogler, L.V. Butov, A.C. Gossard, *PRL* 97, 187402 (2006)

A.A. High, J.R. Leonard, A.T. Hammack, M.M. Fogler, L.V. Butov, A.V. Kavokin, K.L. Campman, A.C. Gossard, *Nature* 483, 584 (2012)

Commensurability effect of IX density wave

fluctuations of exciton density wave
are suppressed when
number of wavelength on wave
confinement length is integer

$$\nu = L / \lambda_{\text{IX-wave}} = N$$



commensurability effect: macroscopic system
of IXs of length $\sim 100 \mu\text{m}$ behaves collectively:
MOES is collective phenomenon

$$l_{\text{commensurability}} \gg \lambda_{\text{IX-wave}} > \xi_{\text{coh}} \gg \lambda_{\text{dB}}$$

↑
↓

MOES is condensate
in momentum space

IX density wave and commensurability effect



instability due to stimulated processes

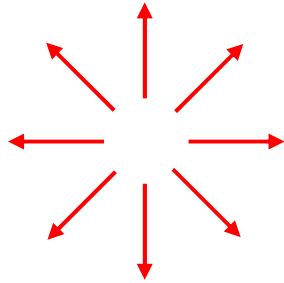
**Pancharatnam-Berry phase,
spin textures,
and long-range coherent spin transport
in IX condensate**

A.A. High, A.T. Hammack, J.R. Leonard, Sen Yang, L.V. Butov,
T. Ostatnický, M. Vladimirova, A.V. Kavokin, K.L. Campman,
A.C. Gossard, *PRL* 110, 246403 (2013)

J.R. Leonard, A.A. High, A.T. Hammack, M.M Fogler, L.V. Butov,
K.L. Campman, A.C. Gossard, *Nature Commun* 9, 2158 (2018)

IX transport and spin precession

IX transport from radial source of cold IXs



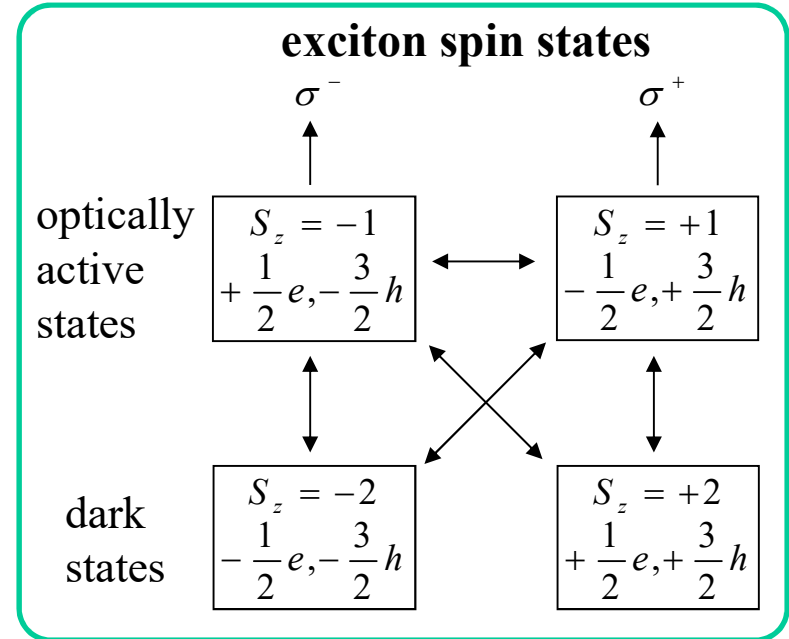
within coherence length



ballistic IX transport with spin precession



due to spin-orbit interaction, splitting of exciton states, Zeeman effect



polarization textures

Pancharatnam-Berry phase

discovered
by Pancharatnam for light

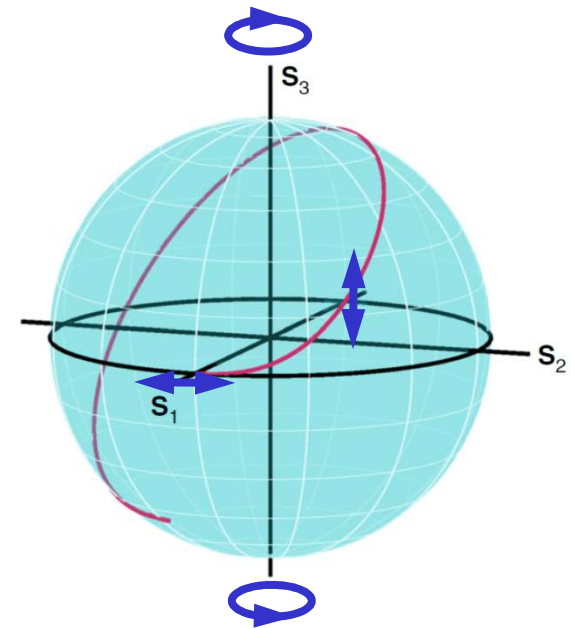
The **Pancharatnam-Berry phase** is a geometric phase

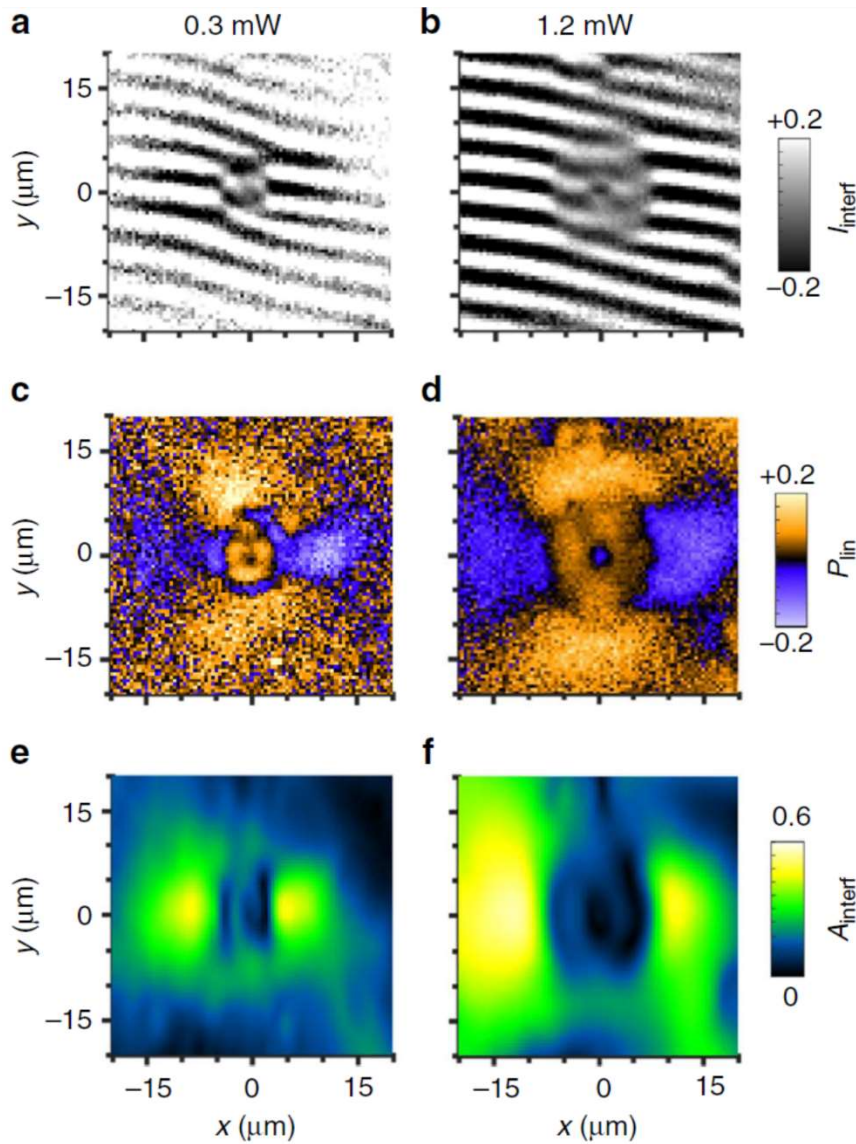
- appearing when the polarization state of light changes
- acquired over a cycle of parameters in the Hamiltonian governing the system

by Berry for matter waves

polarization state of light
goes along closed contour
on Poincaré sphere

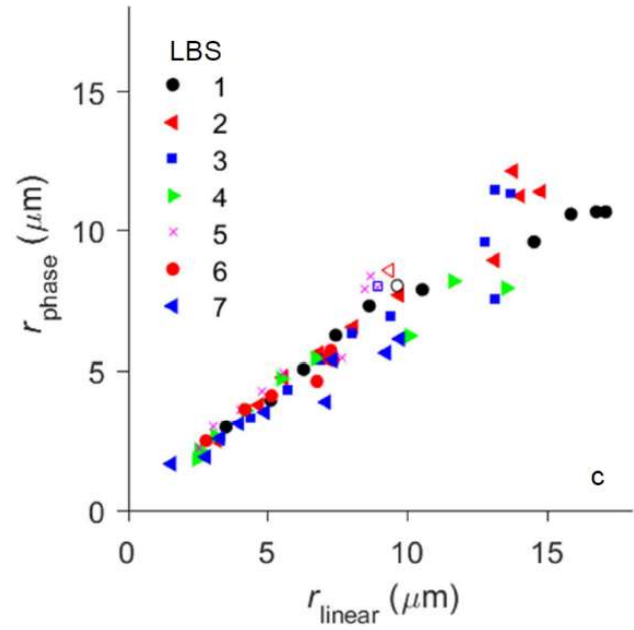
acquired Pancharatnam-Berry
phase $\phi_{PB} = 1/2 \Omega$





correlation with polarization
 identifies the phase as PB phase

correlations between
 phase shifts,
 polarization pattern, and
 onset of spontaneous coherence

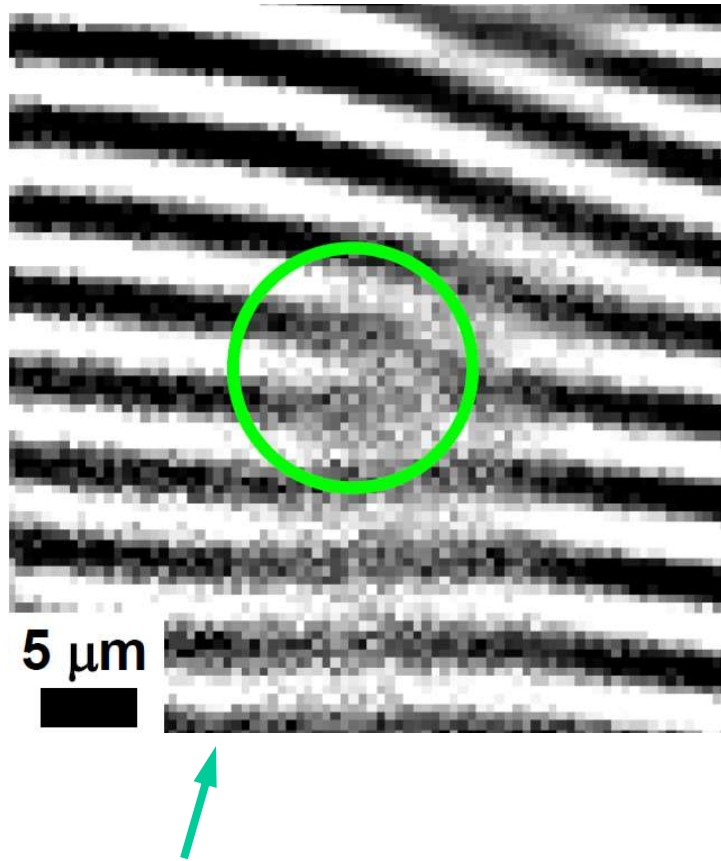


data for different sources
 and different source powers
 collapse on universal line $r_{phase} = r_{linear}$

Phase singularities, interference dislocations in IX condensate

J.R. Leonard, Lunhui Hu, A.A. High, A.T. Hammack, Congjun Wu,
L.V. Butov, K.L. Campman, A.C. Gossard, arXiv:1910.06387

Dislocations (forks) in IX interference patterns



forks in interference patterns are commonly associated with **vortices** in quantum systems

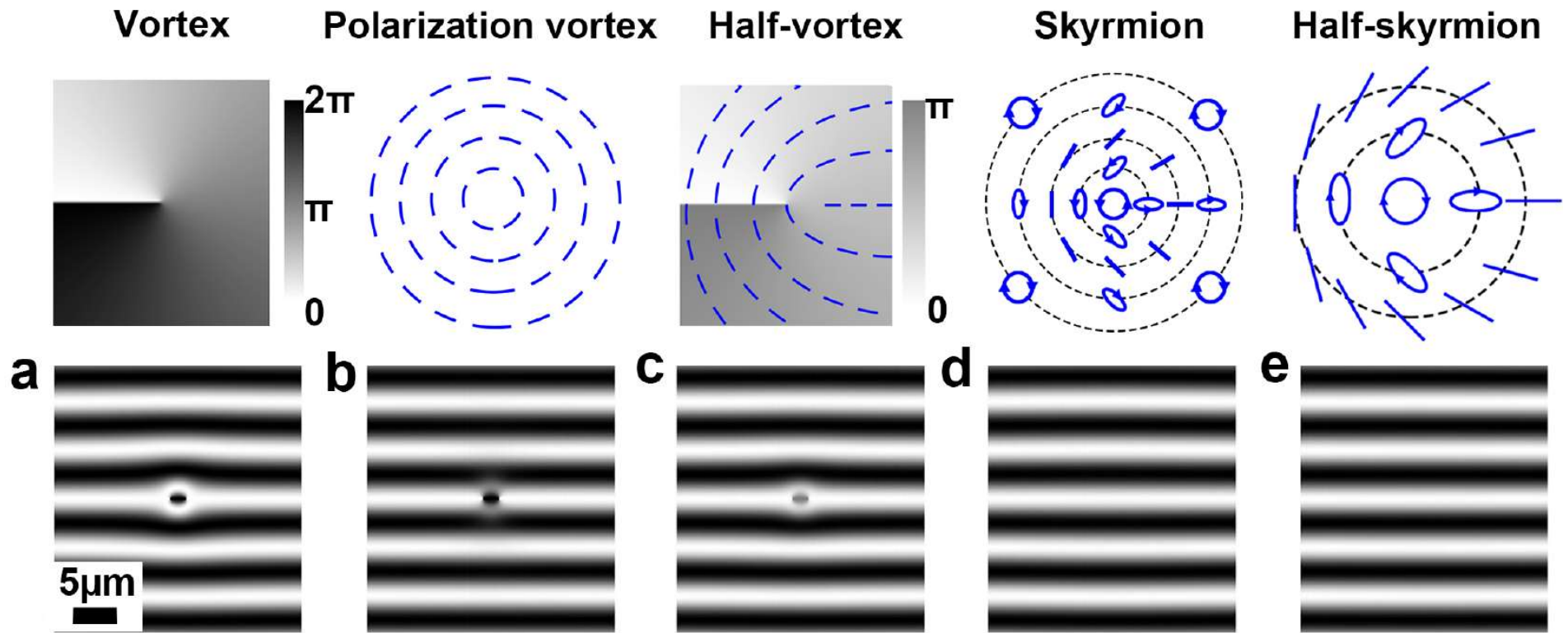
quantized vortex: phase winds by 2π around singularity point ← can be revealed as fork in interference pattern

↑

explored for vortices in **atom** condensates, **optical** vortices, and **polariton** vortices

observed dislocations in interference pattern are **not** associated with conventional phase defects:

not vortices, **not** polarization vortices, **not** half-vortices, **not** skyrmions, **not** half-skyrmions



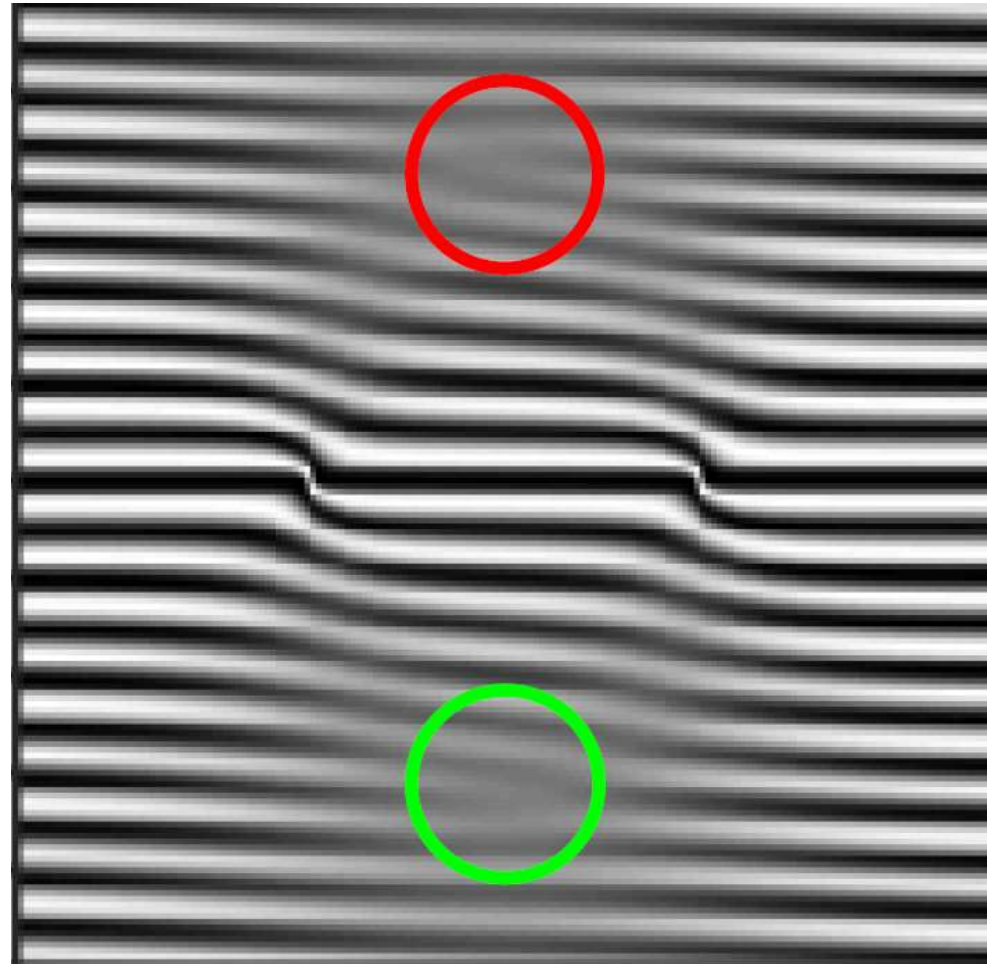
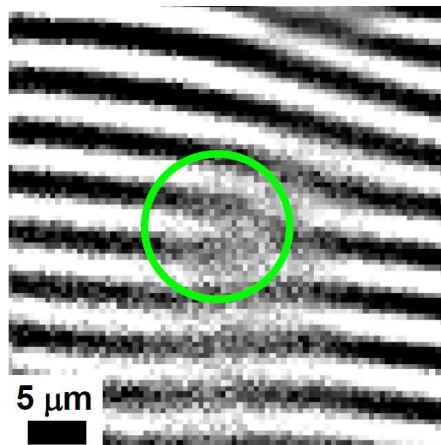
none of these simulated interference patterns is similar to experiment

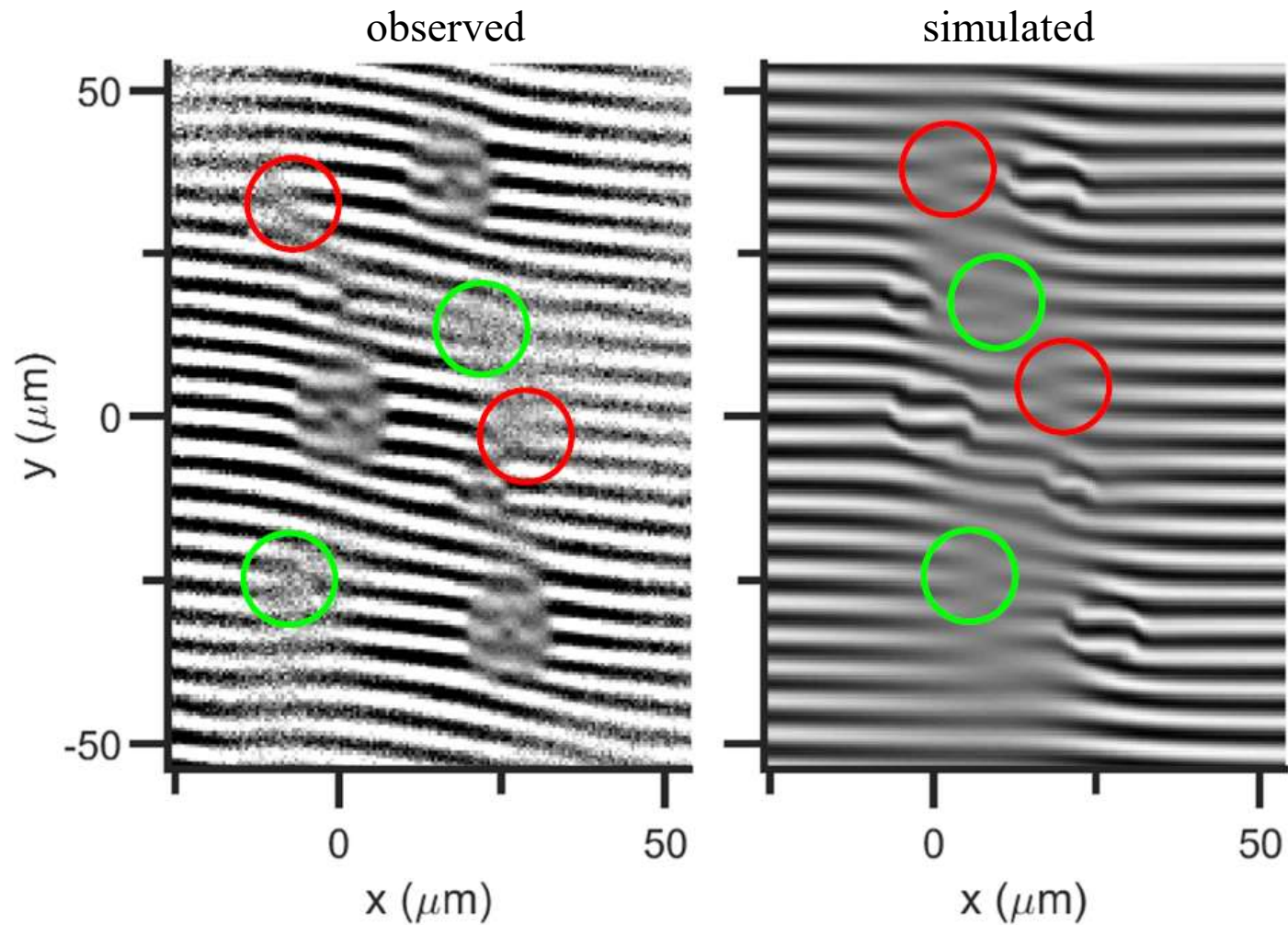
observed singularity in interference pattern is **not** associated with vortex, or polarization vortex, or half-vortex, or skyrmion, or half-skyrmion

origin of observed phase singularities in IX condensate interference patterns:

observed interference dislocations originate from converging of condensate matter waves

simulations reproduce observed “isolated” interference dislocations

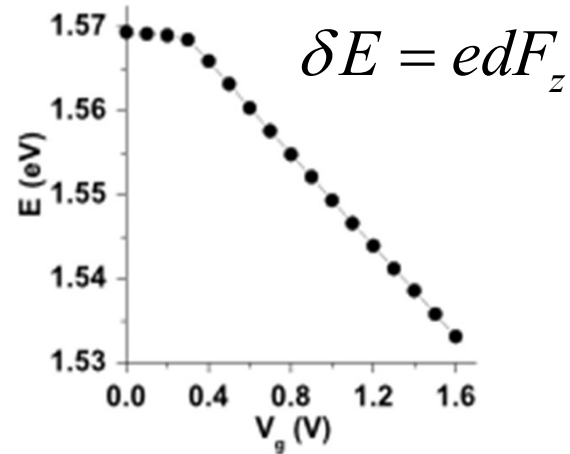
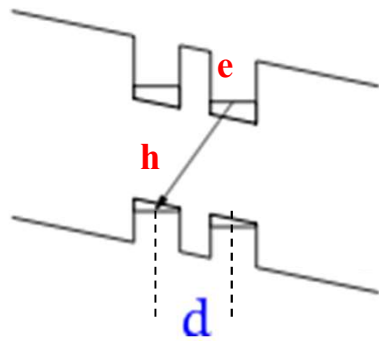




observed and simulated complex interference patterns with multiple
interference dislocations \leftrightarrow [converging of condensate matter waves](#)
and phase domains \leftrightarrow [Pancharatnam-Berry phase](#)

Excitonic devices

Potential energy of IXs can be controlled by voltage



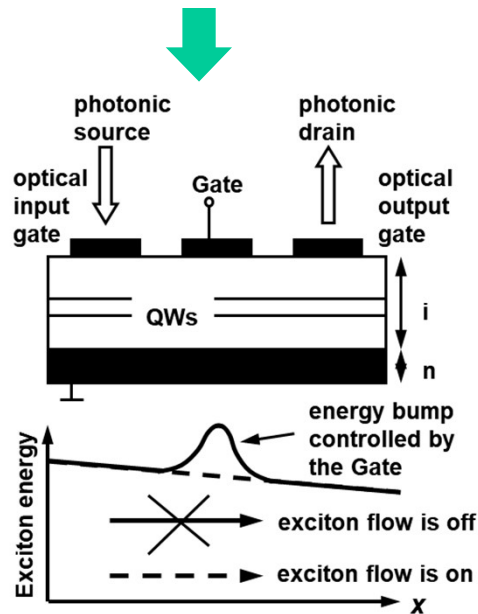
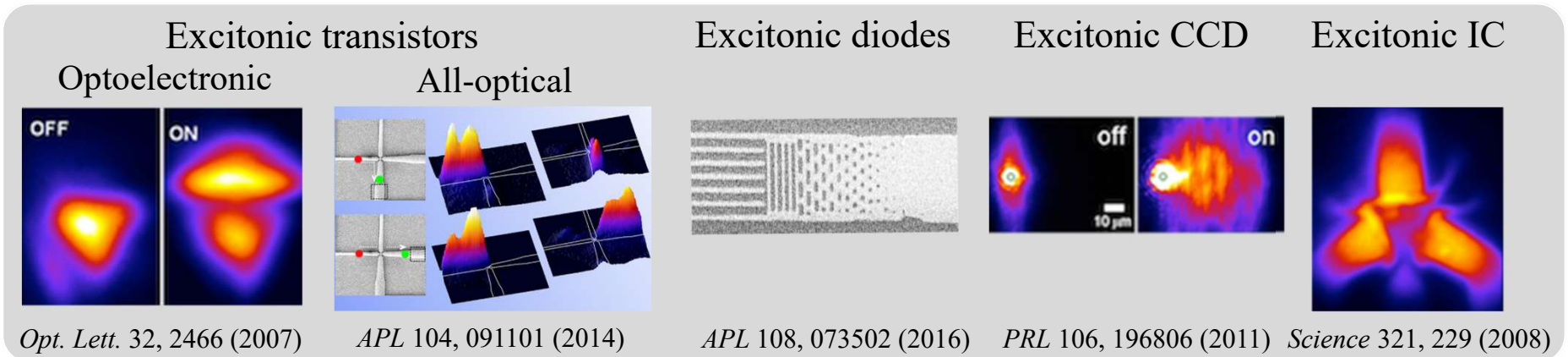
in-plane potential landscapes
for IXs can be created
and controlled by voltage

excitonic devices
operate with excitons
in place of electrons

Proof of principle demonstration of excitonic devices

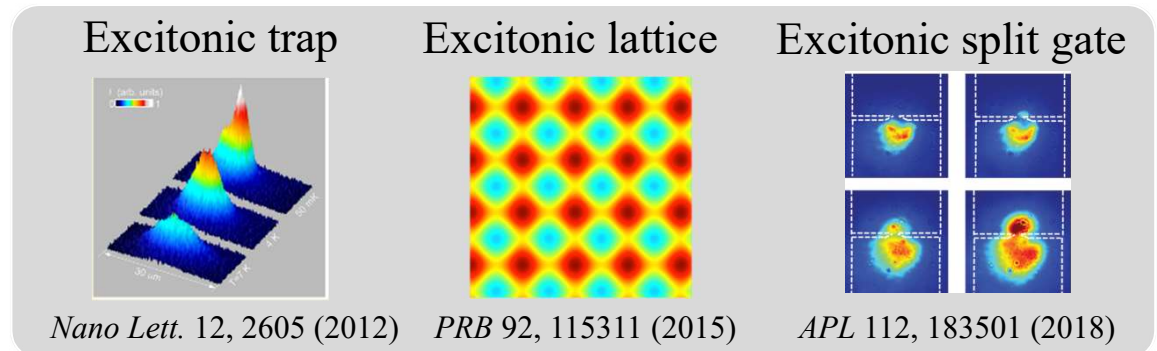
Excitonic Circuit Devices

- low-energy
- seamless coupling to optical communication



Excitonic devices for basic study

- mesoscopic of bosons

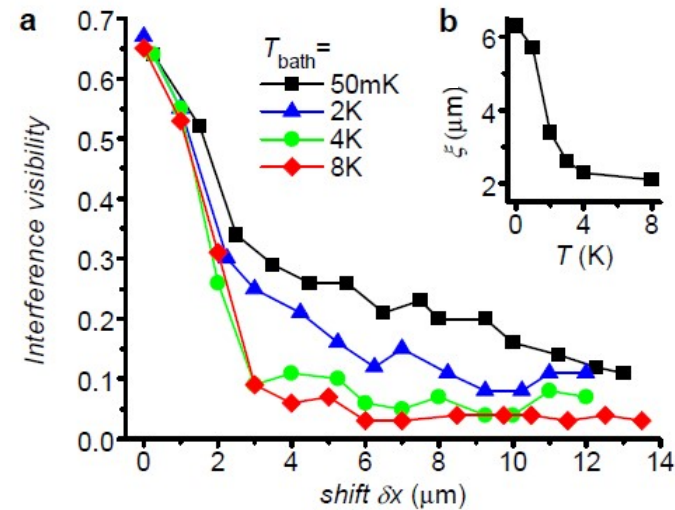
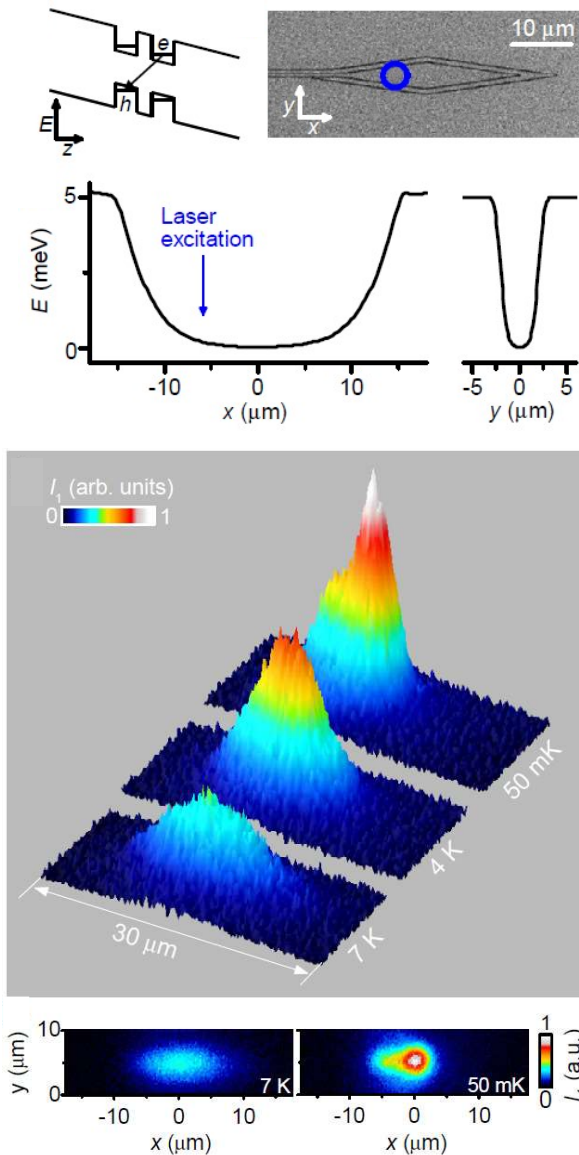


Condensation of IXs in a trap

diamond-shaped trap

With lowering T

- IXs condense at the trap bottom
- IX spontaneous coherence emerges



Experiment: A.A. High, J.R. Leonard, M. Remeika, L.V. Butov, M. Hanson, A.C. Gossard, *Nano Lett.* 12, 2605 (2012)

Theory: S.V. Lobanov, N.A. Gippius, L.V. Butov, *Phys. Rev. B* 94, 245401 (2016)

Agreement between experiment and theory:
measured IX condensation is adequately described by
(quasi)equilibrium BEC of interacting bosons

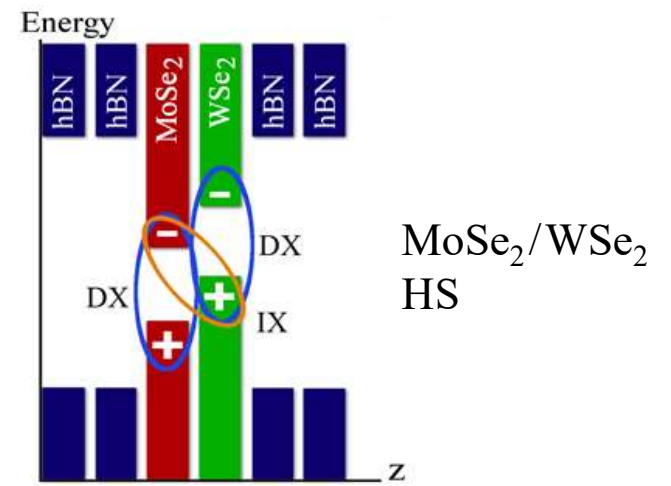
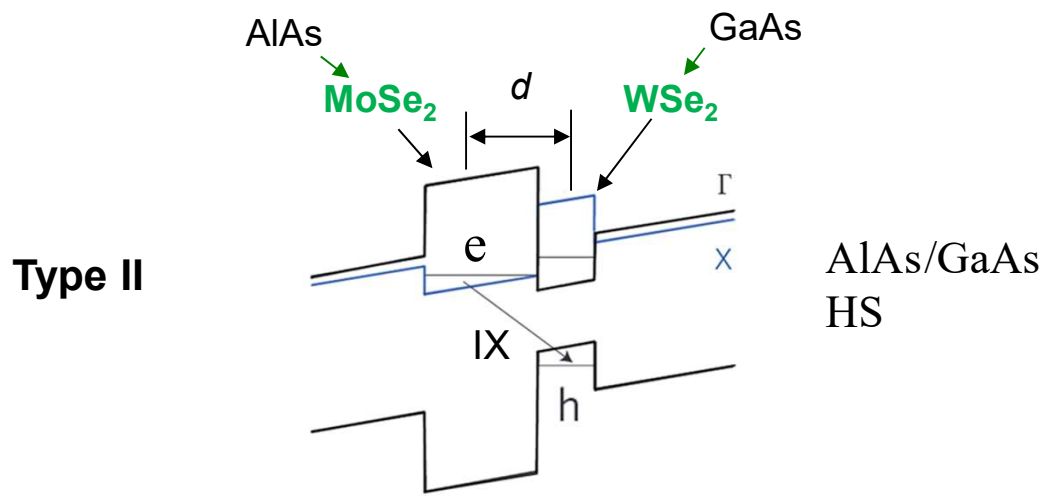
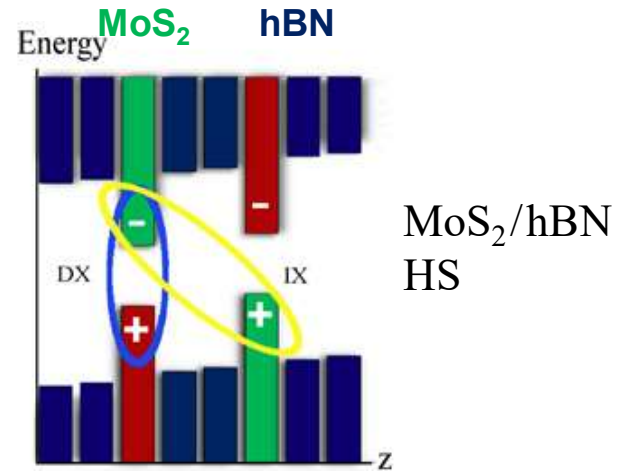
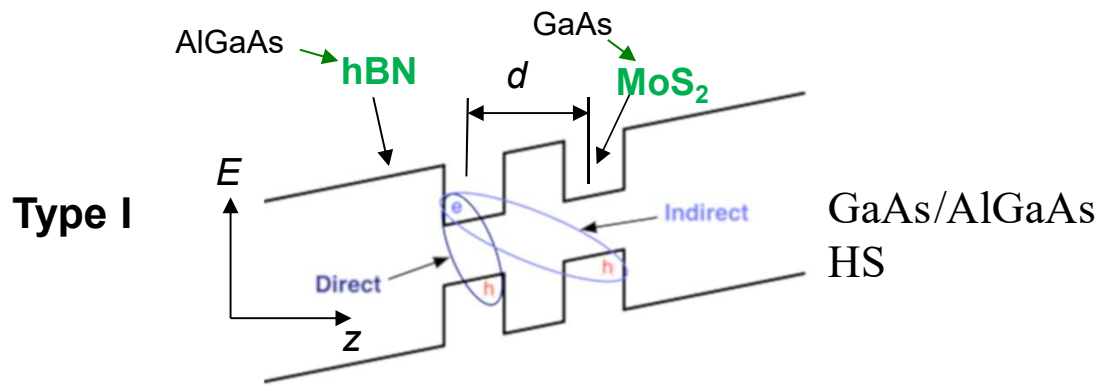
IXs in van der Waals TMD heterostructures

IX is composed of an electron and a hole confined in separated layers

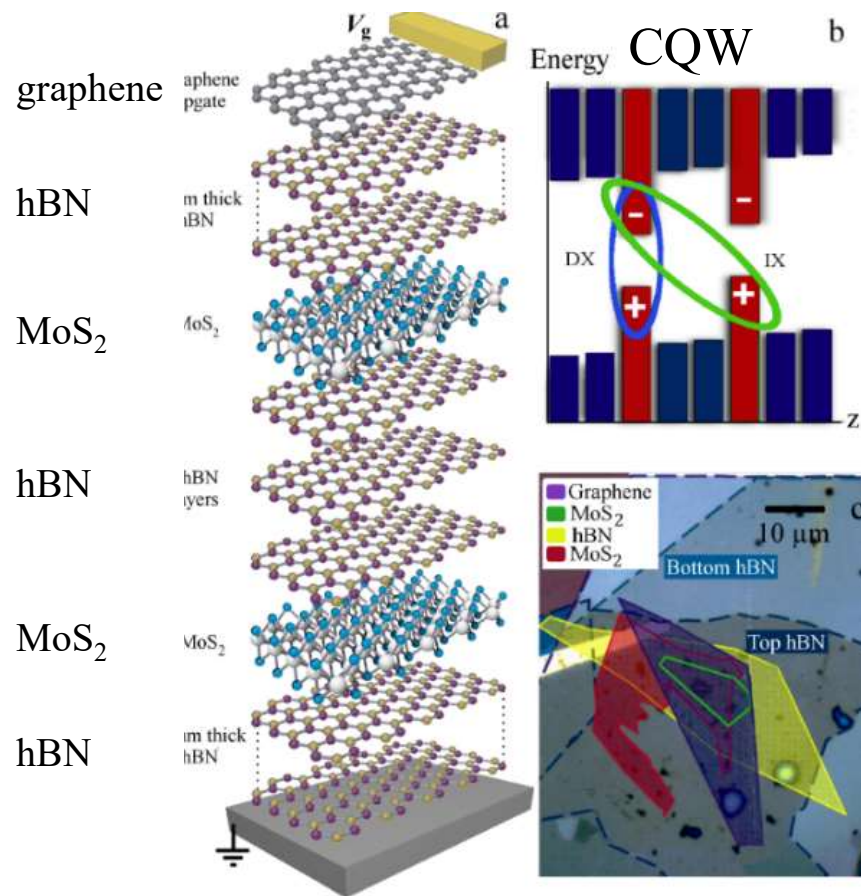
IX heterostructures (HS)

GaAs HS
Low-disorder system

Van der Waals TMD HS
IXs are robust at room T



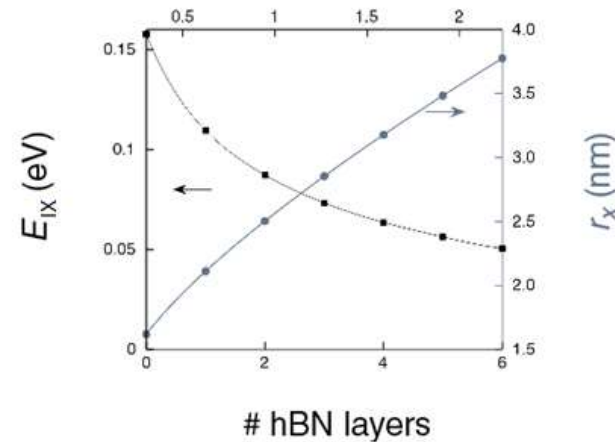
Van der Waals TMD heterostructures for high-T IX condensation



Recent calculations: $E_b \sim 350 \text{ meV} > 10 T_{\text{room}}$

Deilmann, Thygesen, *Nano Lett.* 18, 1460 (2018)

Centre-to-centre distance (nm)



IXs are robust at room T
↓
room-T IX devices

predicted high-T superfluidity in IXs in TMD heterostructures

$$T_0 = \frac{2\pi\hbar^2}{m_x} n = \frac{4\pi m_e m_h}{m_x^2} (n a_x^2) R y_x$$

$$n^{\max} a_x^2 \sim 0.02$$

$$T_0^{\max} \sim 0.06 R y_x$$

high $R y_x \rightarrow$ high T_0

M.M. Fogler, L.V. Butov, K.S. Novoselov, *Nature Commun.* 5, 4555 (2014)

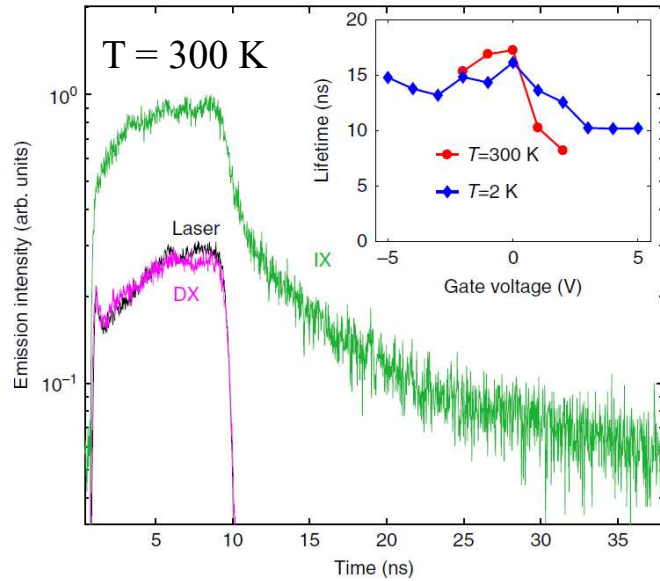
**IXs at room temperature
in van der Waals TMD heterostructures**

E.V. Calman, M.M. Fogler, L.V. Butov,
S. Hu, A. Mishchenko, A.K. Geim,
Nature Commun. 9, 1895 (2018)

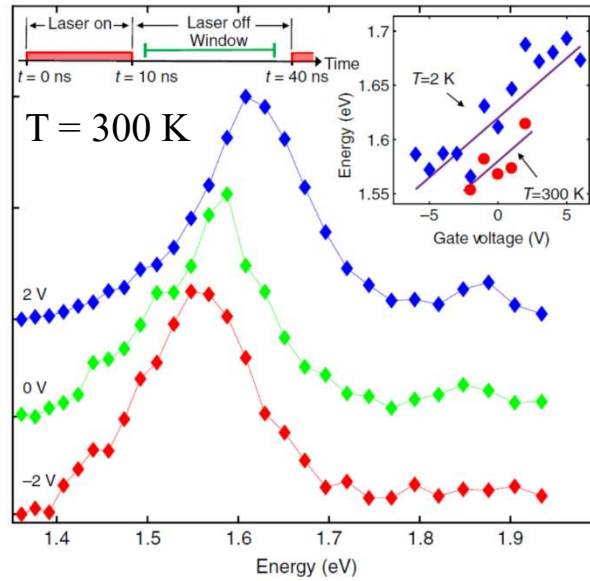
IXs in MoS₂/hBN van der Waals TMD heterostructures

basic IX properties:

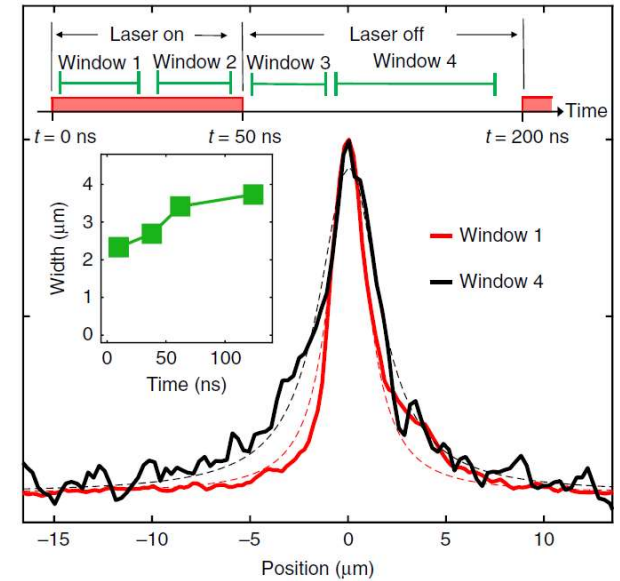
long lifetime



control of energy by voltage



transport



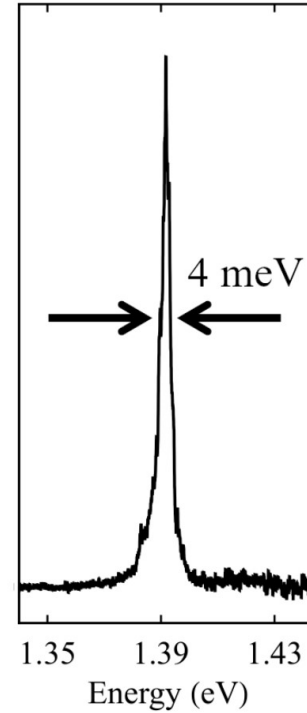
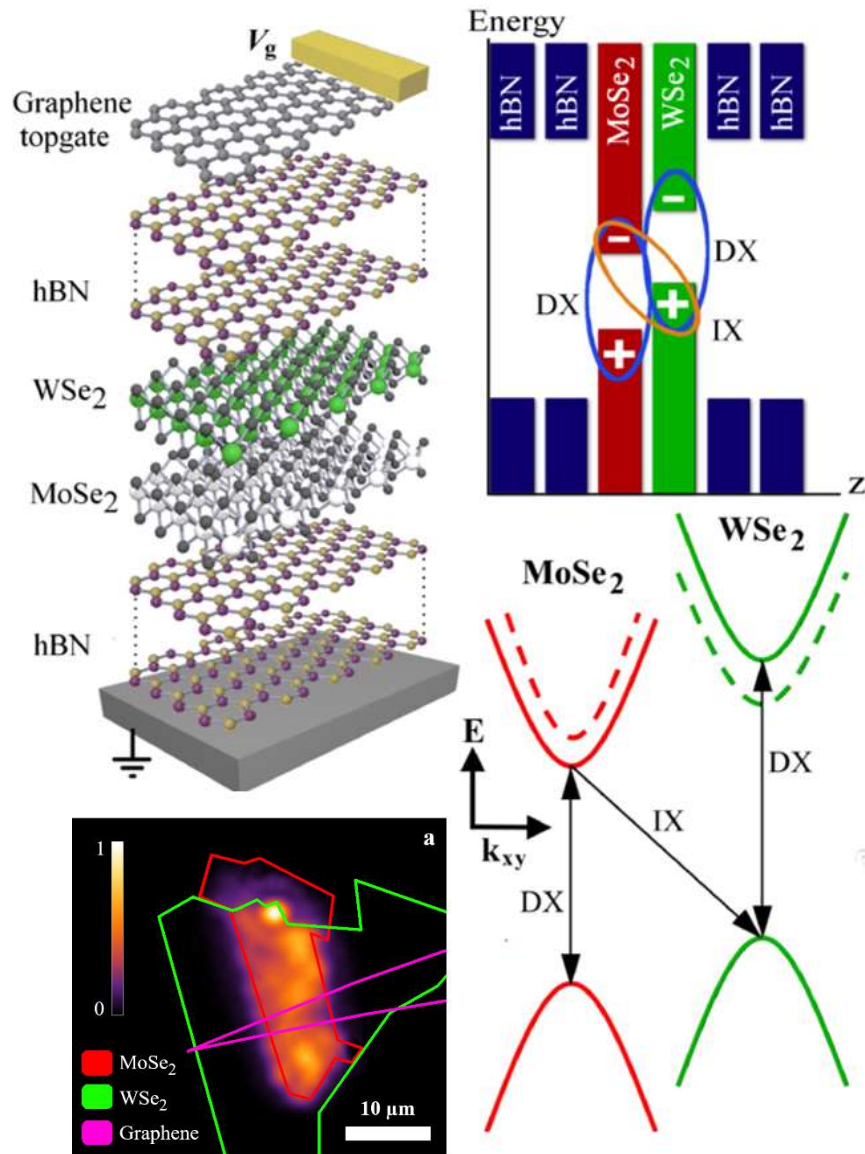
problems: ~ 100 meV broad line

few μm short-range transport

IXs in MoSe₂/WSe₂ van der Waals TMD heterostructures

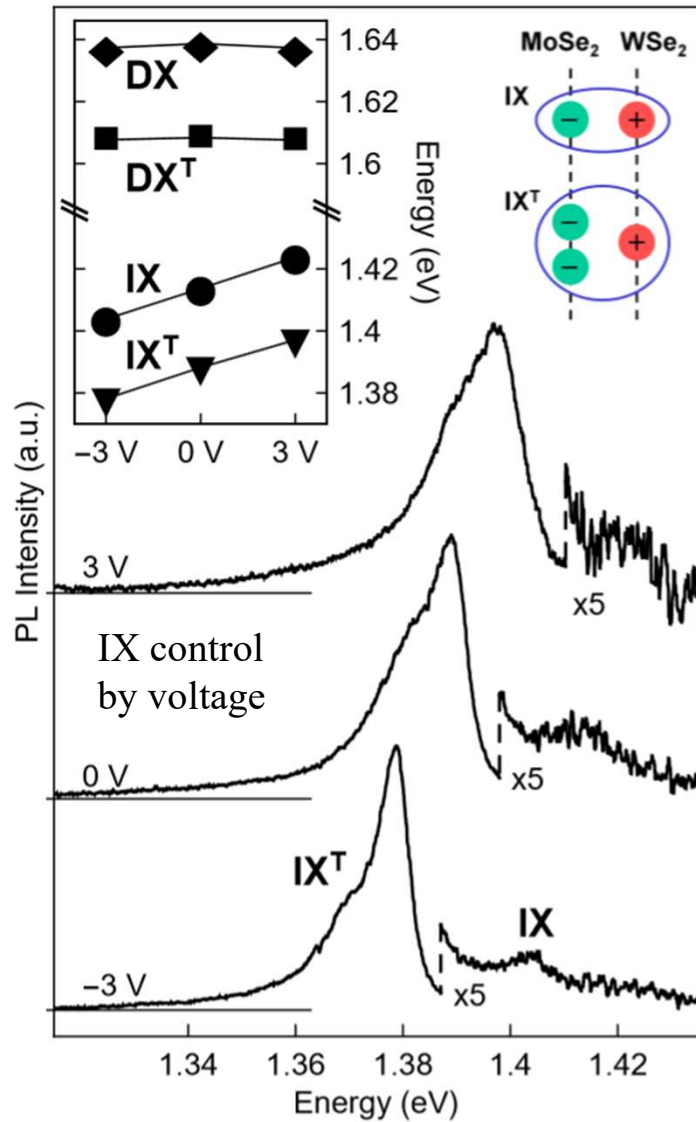
E.V. Calman, L.H. Fowler-Gerace, D.J. Choksy, L.V. Butov,
D.E. Nikonov, I.A. Young, S. Hu, A. Mishchenko, A.K. Geim,
Nano Lett. (2020)

IXs in MoSe₂/WSe₂ van der Waals TMD heterostructures

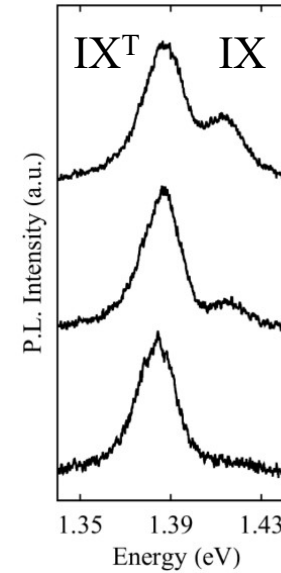


narrow IX linewidth
down to 4 meV

Indirect trions



narrow linewidth
 ↓
 resolve **two peaks** of indirect luminescence



- high energy peak – neutral indirect exciton (IX)
- low energy peak – charged indirect exciton, i.e. indirect trion (IX^T)

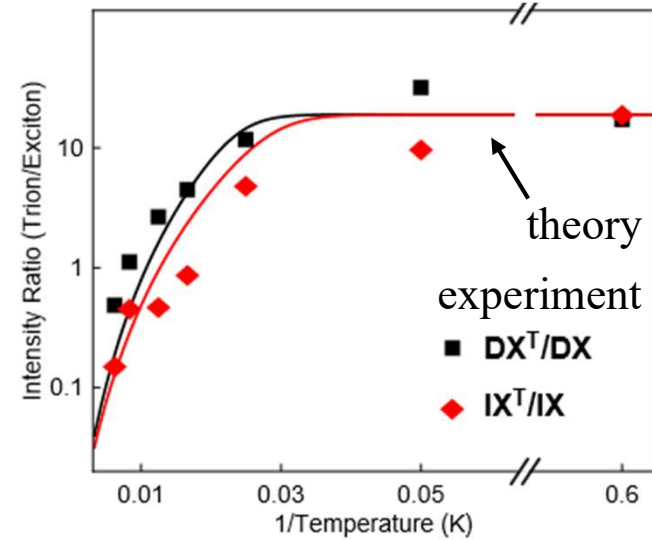
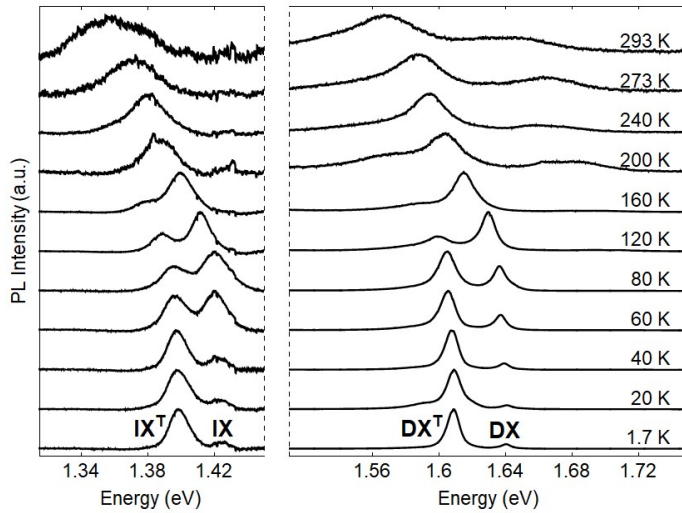
measured binding energy of indirect trion is in agreement with theory

$$E_b = 28 \text{ meV}$$

theory: Deilmann, Thygesen, *Nano Lett.* 18, 1460 (2018)

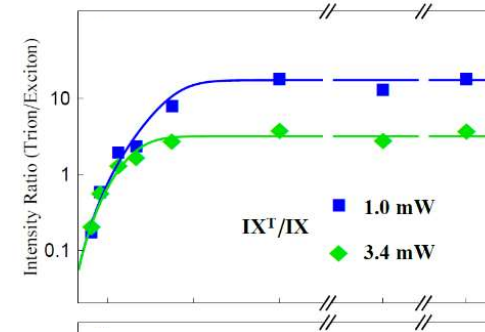
Temperature dependence

observed
IX up to
room T



measured exciton/trion temperature dependence
is in agreement with theory of trions

checked in different samples,
spots, excitation powers

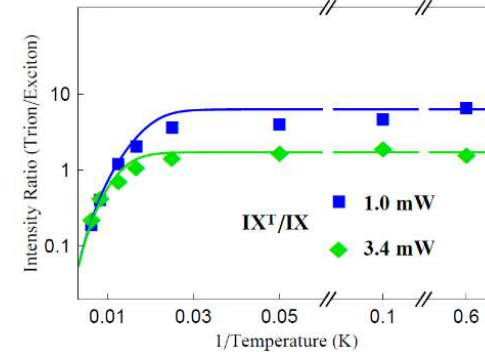


key feature of trion luminescence

is saturation at low T ← observed

↑

at low T: $n_{\text{trion}} \rightarrow n_{\text{background e}}$



Summary

- Spontaneous coherence and condensation of IXs
- Phenomena in IX condensate
 - Density wave, commensurability effect
 - Spin textures
 - Pancharatnam-Berry phase, coherent spin transport
 - Phase singularities, interference dislocations
- IXs in van der Waals heterostructures
 - Opportunity to realize high-T IX condensation
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