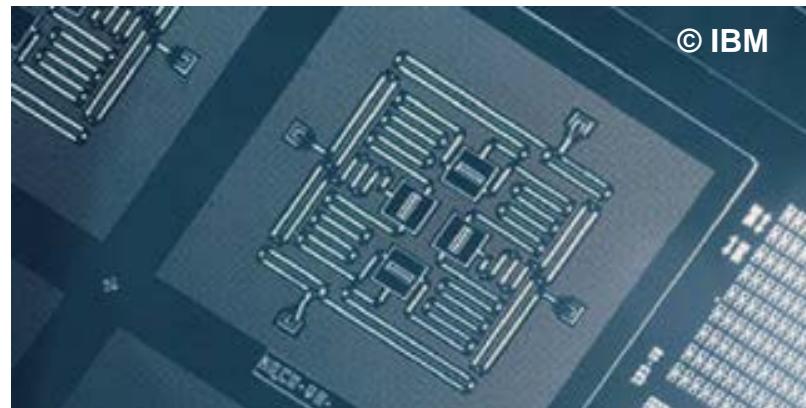
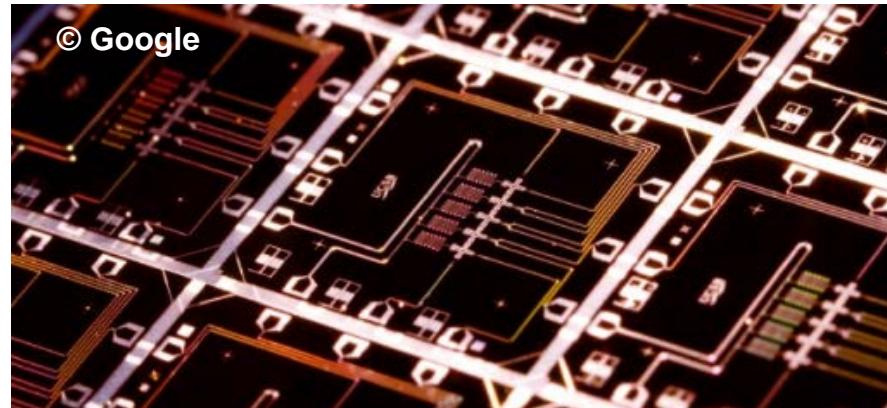


Сверхпроводящие квантовые процессоры: физика, технология и перспективы

Алексей Устинов

Технологический институт Карлсруэ
НИТУ МИСиС
Российский квантовый центр



Superconducting quantum machines

Components

- resonators
- Josephson junctions

Qubits

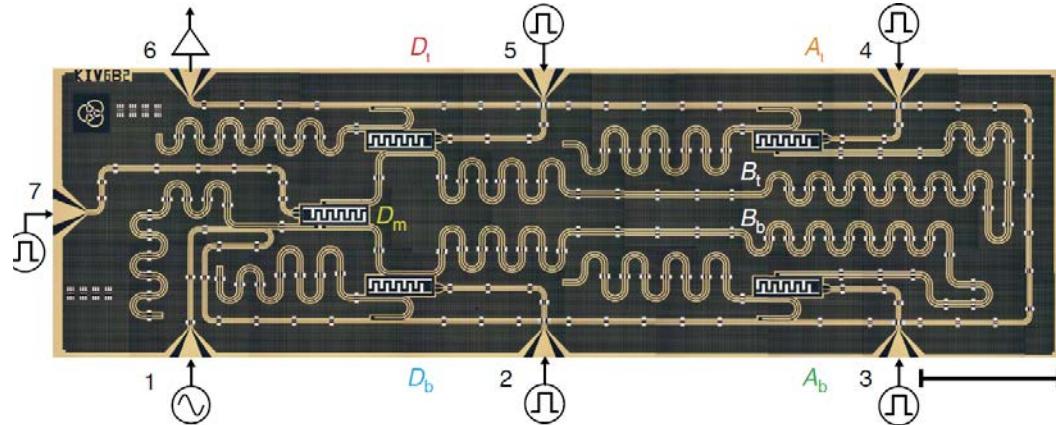
- phase, charge, flux
- -um's and –mon's
- decoherence

Quantum processors

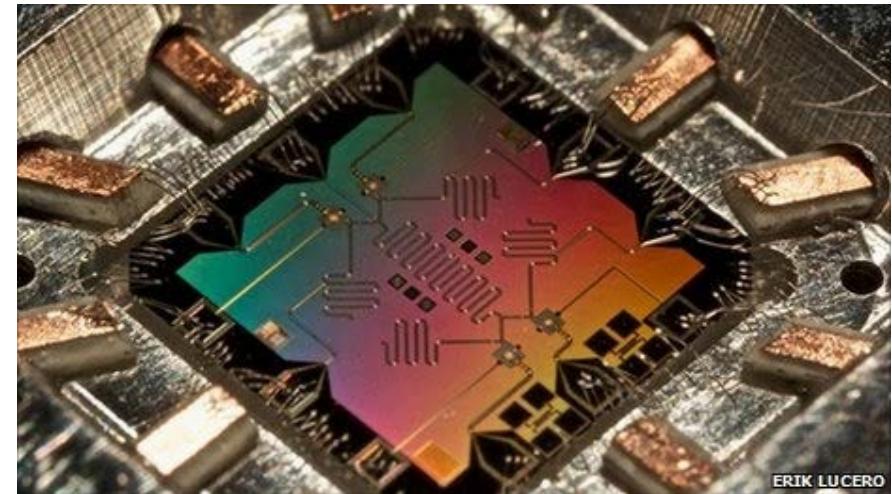
- simple algorithms
- error correction
- simulators vs computers

Not in this talk

- quantum microwaves
- single-photon detectors



© DiCarlo lab TU Delft

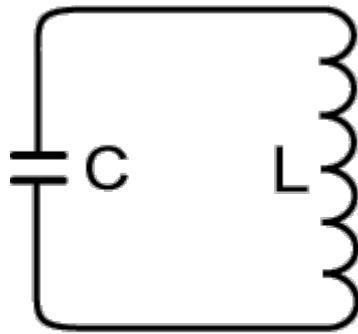


ERIK LUCERO

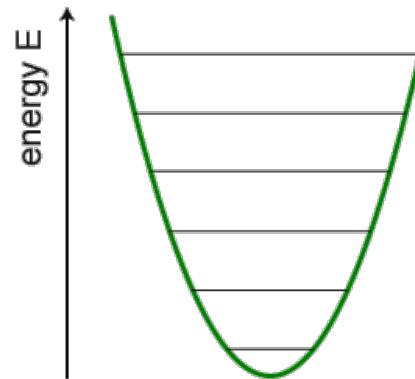
© Martinis lab UCSB/Google

Electromagnetic resonator as a quantum system

■ LC - resonator



■ potential energy

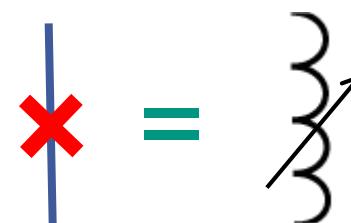


$$\hat{H} = \frac{\hat{Q}^2}{2C} + \frac{\hat{\Phi}^2}{2L}$$

conjugate operators

$$[\hat{Q}, \hat{\Phi}] = i\hbar$$

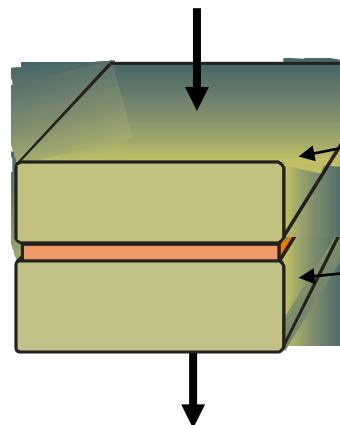
- harmonic potential: equidistant energy levels
- way to get anharmonic potential is by using a nonlinear inductor: a Josephson junction



Josephson junction



superconductor
tunnel barrier
superconductor



$$\Psi_1 = |\Psi_1| \exp(i\theta_1)$$

$$\Psi_2 = |\Psi_2| \exp(i\theta_2)$$

$$\text{superconducting phase difference: } \varphi = \theta_1 - \theta_2$$

Josephson
relations

$$\left\{ \begin{array}{l} I_S = I_C \sin \varphi \\ V = \frac{\hbar}{2e} \frac{d\varphi}{dt} \end{array} \right.$$

Josephson inductance

$$L_J = V \frac{dI_S}{dt} = \frac{\Phi_0}{2\pi I_C \cos \varphi}$$



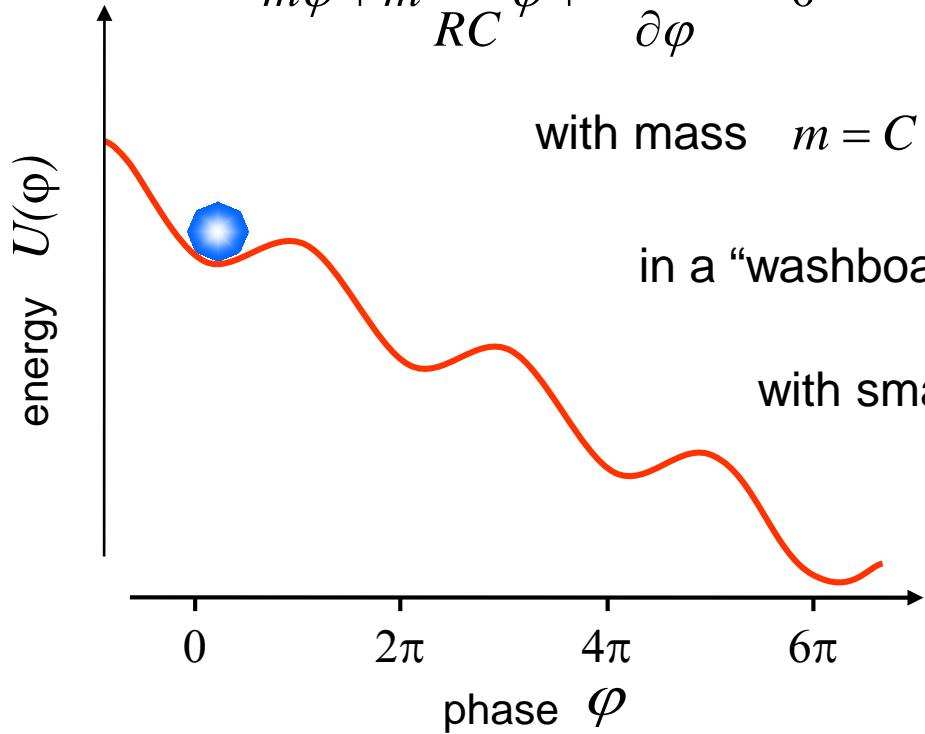
Josephson junction: washboard potential

Dynamics of a small Josephson junction is equivalent to the motion of a particle

$$m\ddot{\phi} + m\frac{1}{RC}\dot{\phi} + \frac{\partial U(\phi)}{\partial \phi} = 0$$

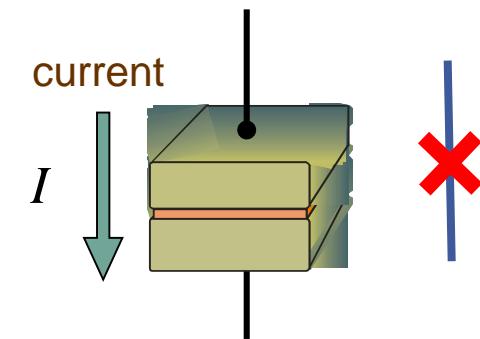
with mass $m = C\left(\frac{\Phi_0}{2\pi}\right)^2$

in a “washboard” potential $U(\phi) = -\frac{I_c\Phi_0}{2\pi}\left(\frac{I\phi}{I_c} + \cos\phi\right)$

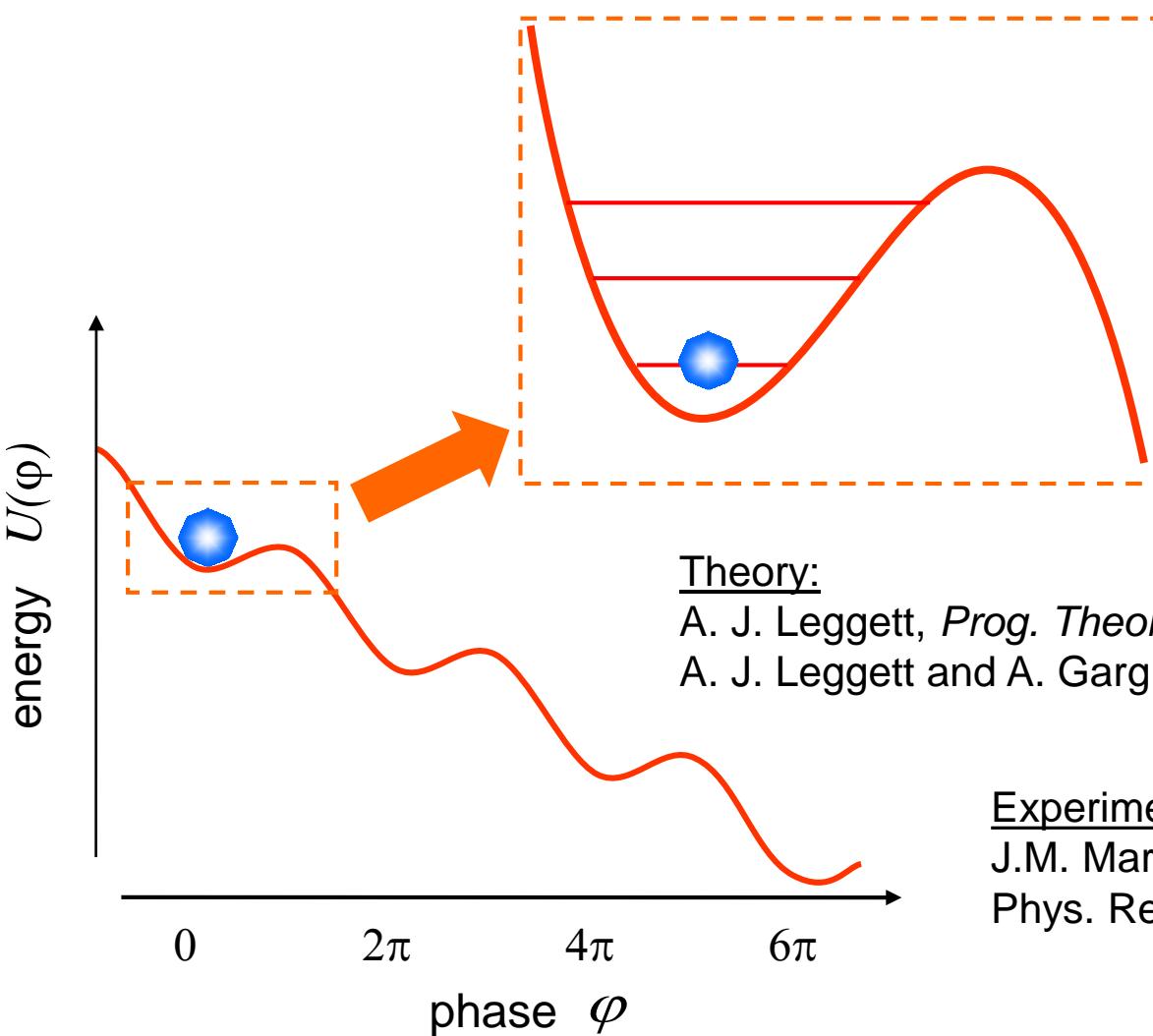


with small-amplitude oscillation frequency

$$\omega(I) = \omega_0 \sqrt[4]{1 - \frac{I^2}{I_c^2}} \quad ; \quad \omega_0 = \sqrt{\frac{2\pi I_c}{\Phi_0 C_J}}$$



Macroscopic quantum system



Theory:

A. J. Leggett, *Prog. Theor. Phys. Suppl.* **69**, 80 (1980).
A. J. Leggett and A. Garg, *Phys. Rev. Lett.* **54**, 857(1985).

Experiment:

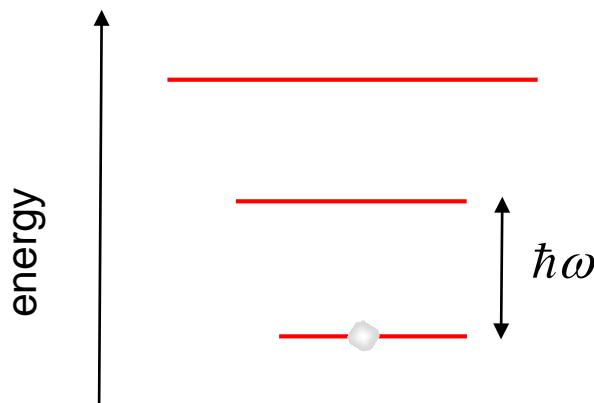
J.M. Martinis, M. H. Devoret, and J. Clarke,
Phys. Rev. B **35**, 4682 (1987).

By reducing the capacitance C the 'mass'

$$m = C \left(\frac{\Phi_0}{2\pi} \right)^2$$

can be made small

How low should be the temperature?



To see the energy levels we need $\hbar\omega \gg k_B T$.

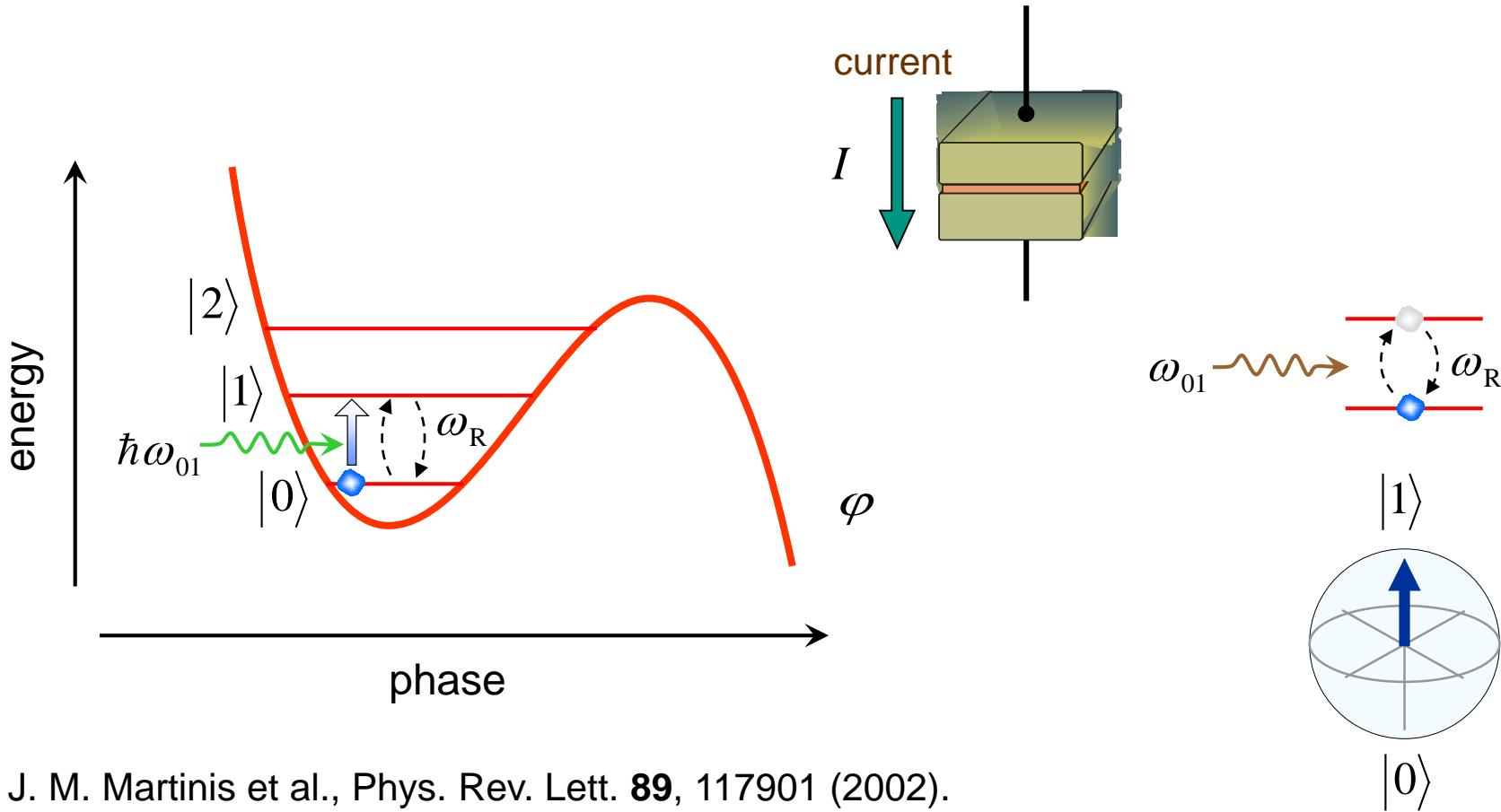
For the level separation frequency of

$$f = \frac{\omega}{2\pi} = 10 \text{ GHz}$$

the condition $T = \frac{\hbar\omega}{k_B}$ corresponds to $T \approx 0.48 \text{ K}$

i.e. $1 \text{ GHz} \longleftrightarrow \sim 50 \text{ mK}$

Josephson phase qubit



J. M. Martinis et al., Phys. Rev. Lett. **89**, 117901 (2002).

R. McDermott et al., Science **307**, 1299 (2005)

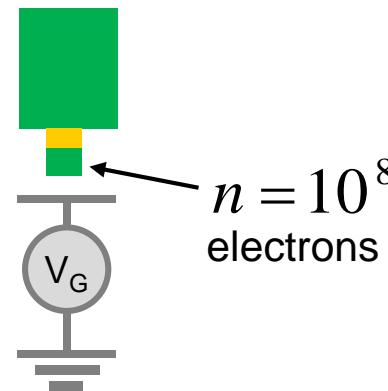
Flux and charge: Two extremes

Uncertainty relation for a superconductor: $\Delta n \cdot \Delta \varphi \geq 1$

Charging energy $E_C = \frac{e^2}{2C_J}$

charge
qubit

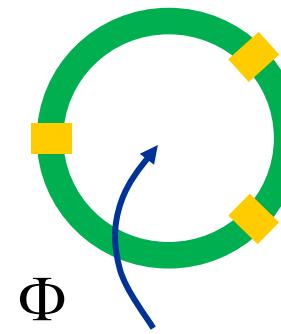
$$E_C \gg E_J$$



Josephson energy $E_J = \frac{I_c \Phi_0}{2\pi}$

flux
qubit

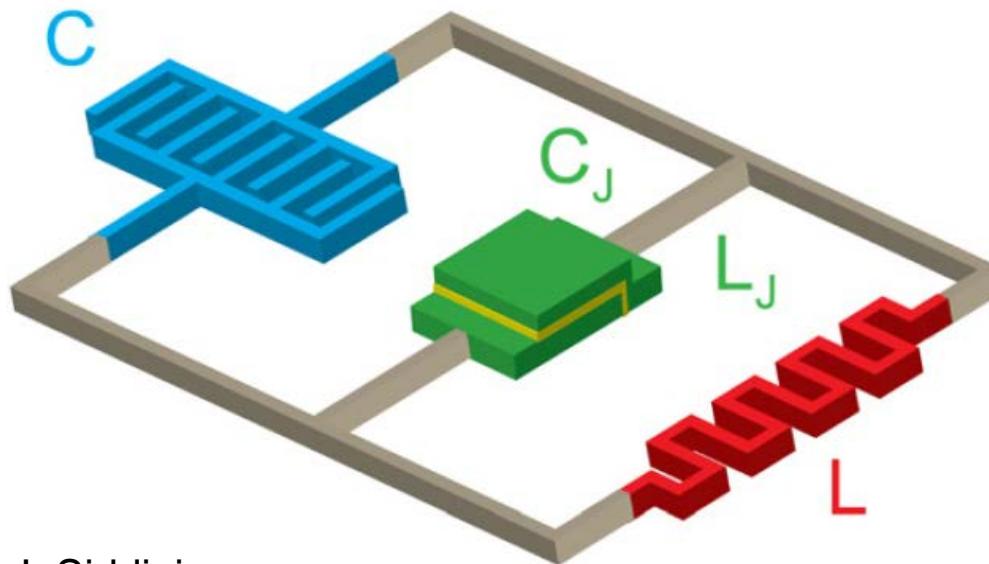
$$E_C \ll E_J$$



- Makhlin et al.,
Nature 398, 305 (1999)
- Nakamura et al.,
Nature 398, 786 (1999)

- Friedman et al.,
Nature 406, 43 (2000)
- van der Wahl et al.,
Science 290, 773 (2000)

Prototypical nonlinear equivalent circuit of a superconducting qubit



© I. Siddiqi,
Supercond. Sci. Technol.
24, 091002 (2011)

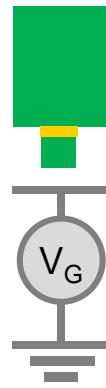
|number>
↓
|phase>

The central element is a **Josephson tunnel junction** shunted by a **capacitor** and an **inductor**. The junction has a nonlinear Josephson inductance L_J and a linear capacitance C_J .

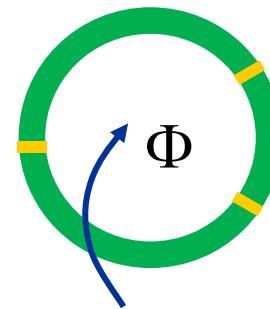
Qubit	E_J/E_C	E_L/E_J
charge	< 1	0
transmon	~ 100	0
flux	~ 100	~ 0.5
phase	~ 10^4	~ 0.2

Overview of superconducting qubits

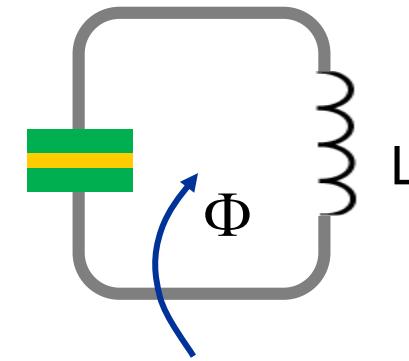
charge qubit
1999



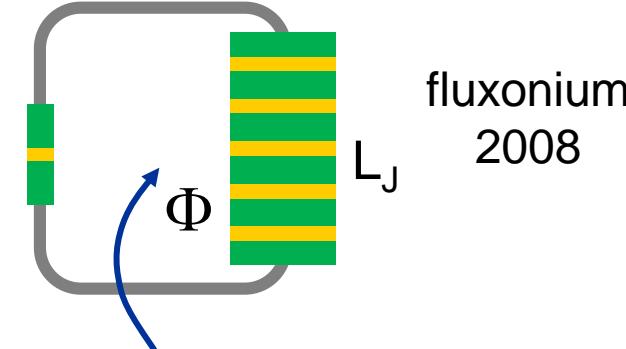
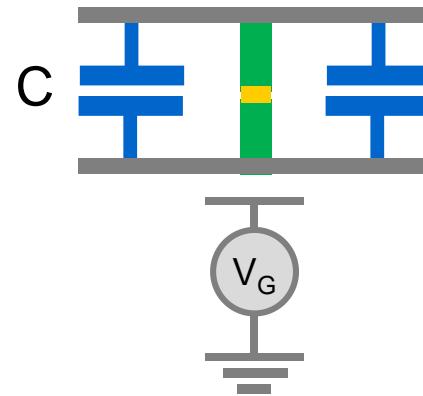
flux qubit
2000



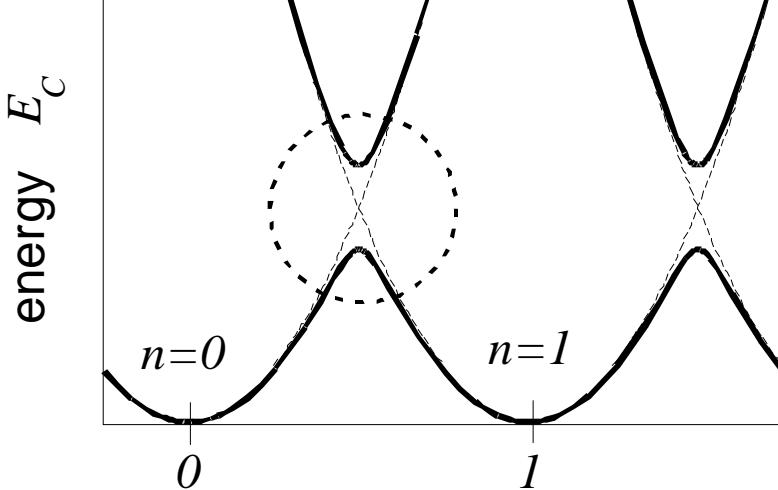
phase qubit
2002



transmon
2007



Charge qubit: NEC experiments



charge of the box $n_G = q/(2e)$

Nakamura et al., Nature

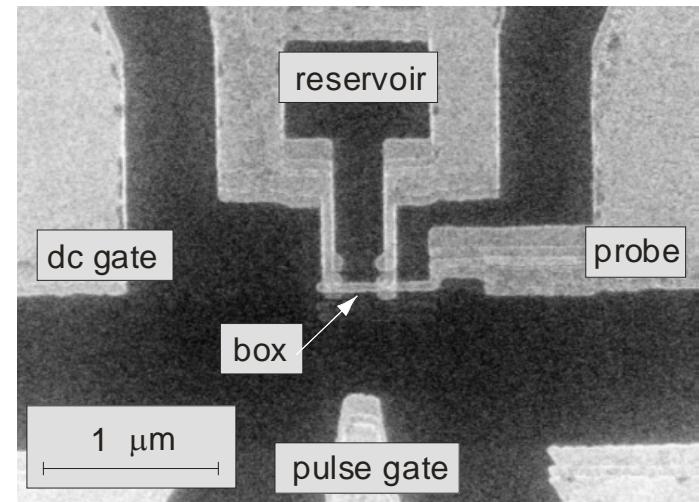
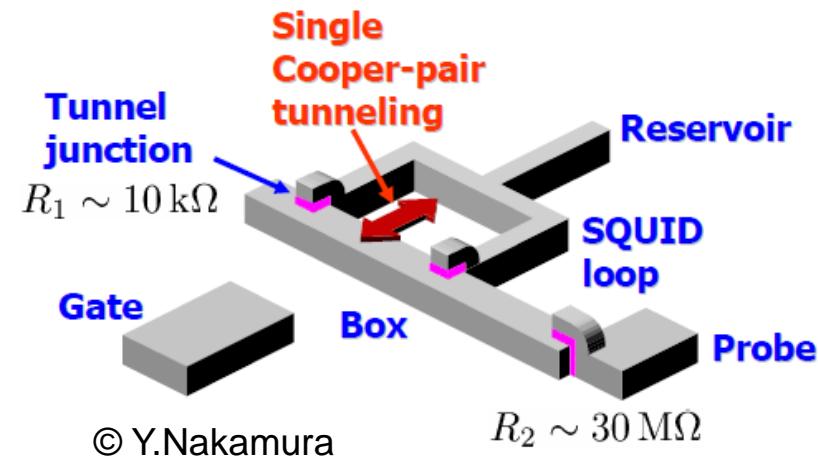
398, 786 (1999)

Nakamura et al.,

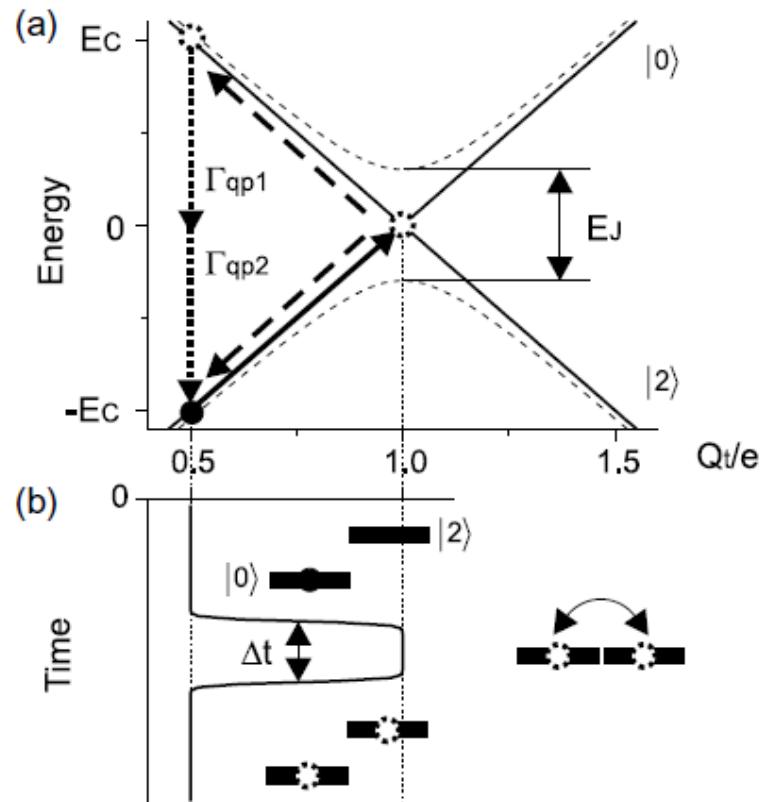
PRL 87, 246601 (2001)

Nakamura et al.,

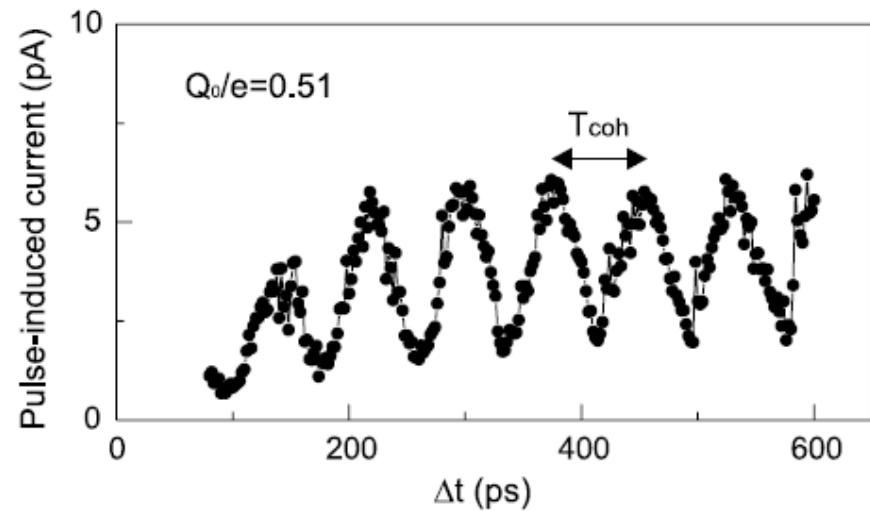
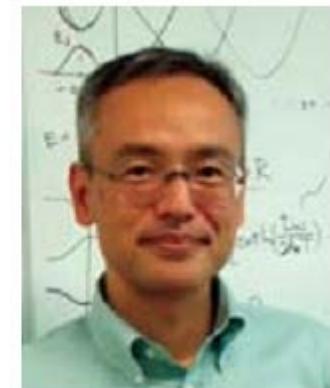
PRL 88, 047901 (2002)



Charge qubit: NEC experiments

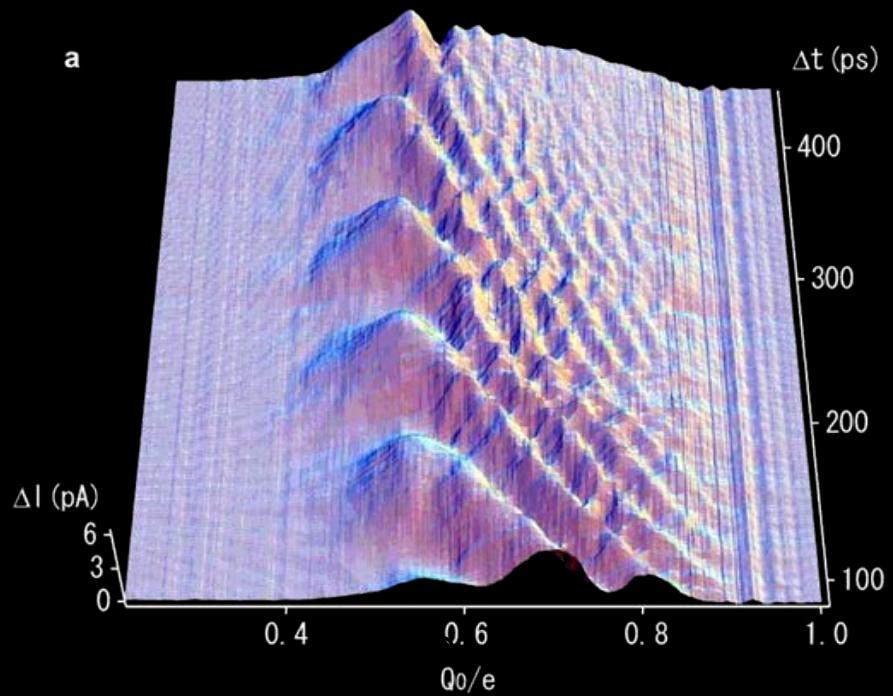


Yasunobu
Nakamura



Nakamura, Pashkin and Tsai, *Nature* 398, 786 (1999)

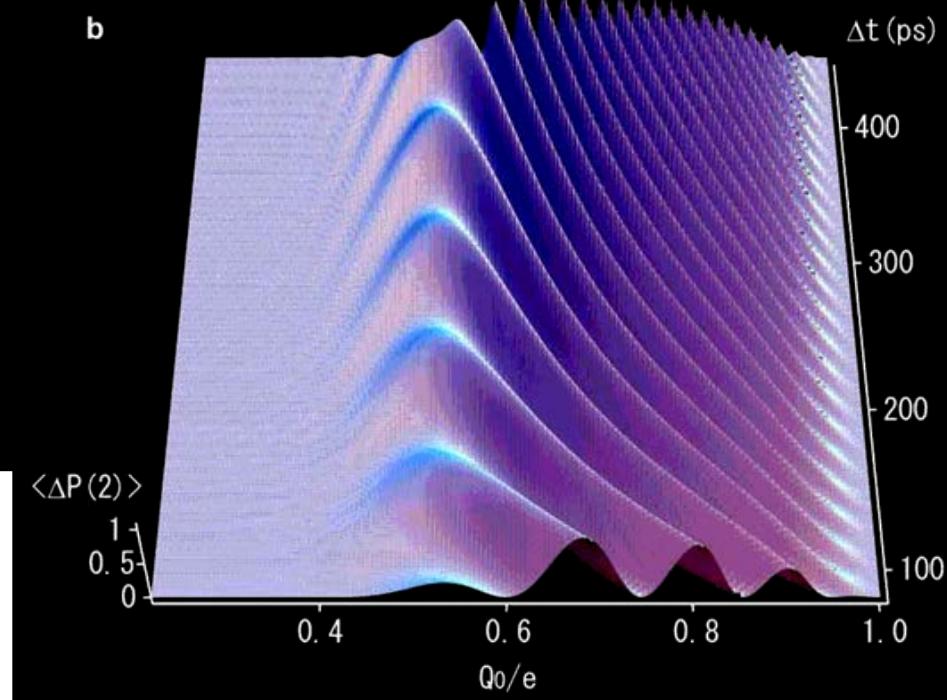
Charge qubit: NEC experiments



Measurement

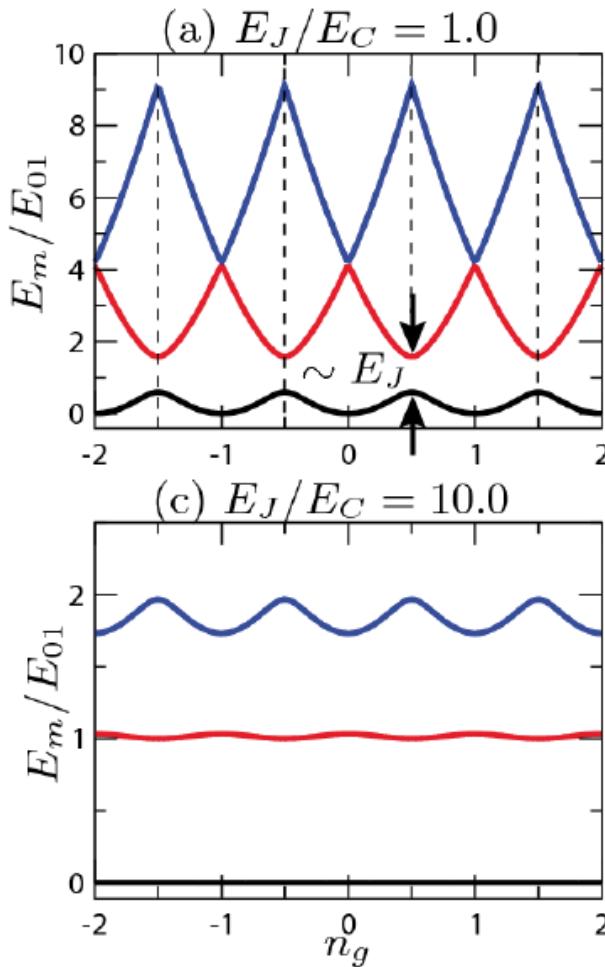
Nakamura et al.,
Nature 398, 786 (1999)

Simulation

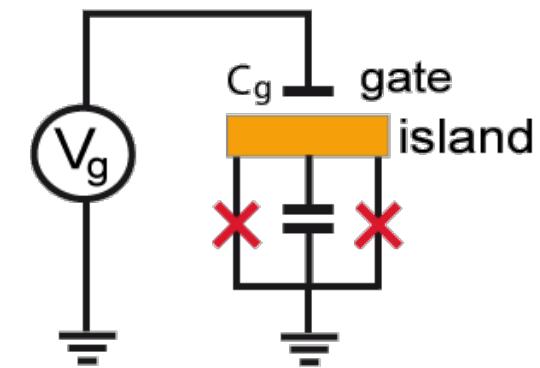


Transmon qubit

J. Koch et al., Phys. Rev. A 76, 042319 (2007)



by **increasing E_J/E_C**
(e.g. by shunting the junction with a capacitor)
the energy bands become
more flat



there is a **sweet spot**
everywhere ! 😊

but the **anharmonicity**
decreases. 😊

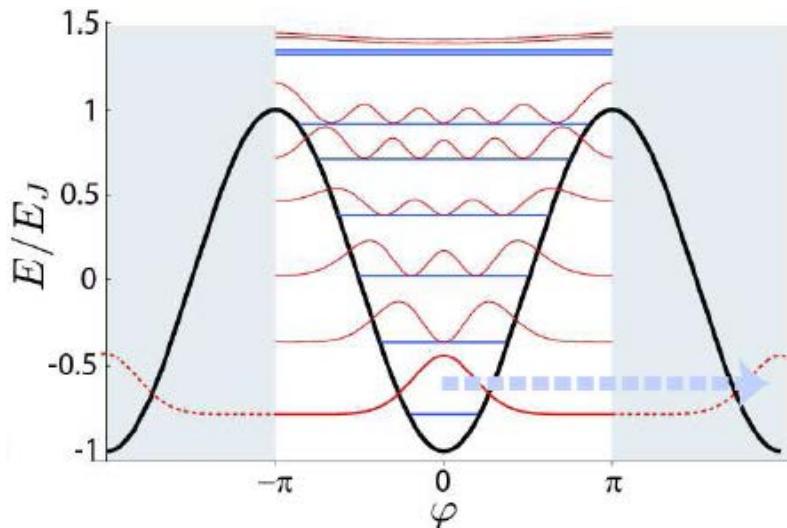
Transmon = TRANSMission-line shunted plasma oscillatiON qubit

Transmon qubit

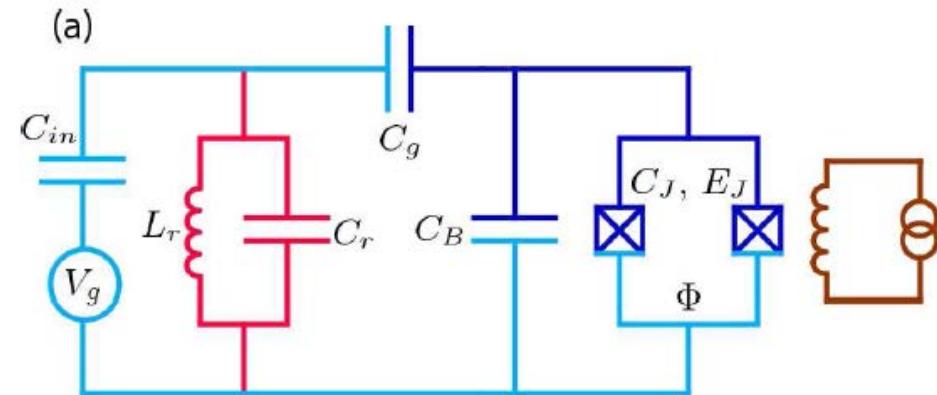
charge qubit Hamiltonian:

$$\hat{H} = E_C (\hat{N} - N_g)^2 - E_J \cos \hat{\phi}$$

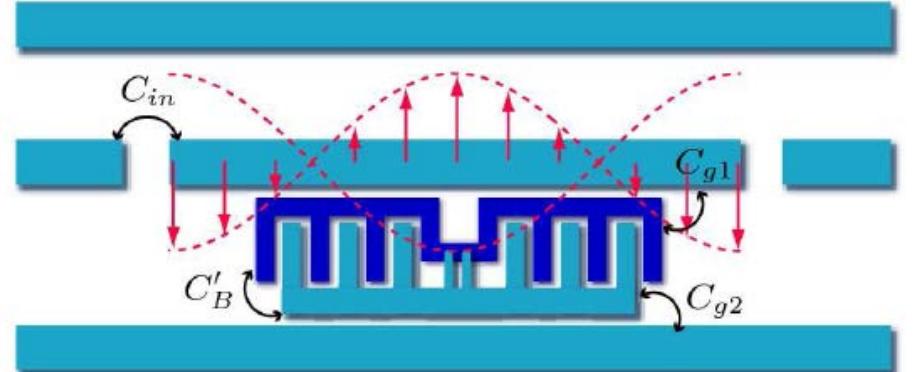
- for $E_J \gg E_C$, the eigenstates are in a cosine potential:



schematic of a transmon coupled to a resonator



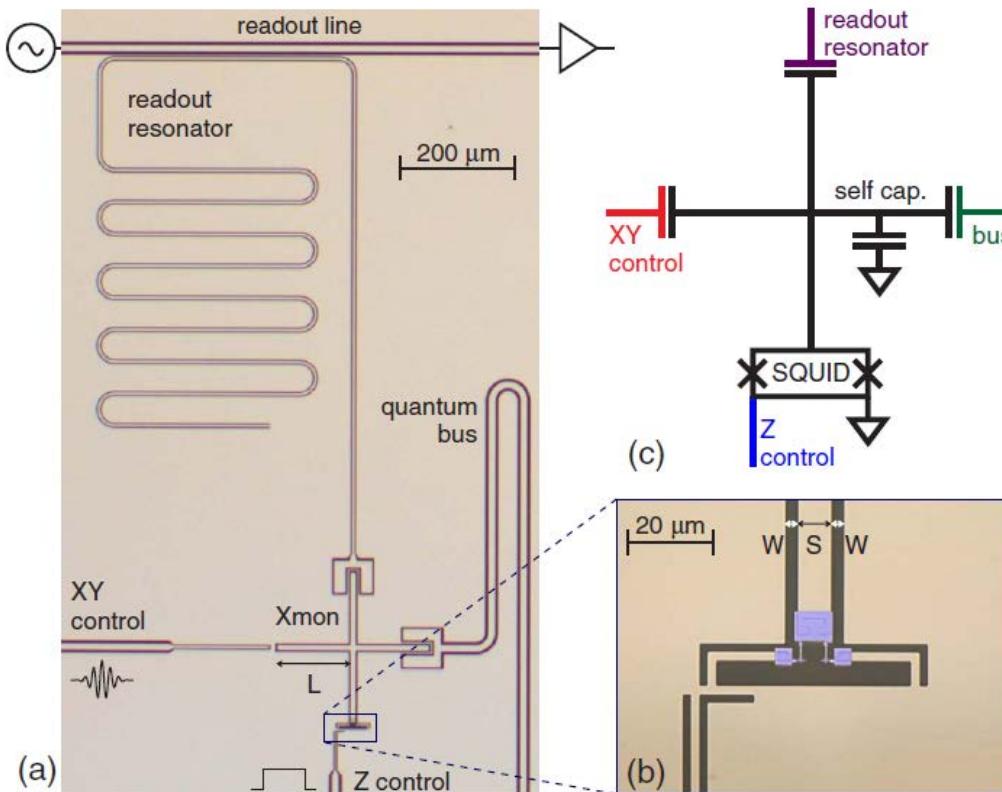
(b) transmon embedded in a coplanar resonator



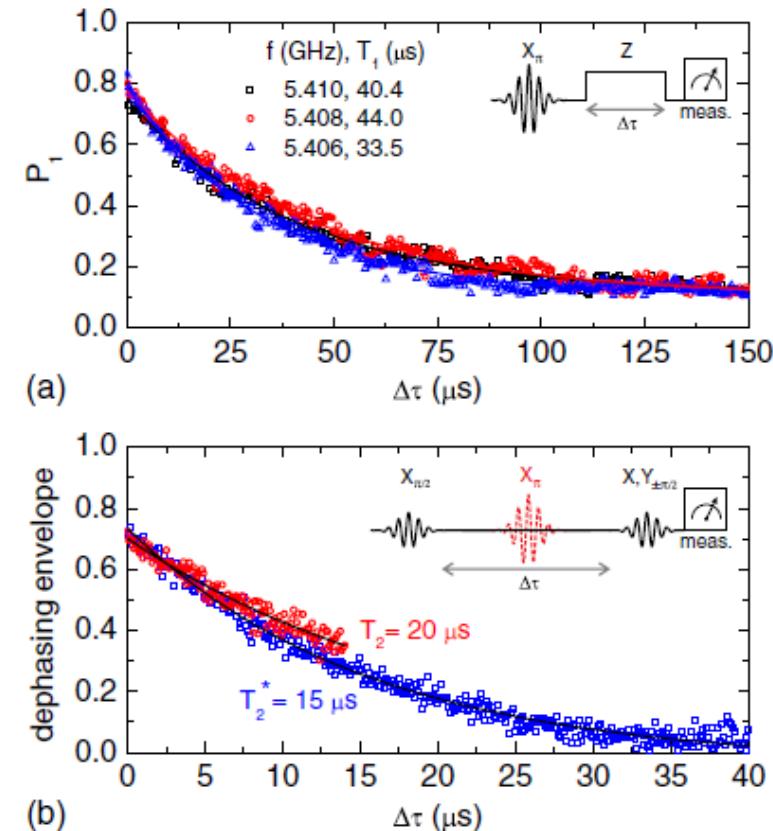
J. Koch et al., Phys. Rev. A 76, 042319 (2007)

X-mon qubits for scalable quantum integrated circuits

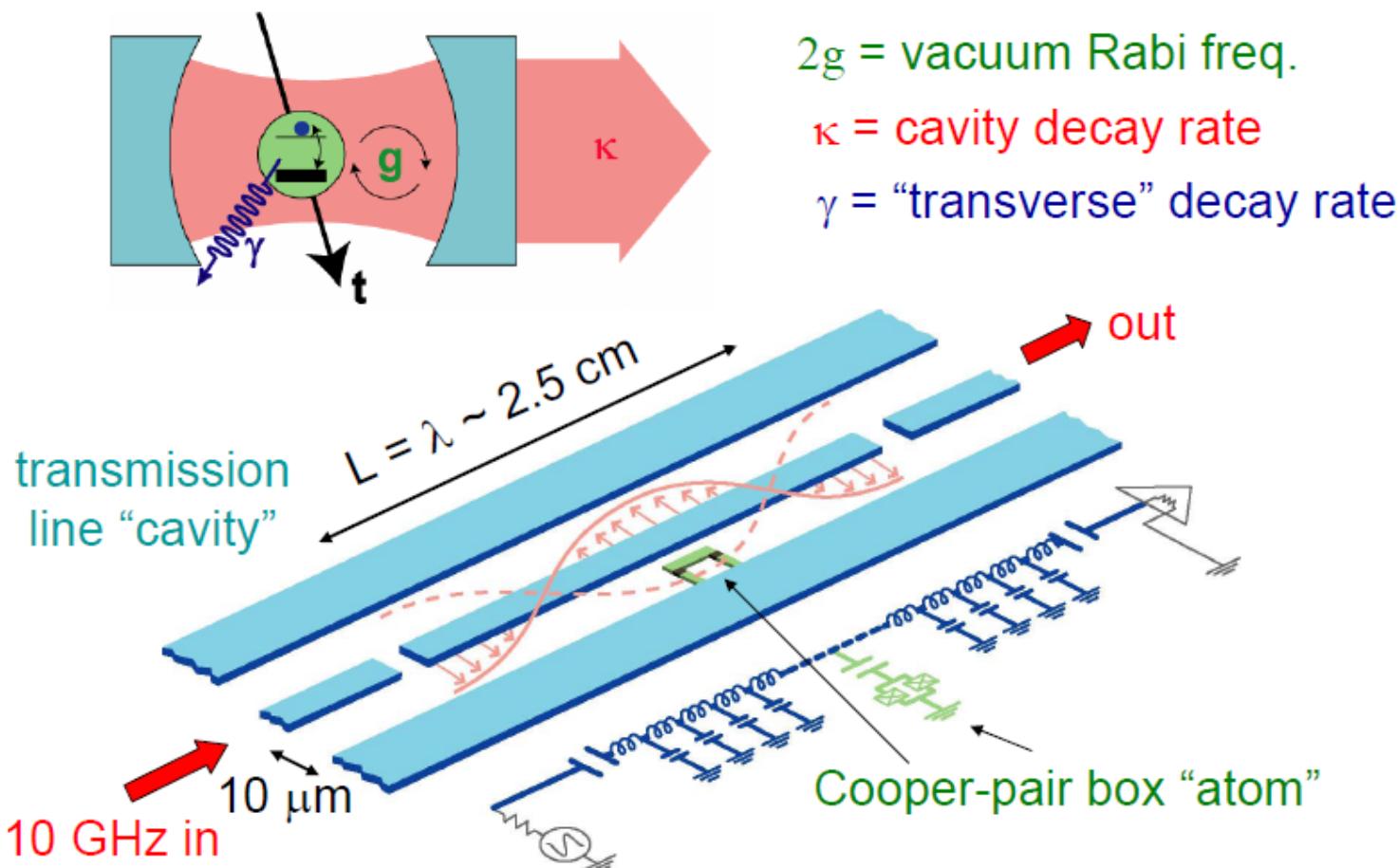
R. Barends et al., Phys. Rev. Lett. **111**, 080502 (2013)



X-mon qubit



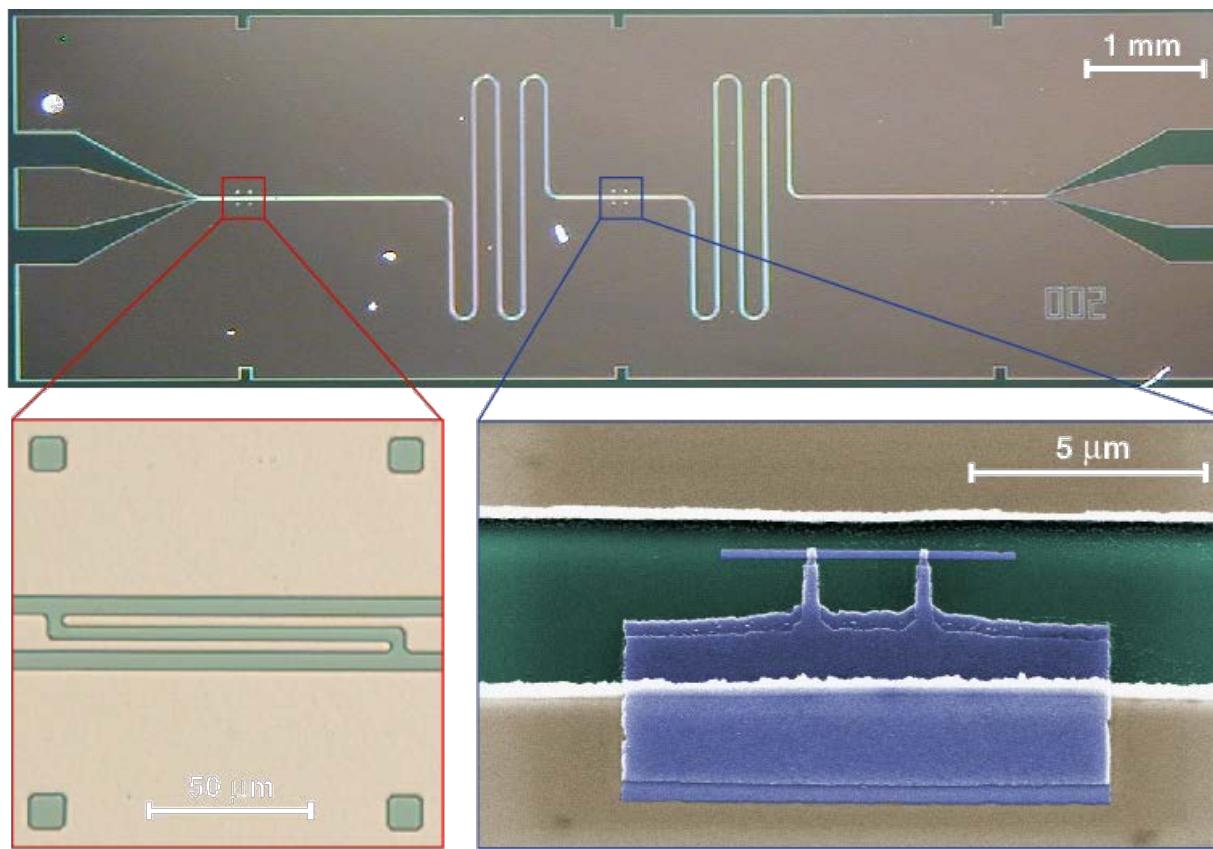
A circuit analog for cavity QED (Yale)



A. Wallraff, D. I. Schuster, A. Blais, et al., *Nature* 431, 162 (2004)

© A. Wallraff

Charge qubit in a cavity (Yale)



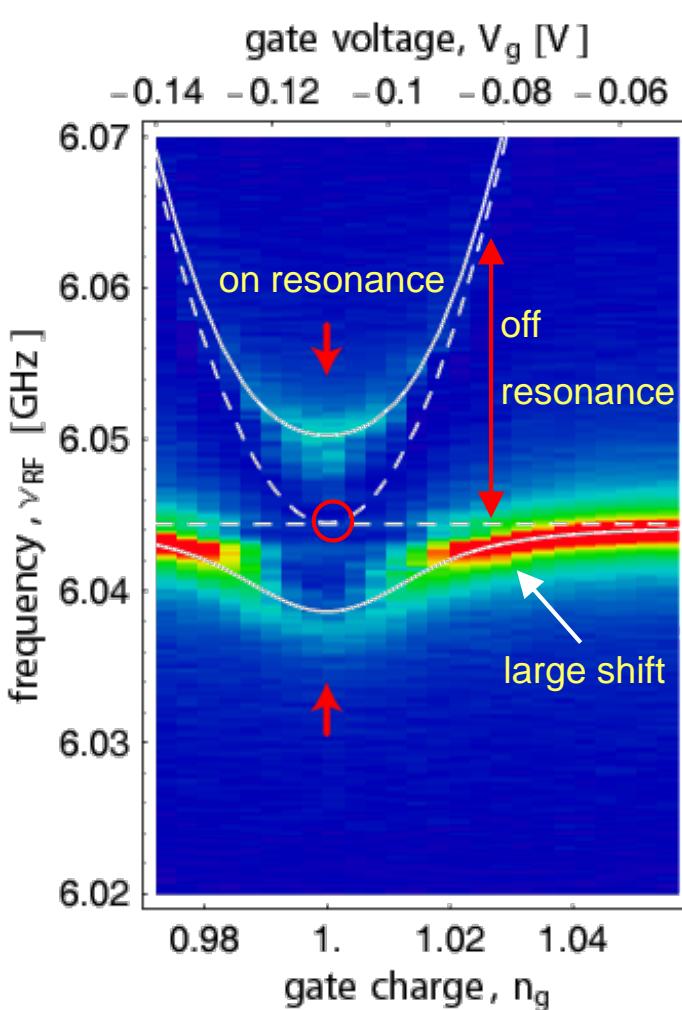
Andreas Wallraff



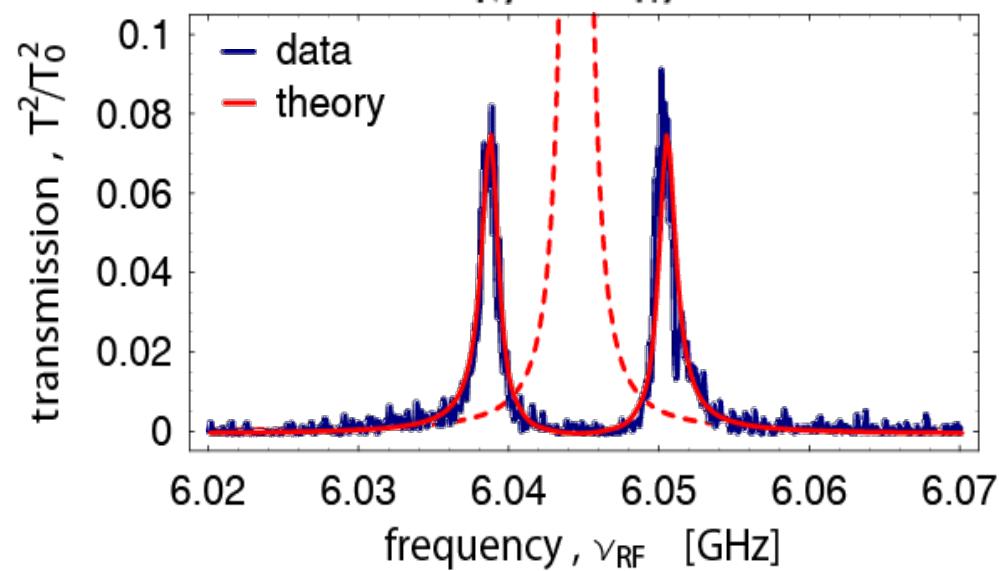
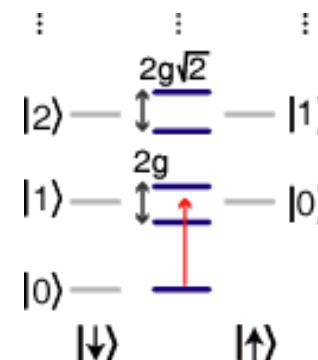
Robert Schoelkopf

A. Wallraff, D. I. Schuster, A. Blais, et al., *Nature* 431, 162 (2004)

Dispersive qubit-field interaction



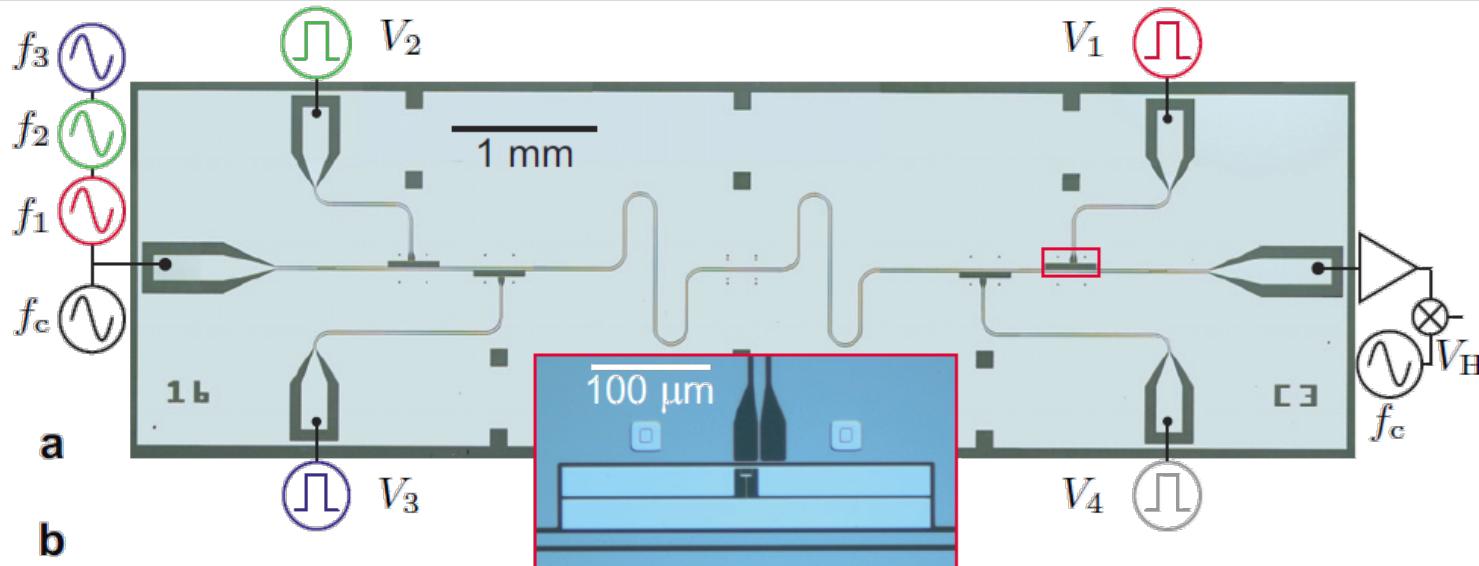
Photon-qubit
anti-crossing:
vacuum Rabi
splitting



A. Wallraff, D. I. Schuster, A. Blais, et al., *Nature* 431, 162 (2004)

Circuit QED with Transmon Qubits

L. DiCarlo et al., Nature 467, 574 (2010)



- each qubit interacts with the resonator. This is described by the **Jaynes-Cummings Hamiltonian**:

$$H = \hbar\omega_R \left(a^\dagger a + \frac{1}{2} \right) + \hbar\omega_Q \frac{\sigma_z}{2} + \hbar g (a^\dagger \sigma^- + a \sigma^+)$$

cavity (resonator) qubit coupling cavity-qubit

a^\dagger and a are photon creation / annihilation operators.

Flux qubit

quantum states:

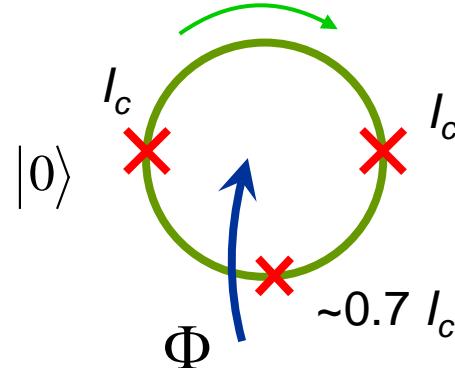
$$|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$$



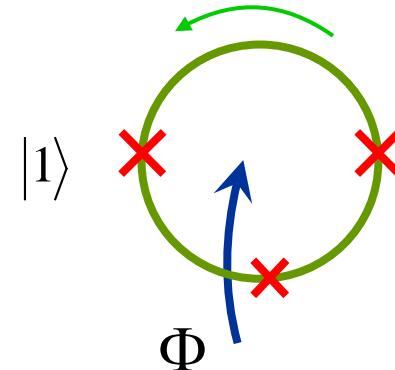
Hans
Mooij

J.E. Mooij *et al.*, *Science* 285, 1036 (1999)
C.H. van der Wal *et al.*, *Science* 290, 773 (2000)

clockwise current

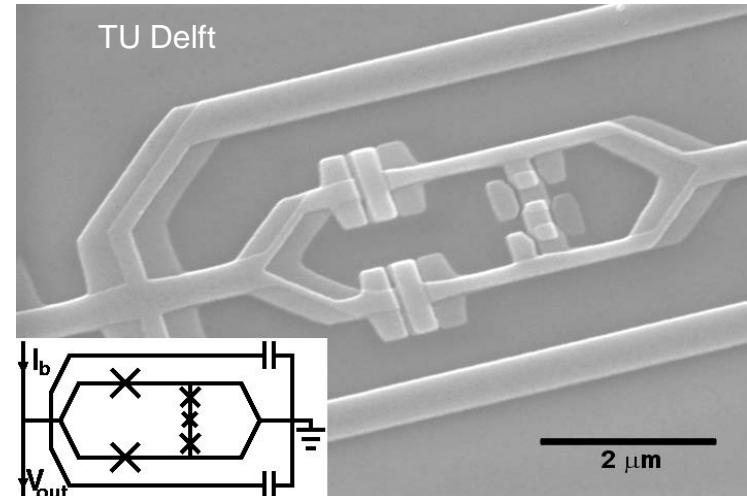


counterclockwise current

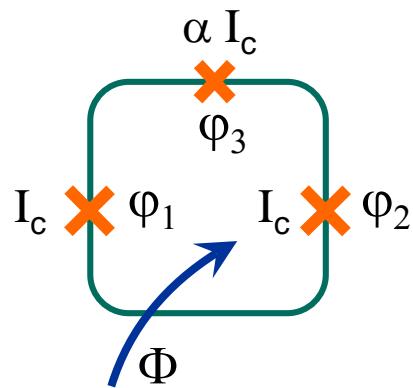


degeneracy point at

$$\Phi = \Phi_0/2$$



Superconducting 3-junction flux qubit



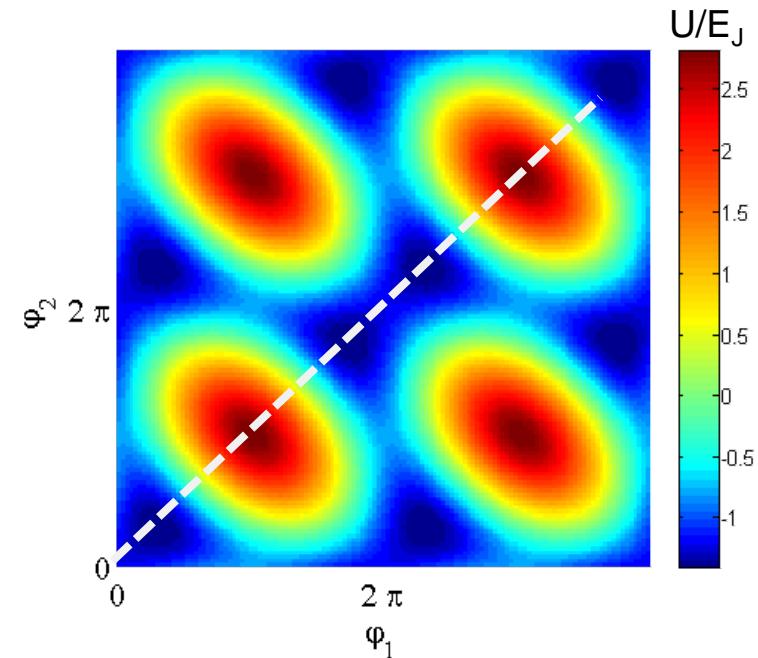
flux quantization:

$$\varphi_1 + \varphi_2 + \varphi_3 + 2\pi \frac{\Phi}{\Phi_0} = 2\pi n$$

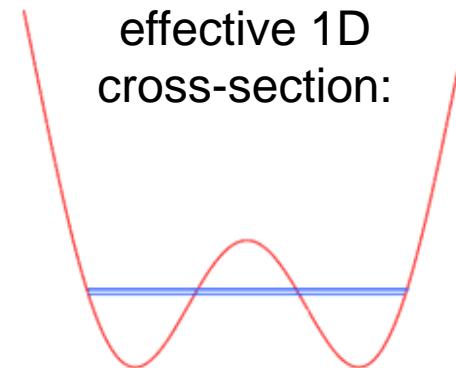
effective 2D potential:

$$\frac{U}{E_J} = \cos \varphi_1 + \cos \varphi_2 + \alpha \cos \left(-\varphi_1 - \varphi_2 - 2\pi \frac{\Phi_{\text{ext}}}{\Phi_0} \right)$$

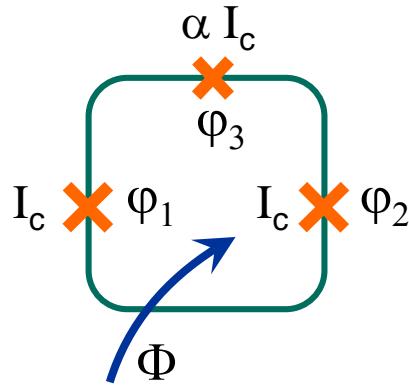
Mooij et al. Science 285, 1036 (1999)
Van der Wal et al. Science 290, 1140 (2000)



effective 1D cross-section:



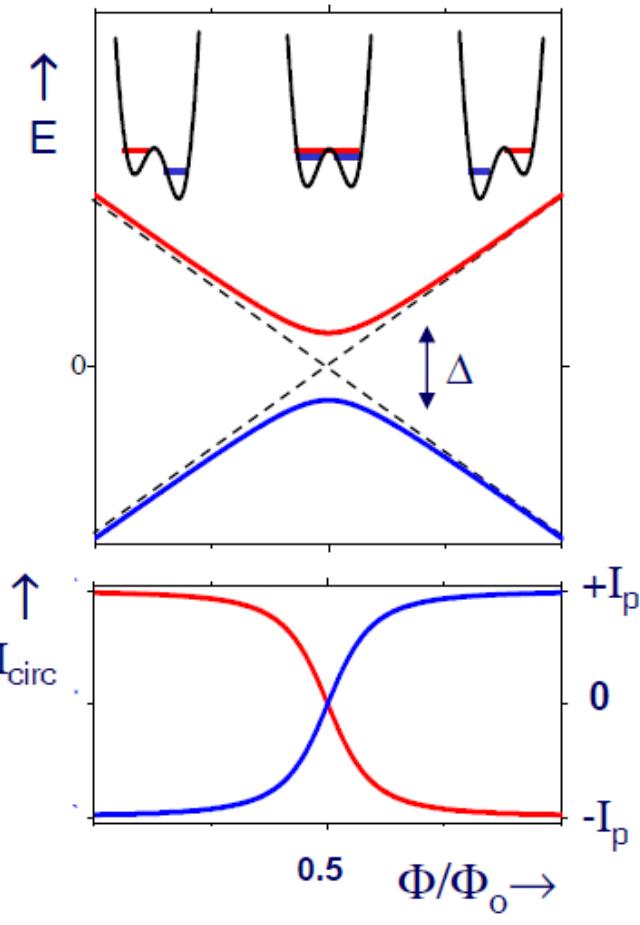
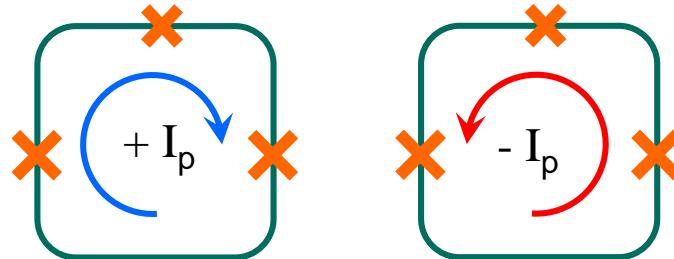
Superconducting flux qubit as a two-level system (artificial atom)



magnetic flux bias $\Phi \sim \Phi_0/2$

$$H = \frac{1}{2} (\varepsilon \sigma_z + \Delta \sigma_x)$$

persistent current states $\pm I_p$

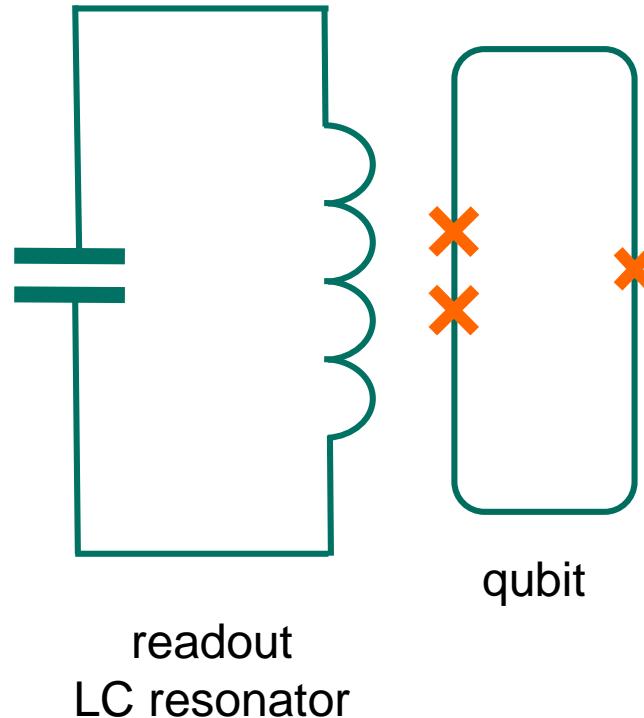
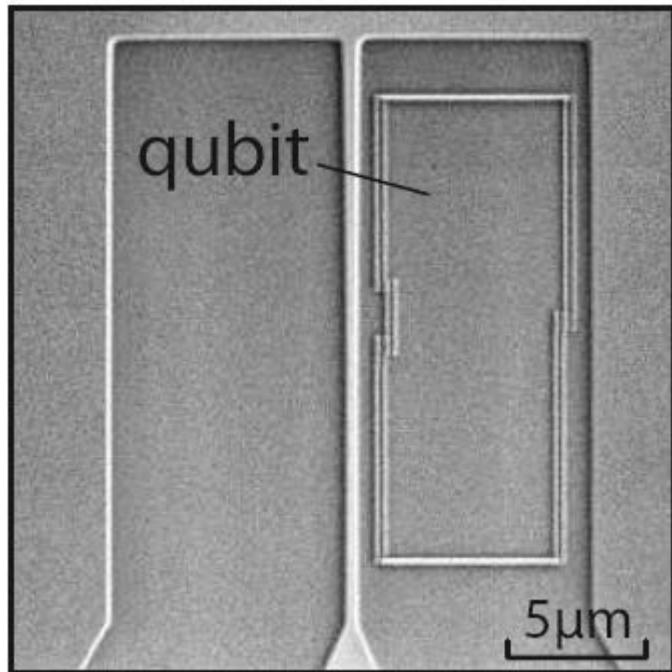


J. Clarke and F. K. Wilhelm, *Nature* 453, 1031 (2008)

— ground state
— excited state

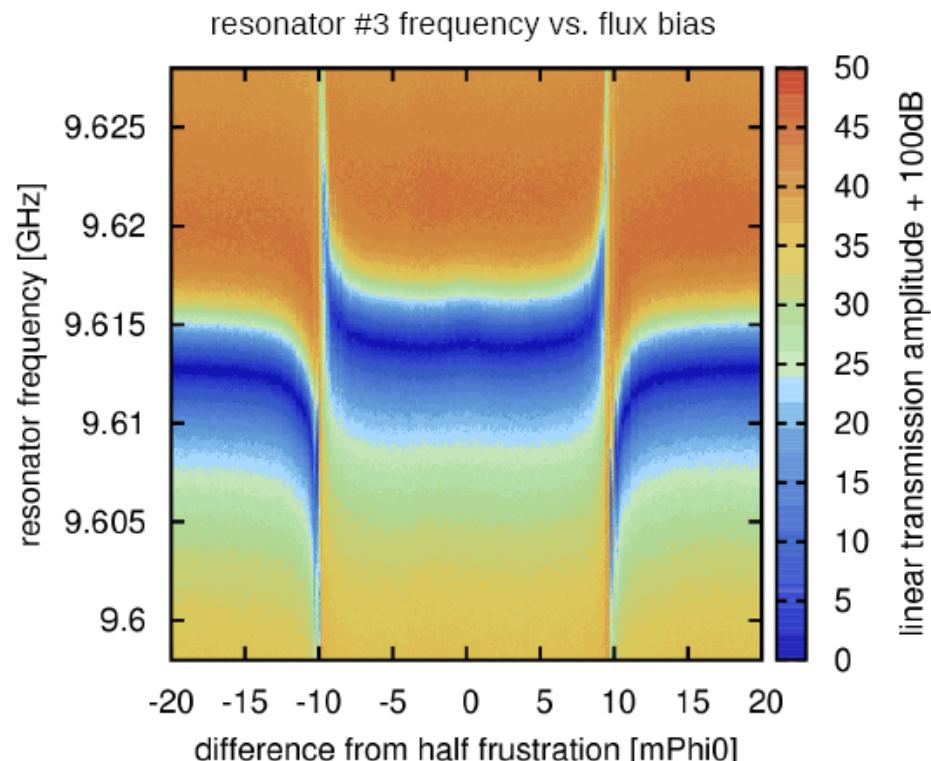
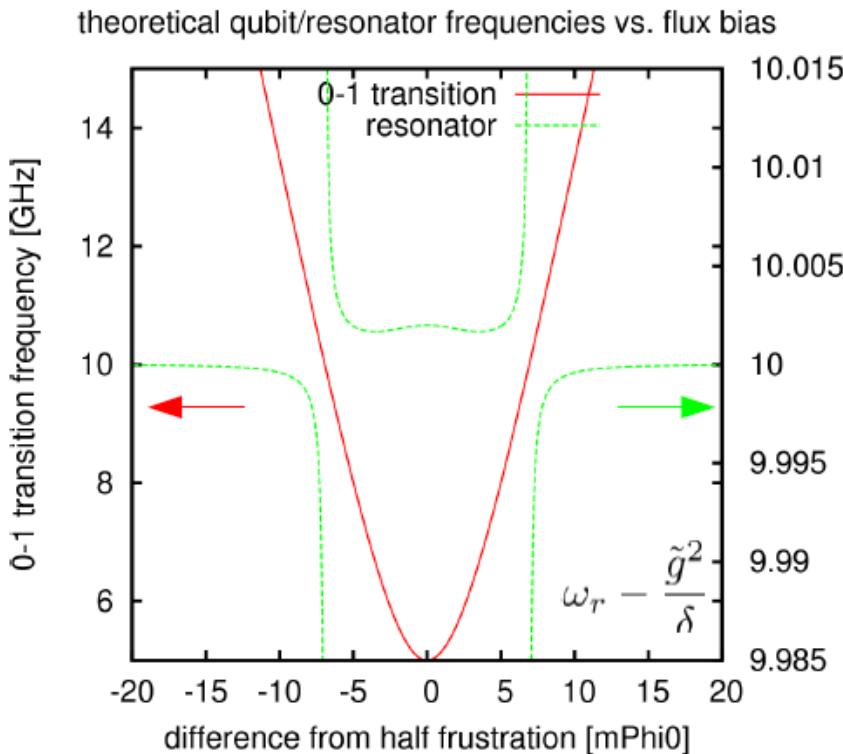
QED readout of a flux qubit

Al/AIO_X/Al Josephson junctions



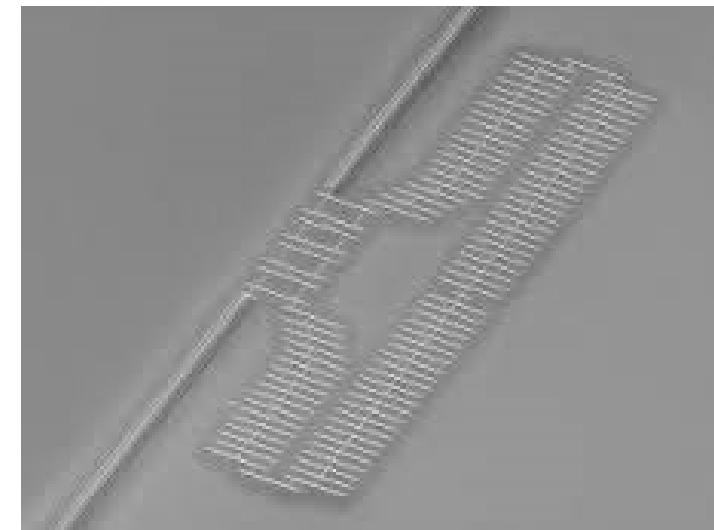
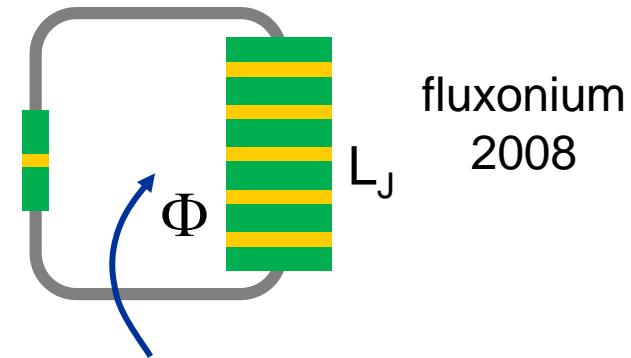
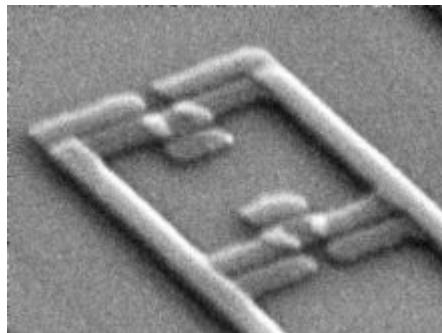
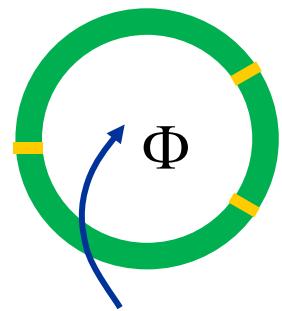
Anticrossings

- Dispersive shift of the resonator due to the qubit in the ground state
- Evaluate qubit-resonator coupling, if gap Δ is known



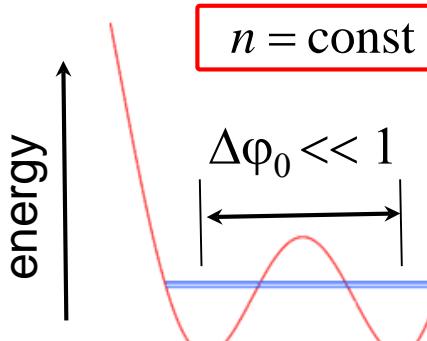
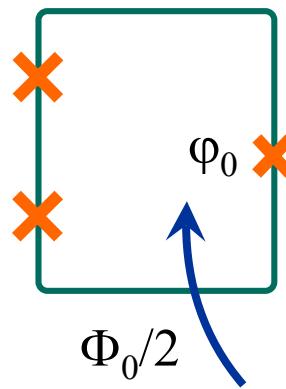
Flux qubit vs „fluxonium“

flux qubit
2000



Flux qubit vs „fluxonium“

flux qubit

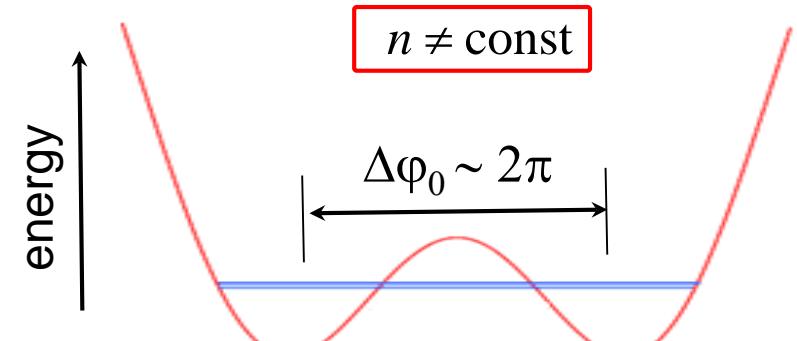
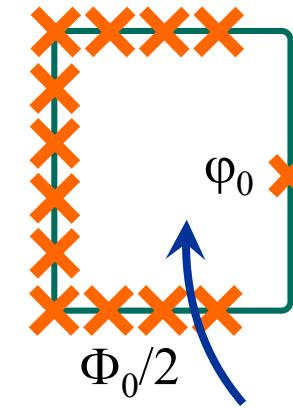


tunneling
between two
states
separated by
phase difference
 $\Delta\phi_0$

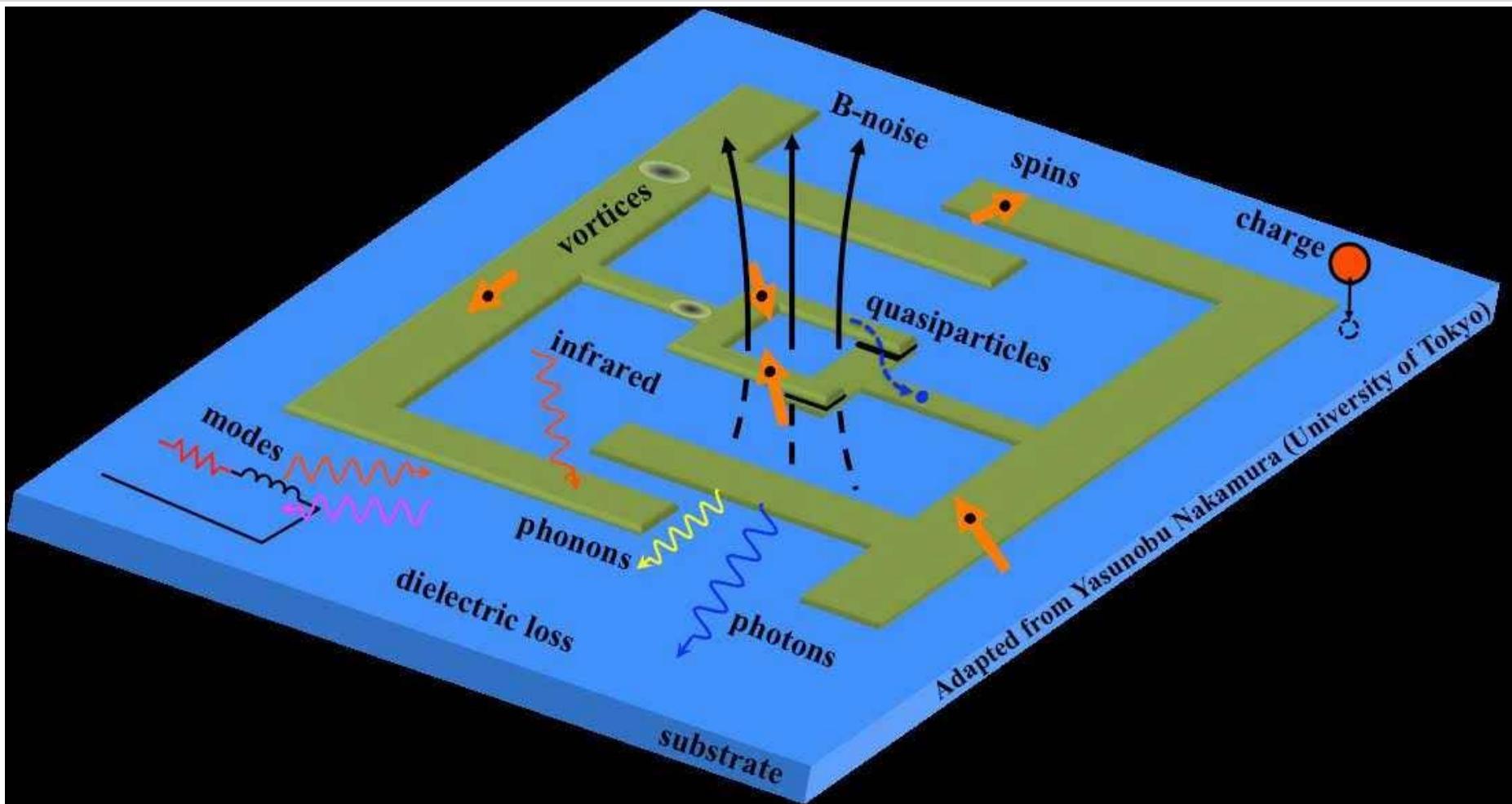
flux quantization:

$$\sum \varphi_i + 2\pi \frac{\Phi}{\Phi_0} = 2\pi n$$

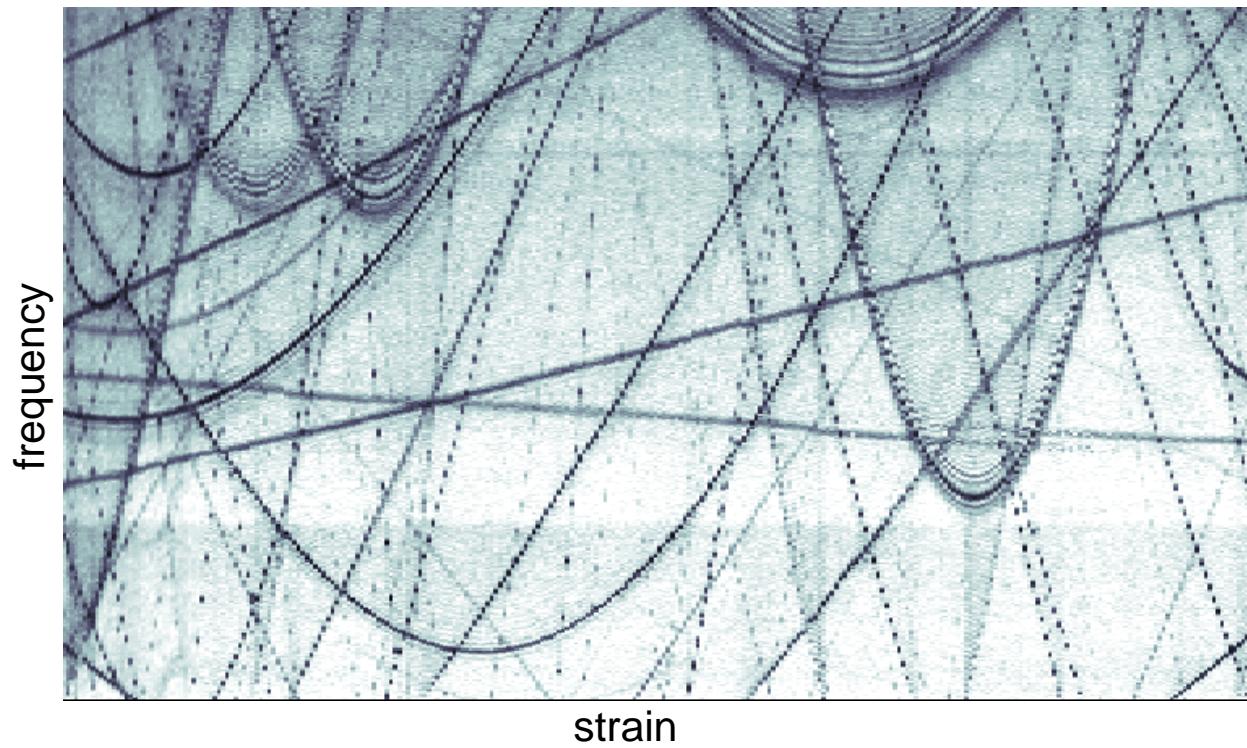
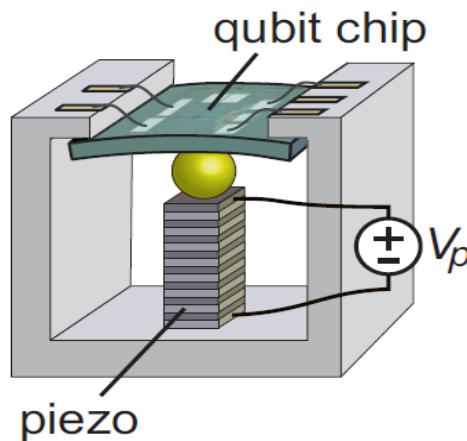
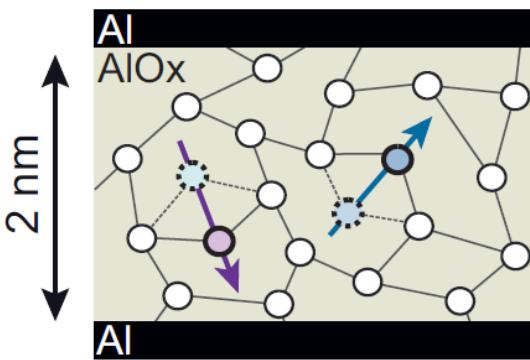
fluxonium



Decoherence



Microscopic defects in qubits

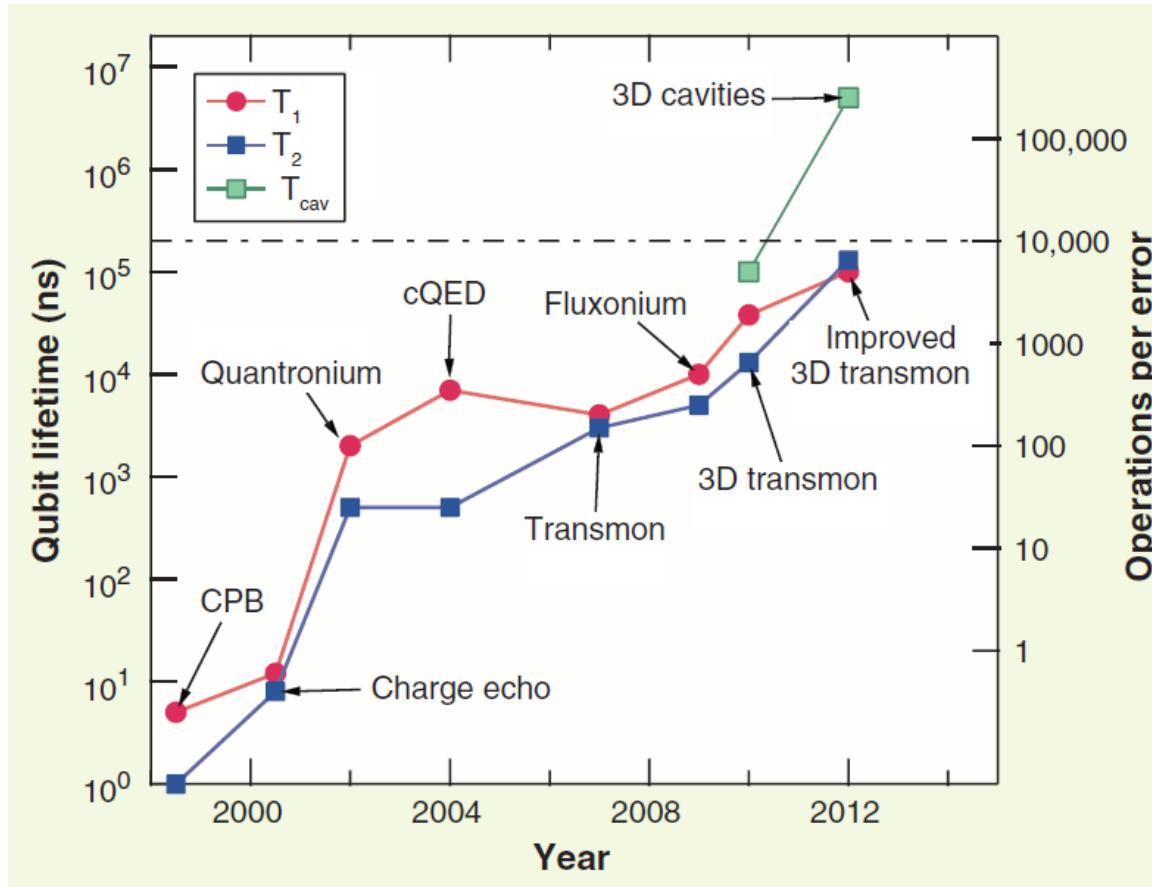


applying mechanical strain tunes the energy splitting of defects in qubits

G. Grabovskij, T. Pecl, J. Lisenfeld, G. Weiss, and A. V. Ustinov, *Science* **338**, 232 (2012)

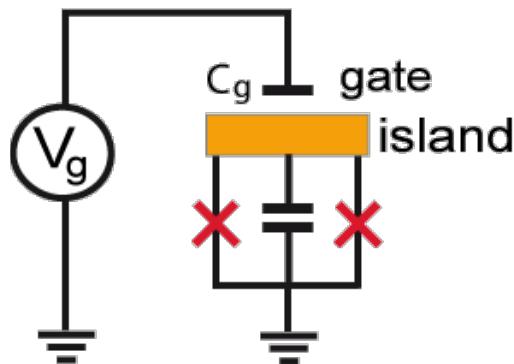
Progress of superconducting qubits

“Moore’s law” of superconducting qubits

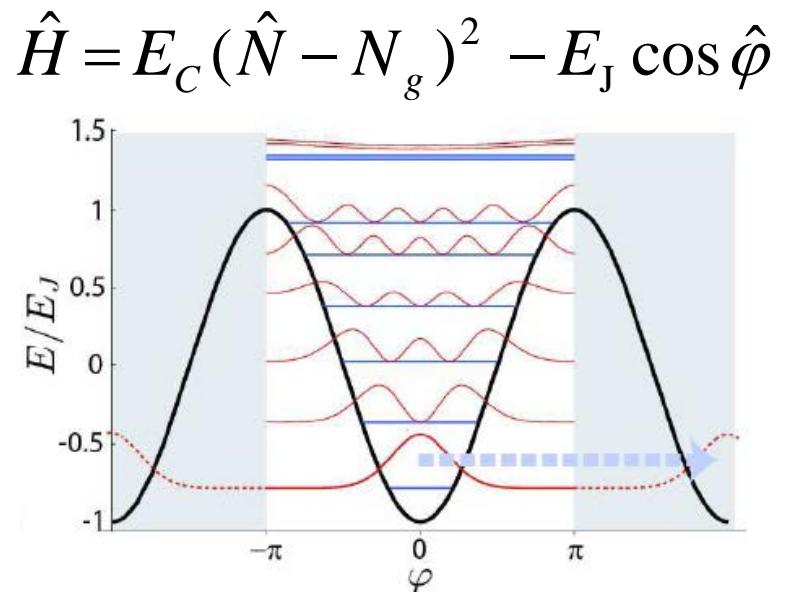
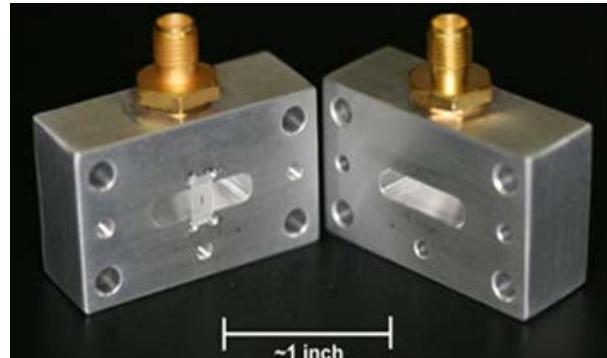
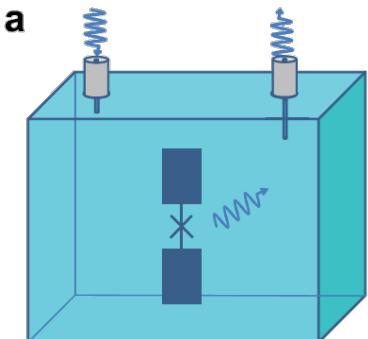


M. H. Devoret and R. J. Schoelkopf, *Science* **339**, 1169 (2013)

The 3D – Transmon Qubit



- no inductance → minimized flux noise
- large capacitance → avoids charge noise
- low field intensities → no excitation of defects
- avoids dielectrics → low energy dissipation



$$T_1 = 70 \text{ } \mu\text{s}$$
$$T_2^* = 95 \text{ } \mu\text{s}$$

Pictures: R. Schoelkopf, Yale / IBM
Ch. Rigetti et al., arXiv: 1202.5533 (2012)

H. Paik et al., PRL 107, 240501 (2011)

Transmon basics:

J. Koch et al., Phys. Rev. A **76**, 042319 (2007)

New development: gatemons

PRL 116, 150505 (2016)

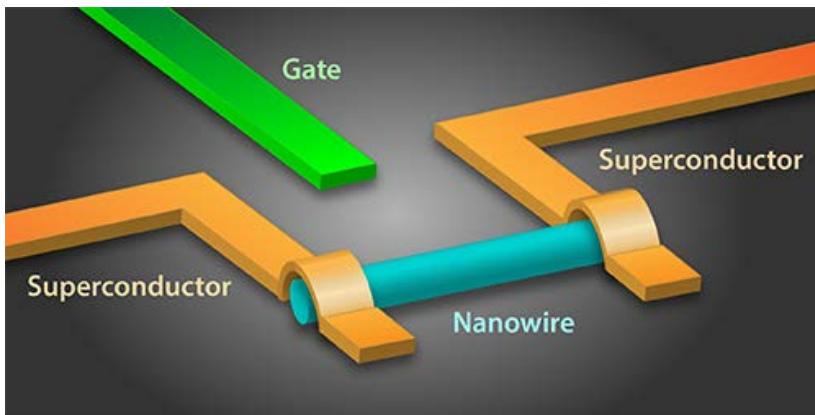
PHYSICAL REVIEW LETTERS

week ending
15 APRIL 2016

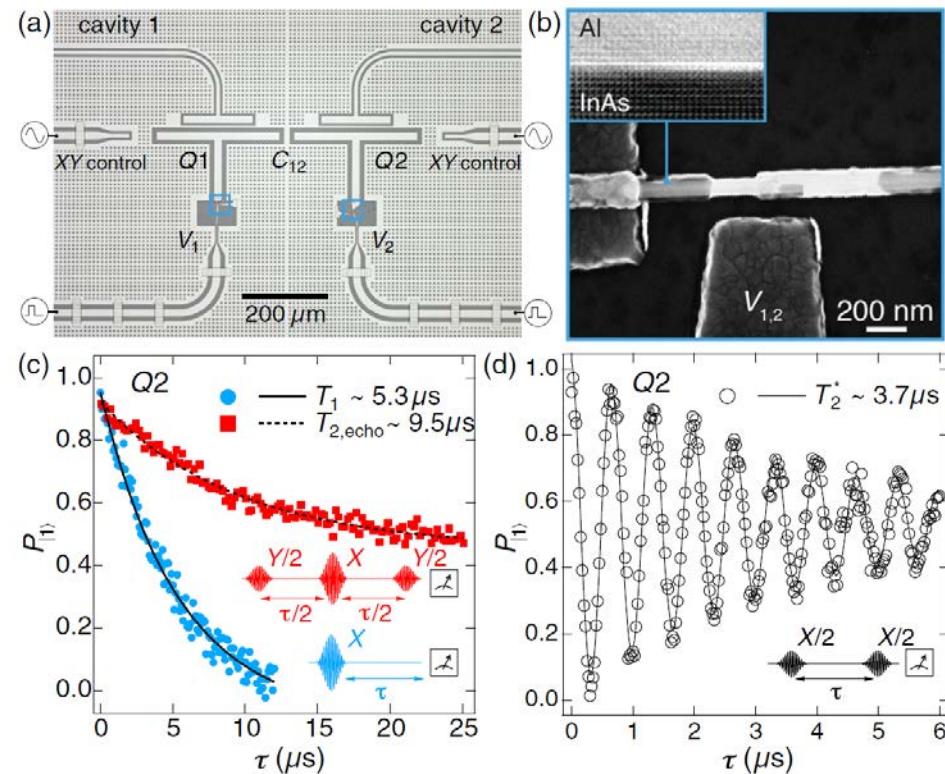
Gatemon Benchmarking and Two-Qubit Operations

L. Casparis,¹ T. W. Larsen,¹ M. S. Olsen,¹ F. Kuemmeth,¹ P. Krogstrup,¹ J. Nygård,^{1,2}
K. D. Petersson,¹ and C. M. Marcus¹

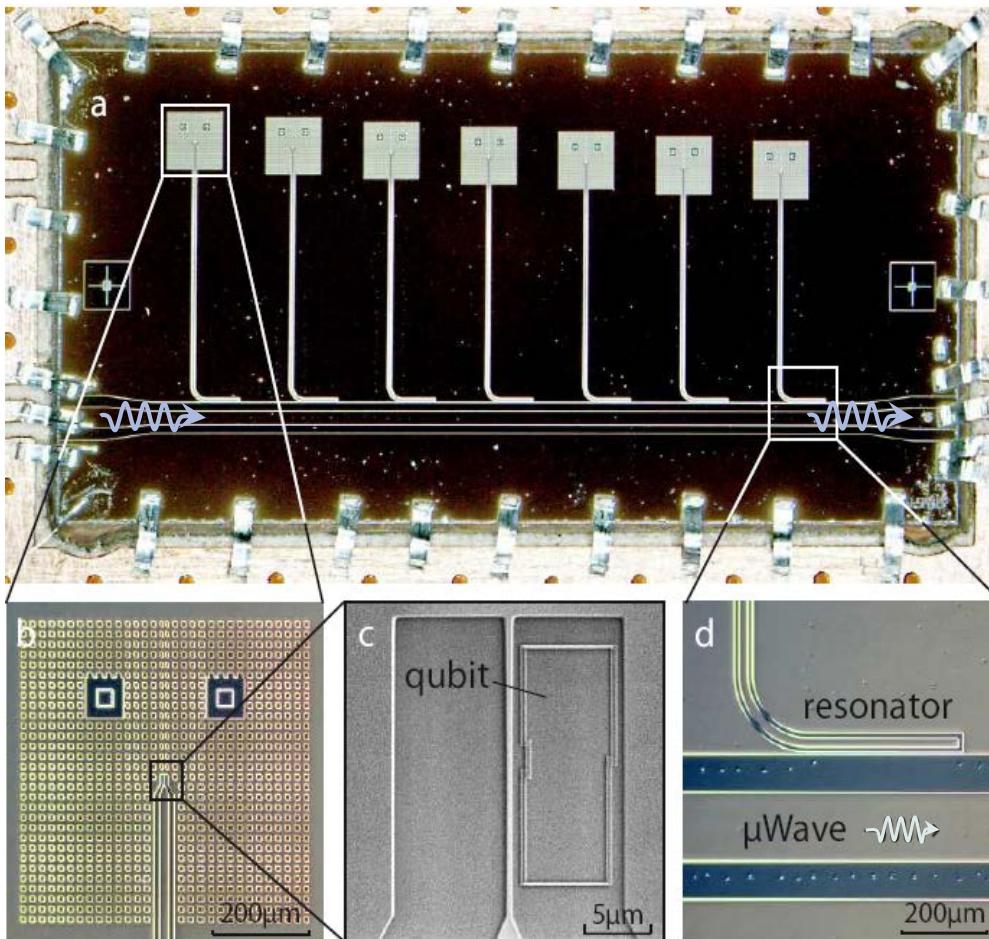
- „gatemon“ – transmon-like qubit tunable by dc voltage gate



- single-qubit fidelity 99.3%
- two-qubit fidelity 91%

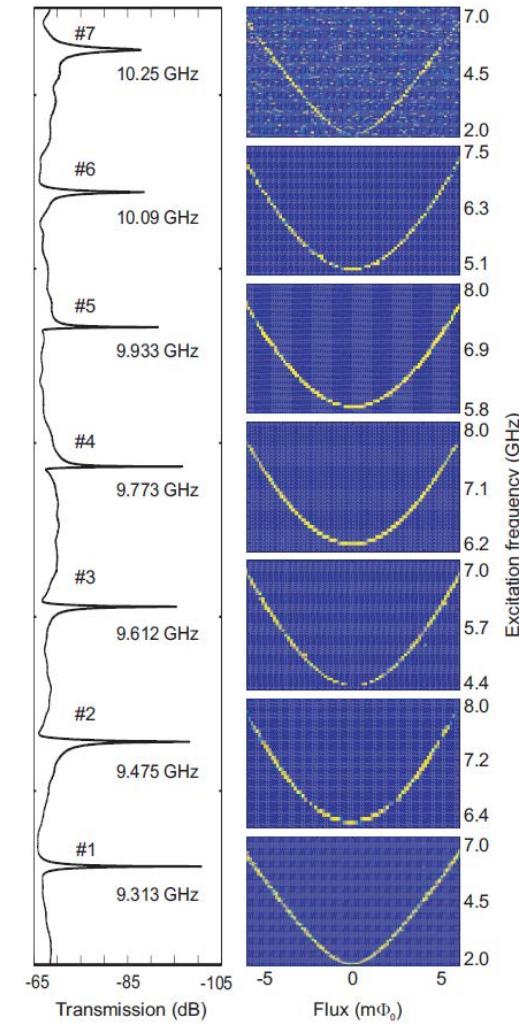


Multiplexed Readout of Superconducting Qubits



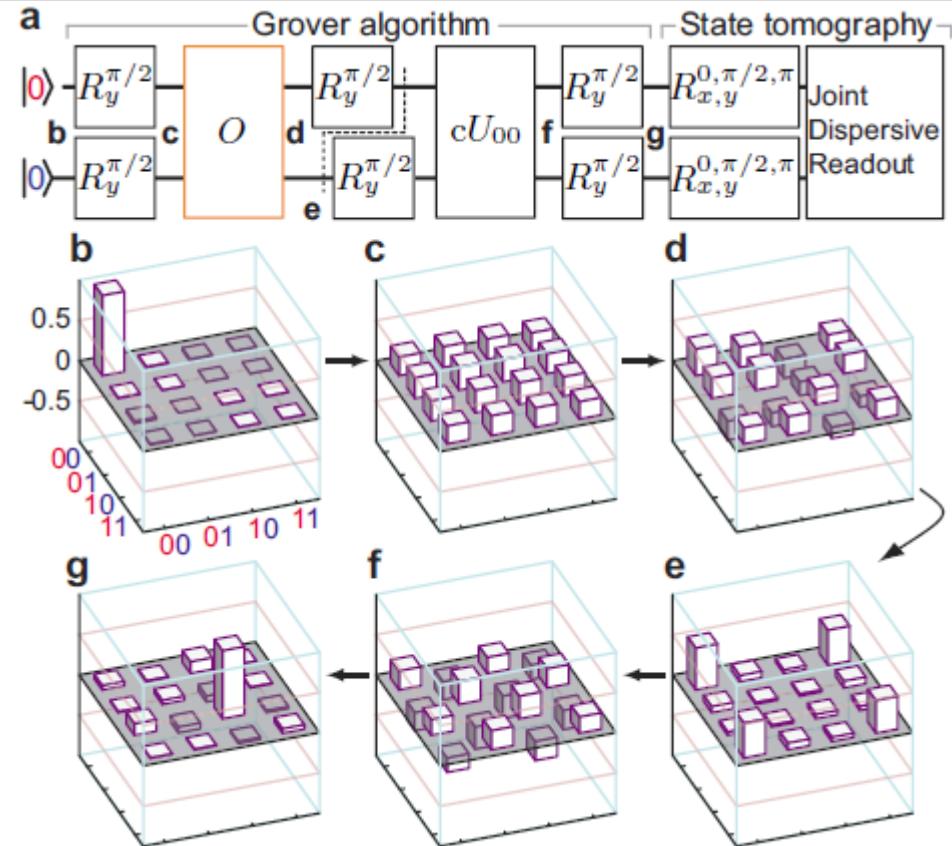
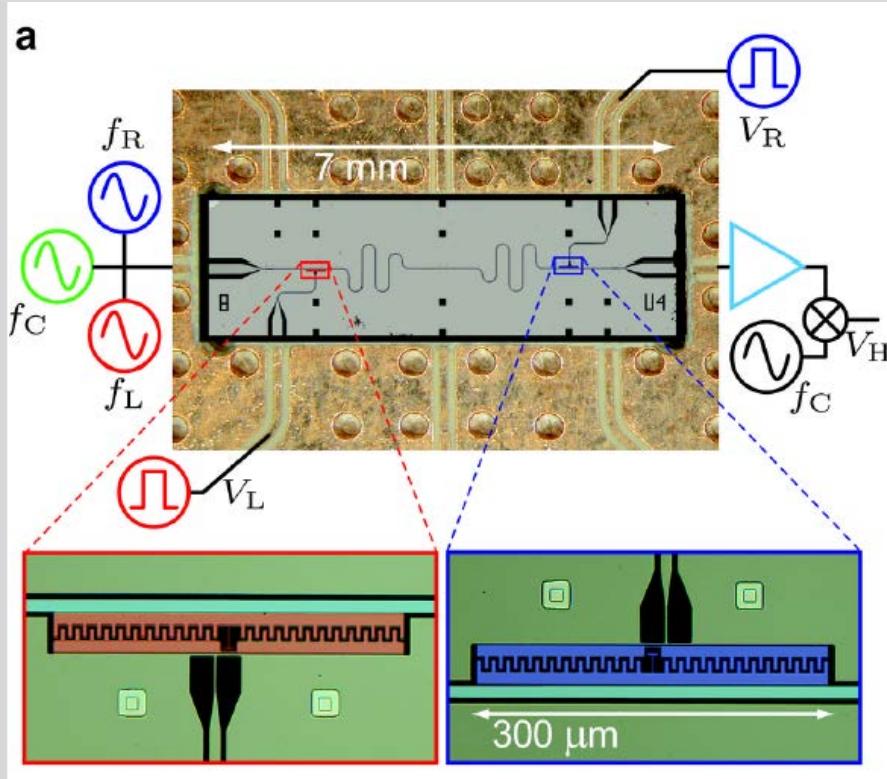
M. Jerger et al., *Europhys. Lett.* **96**, 40012 (2011)

M. Jerger et al., *Appl. Phys. Lett.* **101**, 042604 (2012)



Demonstration of 2-Qubit Algorithms

L. DiCarlo et al., Nature 460, 240 (2009)

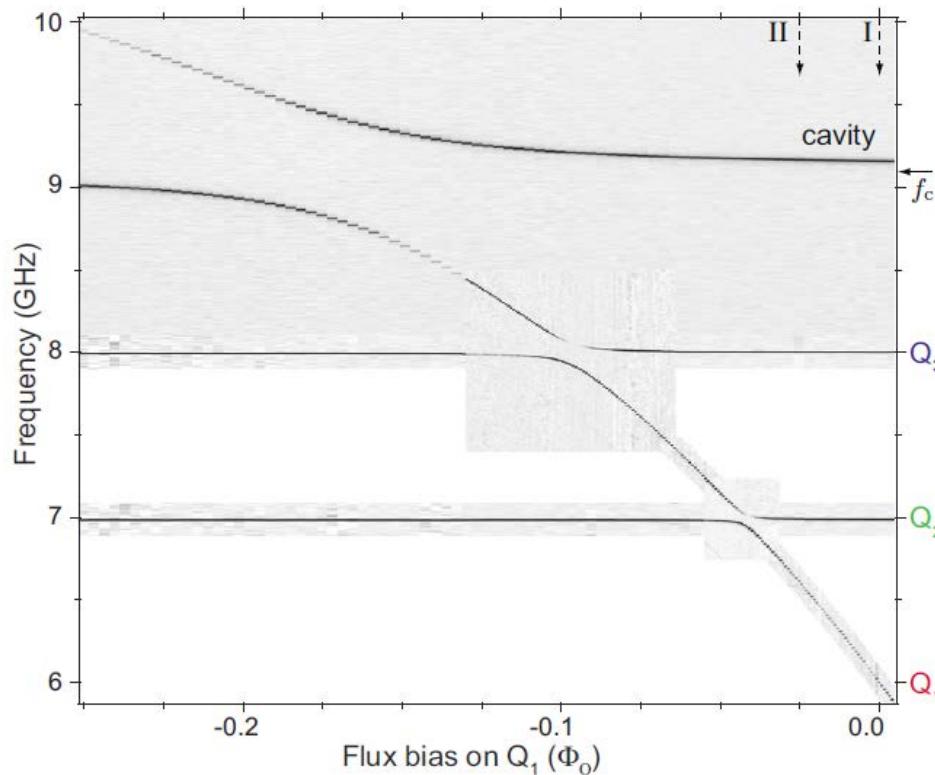


- the Grover algorithm for searching an unsorted database is demonstrated with a fidelity of 85 %.
- (b) starting state (00)
- (c) equal superposition of all 4 states
- (d) rotation of the phase of the searched state (10)
- (g) maximal amplitude of the searched state.

3 coupled Transmon Qubits

L. DiCarlo et al., Nature 467, 574 (2010)

- Spectroscopy:

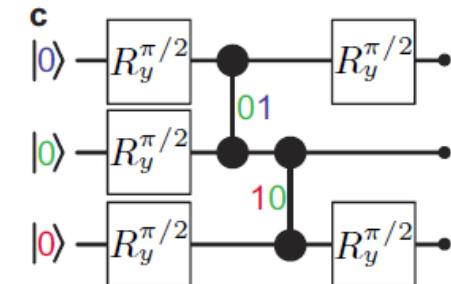


fidelity of GHZ preparation : 88%

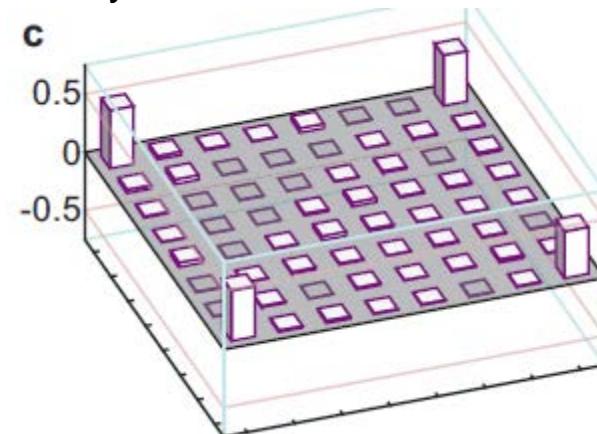
- demonstration of **GHZ** (Greenberger, Horne, and Zeilinger) state of 3 entangled qubits:

$$\Psi_{\text{GHZ}} = \frac{1}{\sqrt{2}} (\lvert 000 \rangle + \lvert 111 \rangle)$$

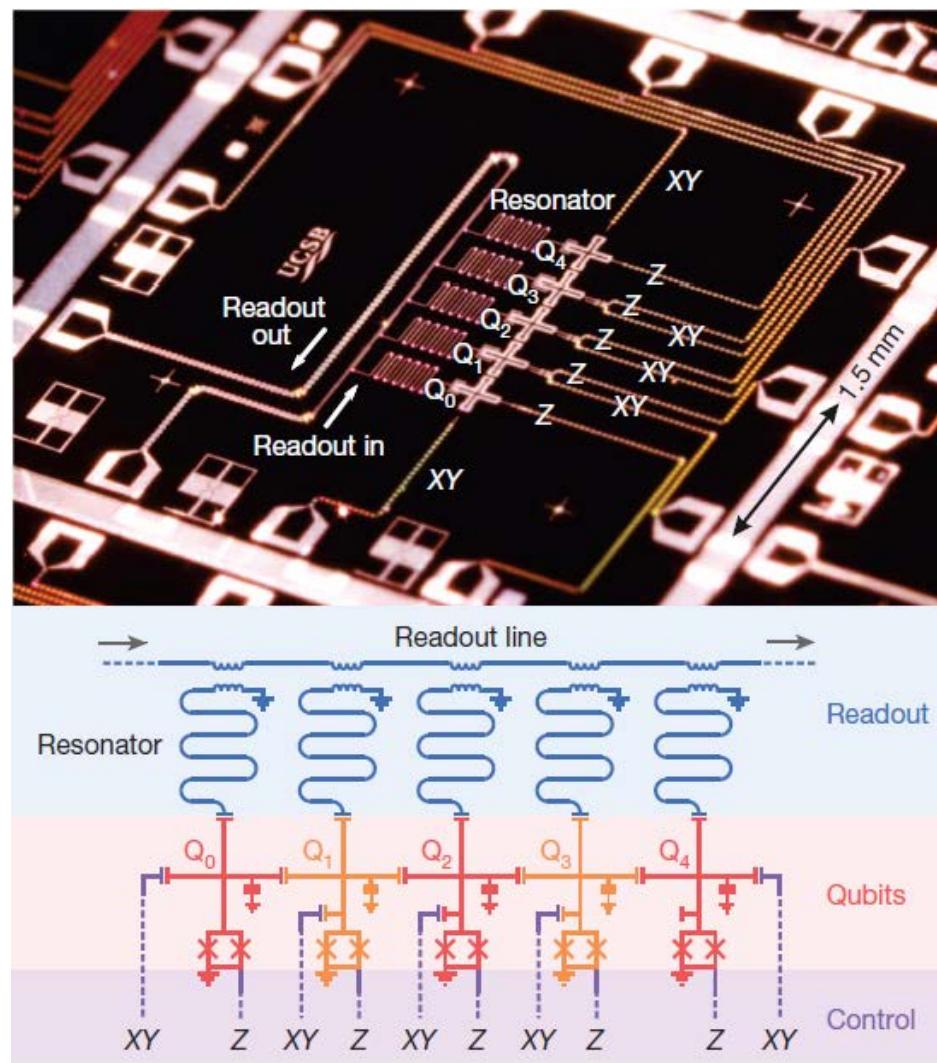
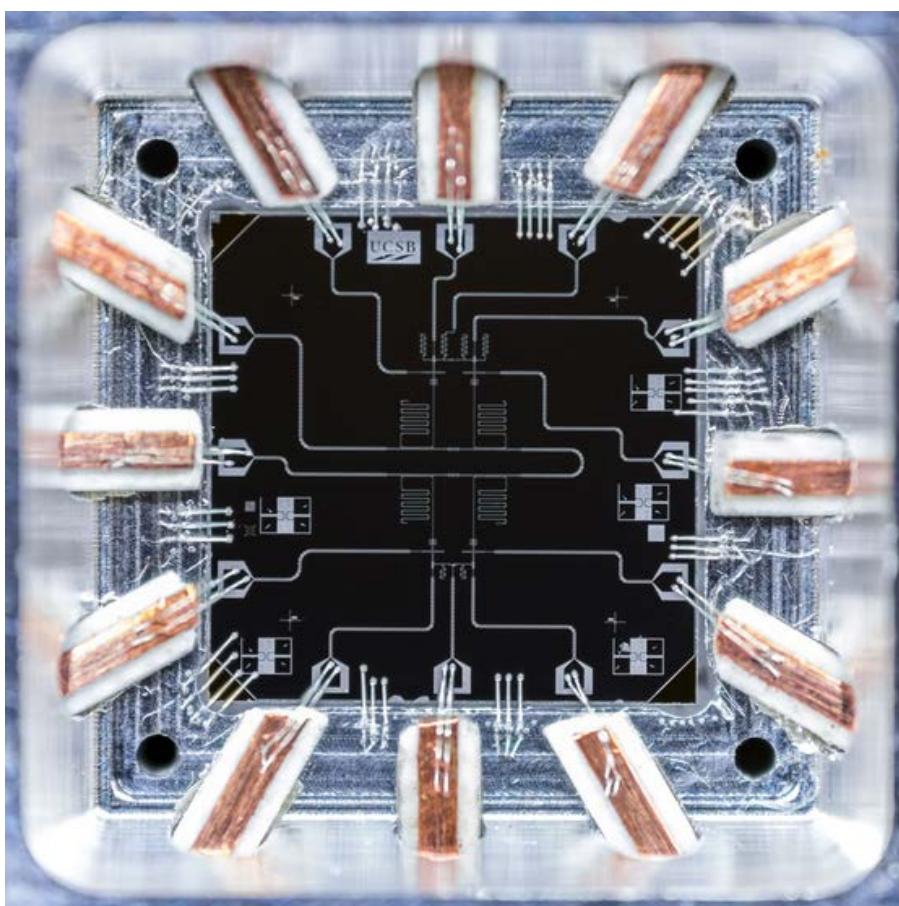
algorithm for preparation of GHZ:



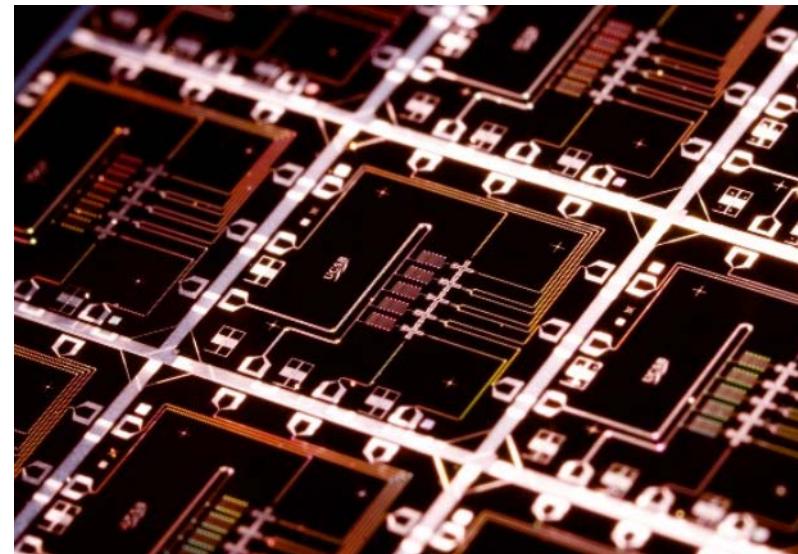
result: density matrix



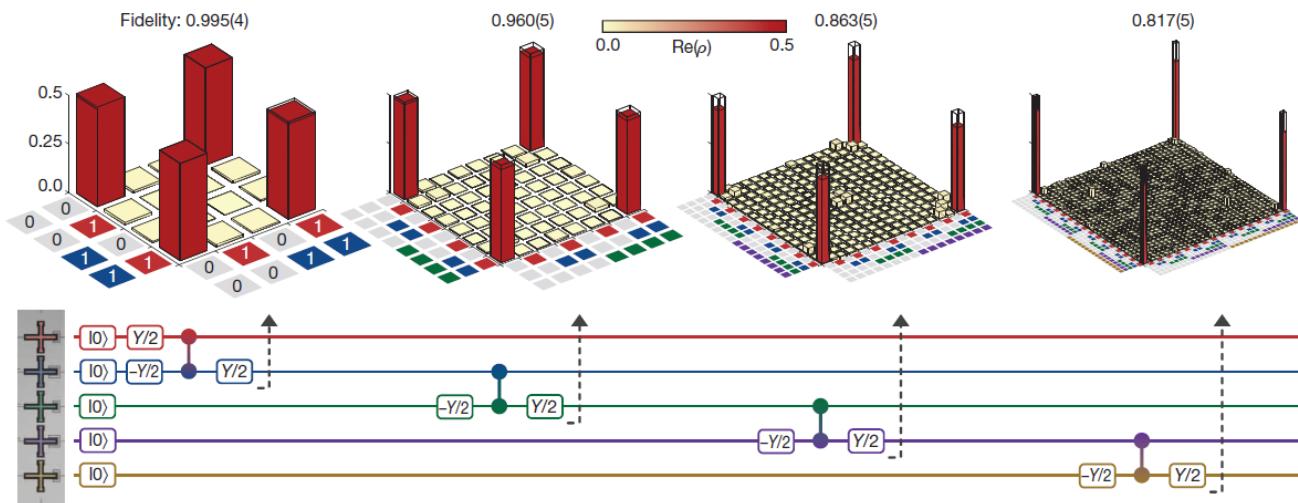
2014: Reaching the surface code threshold for fault tolerance of > 99.4% fidelity



John Martinis' UCSB team hired by Google

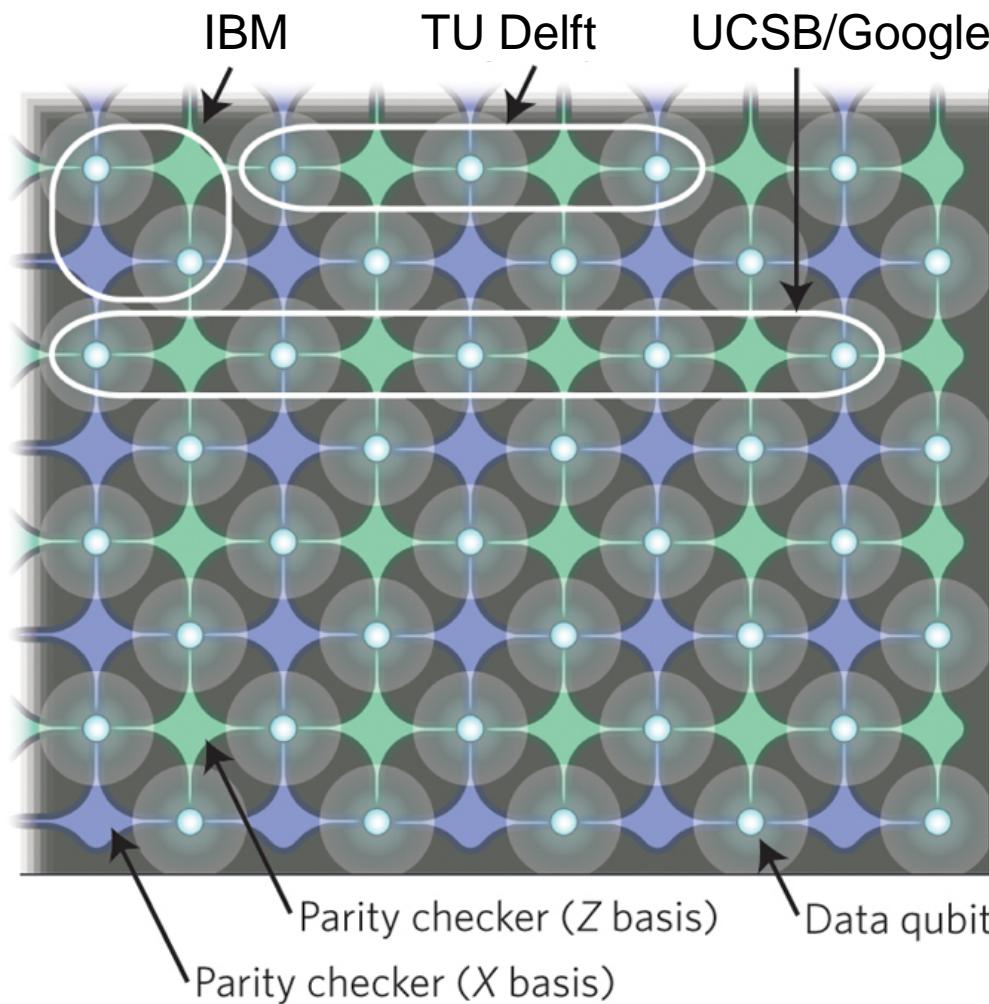


September 2014

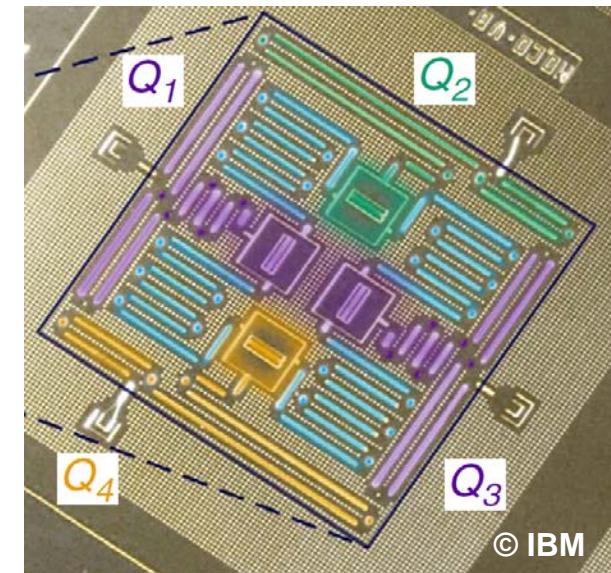


Alexey Ustinov

Kitaev's surface code implementation

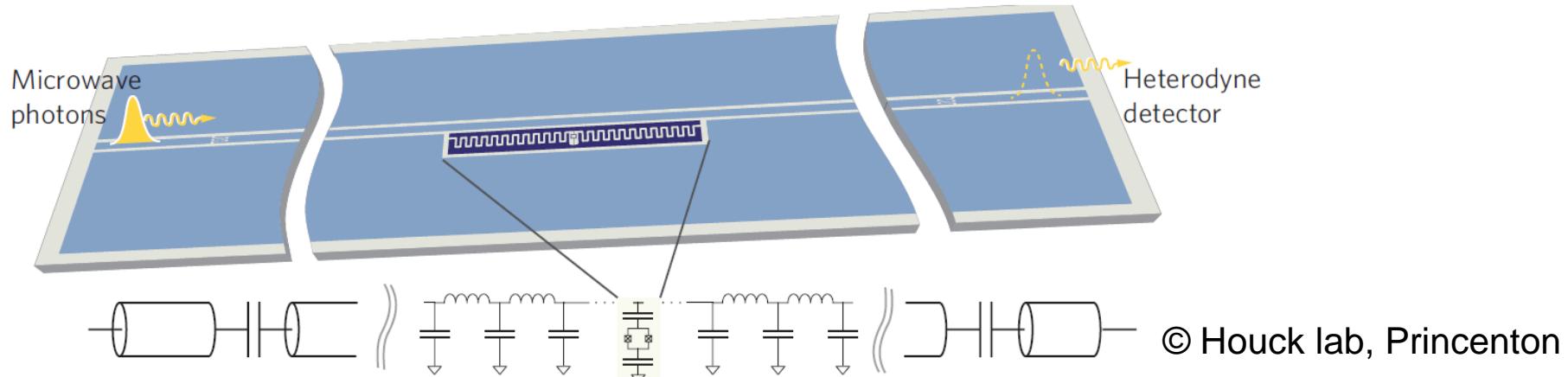


In 2015 three teams have demonstrated the basic parts of the Kitaev's surface code

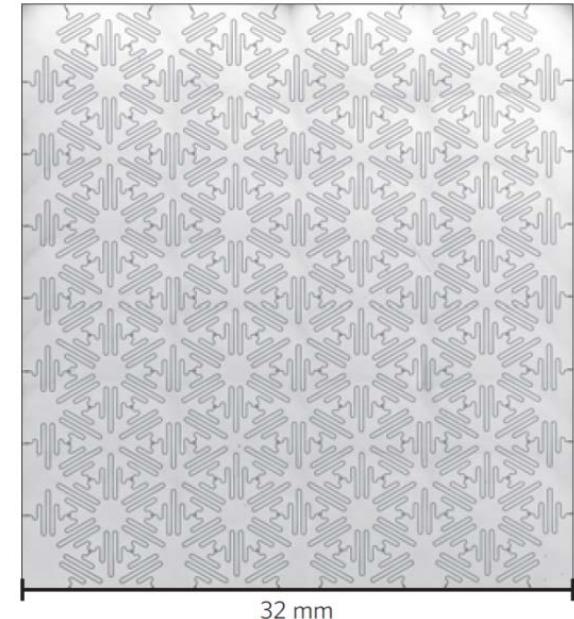


S. Benjamin and J. Kelly,
Nature Materials **14**, 561 (2015)

Quantum simulators

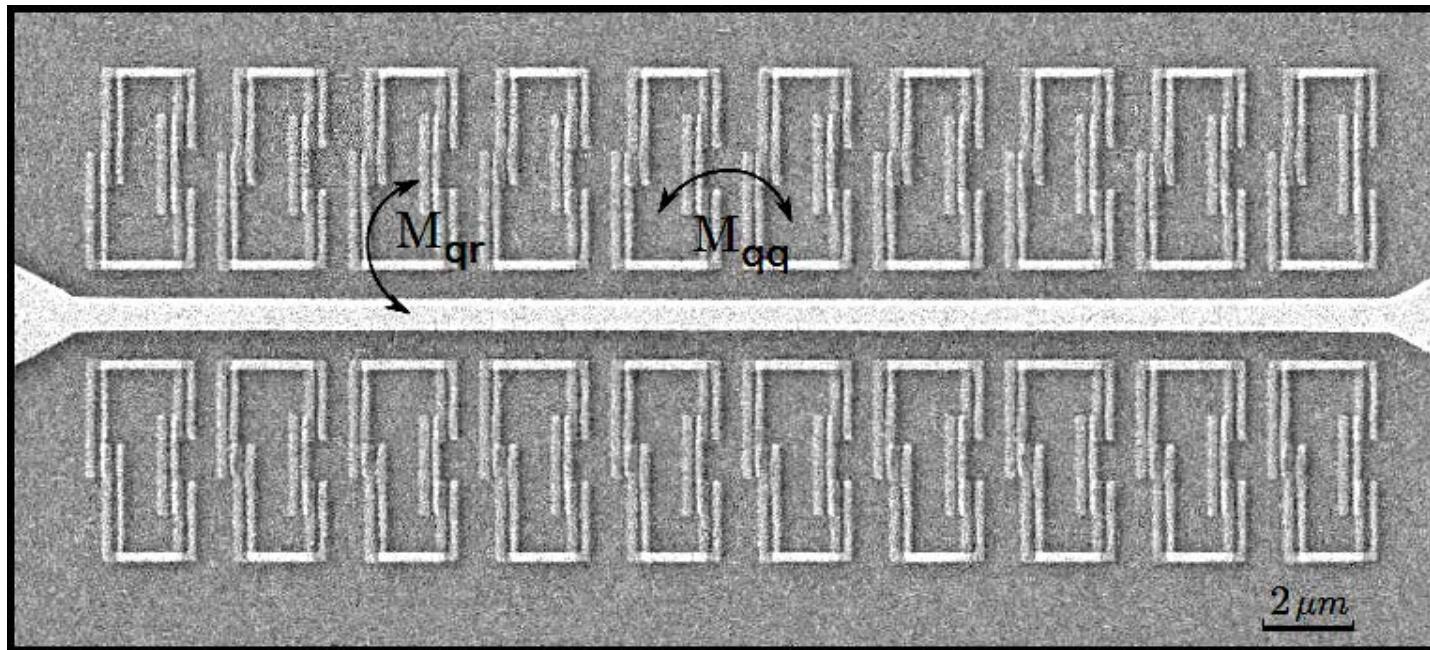


Main idea: Measure the result of interaction of a superconducting qubit with an environment composed of an array of spectrally tailored microwave resonators



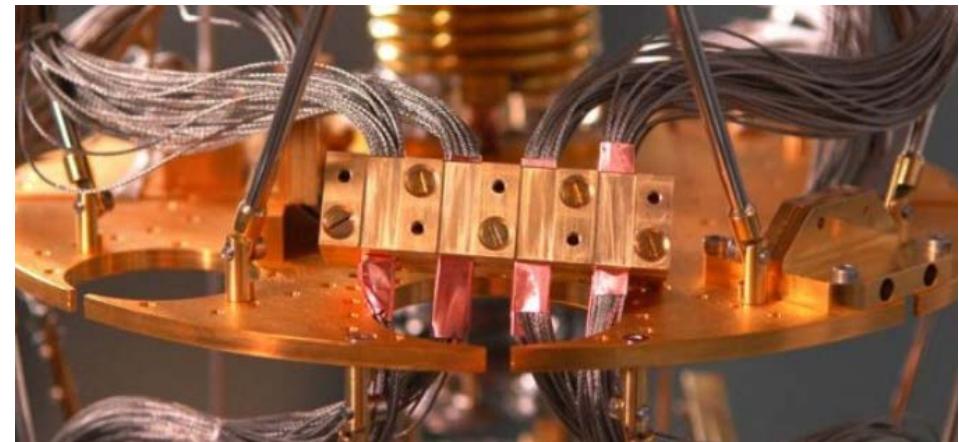
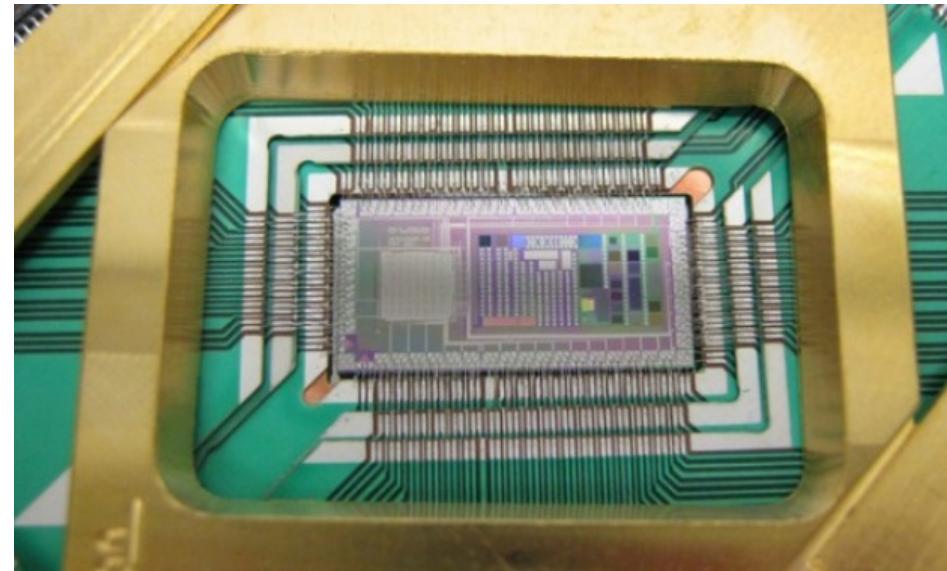
Superconducting quantum metamaterial: the simplest quantum simulator

20 flux qubits

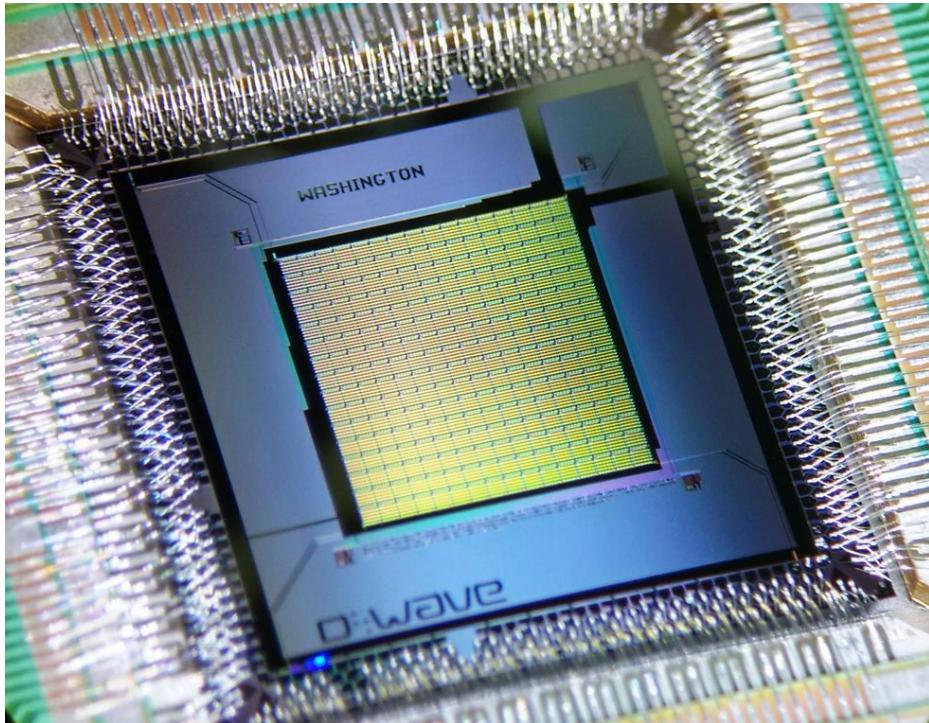


P. Macha, G. Oelsner, J.-M. Reiner, M. Marthaler, S. André, G. Schön, U. Huebner, H.-G. Meyer, E. Il'ichev, and A. V. Ustinov, *Nature Commun.* **5**, 5146 (2014)

D-Wave quantum computer: adiabatic quantum annealer-simulator



D-Wave Quantum Computer controversy



D-WAVE

The Quantum Computing Company™

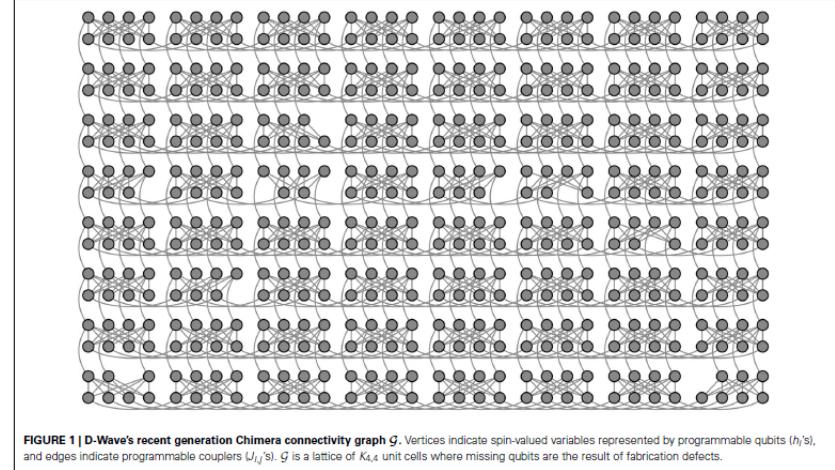


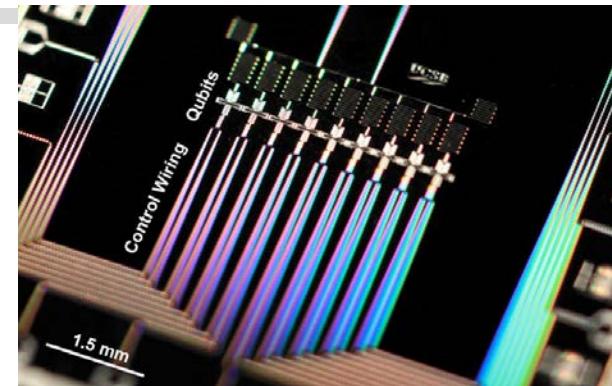
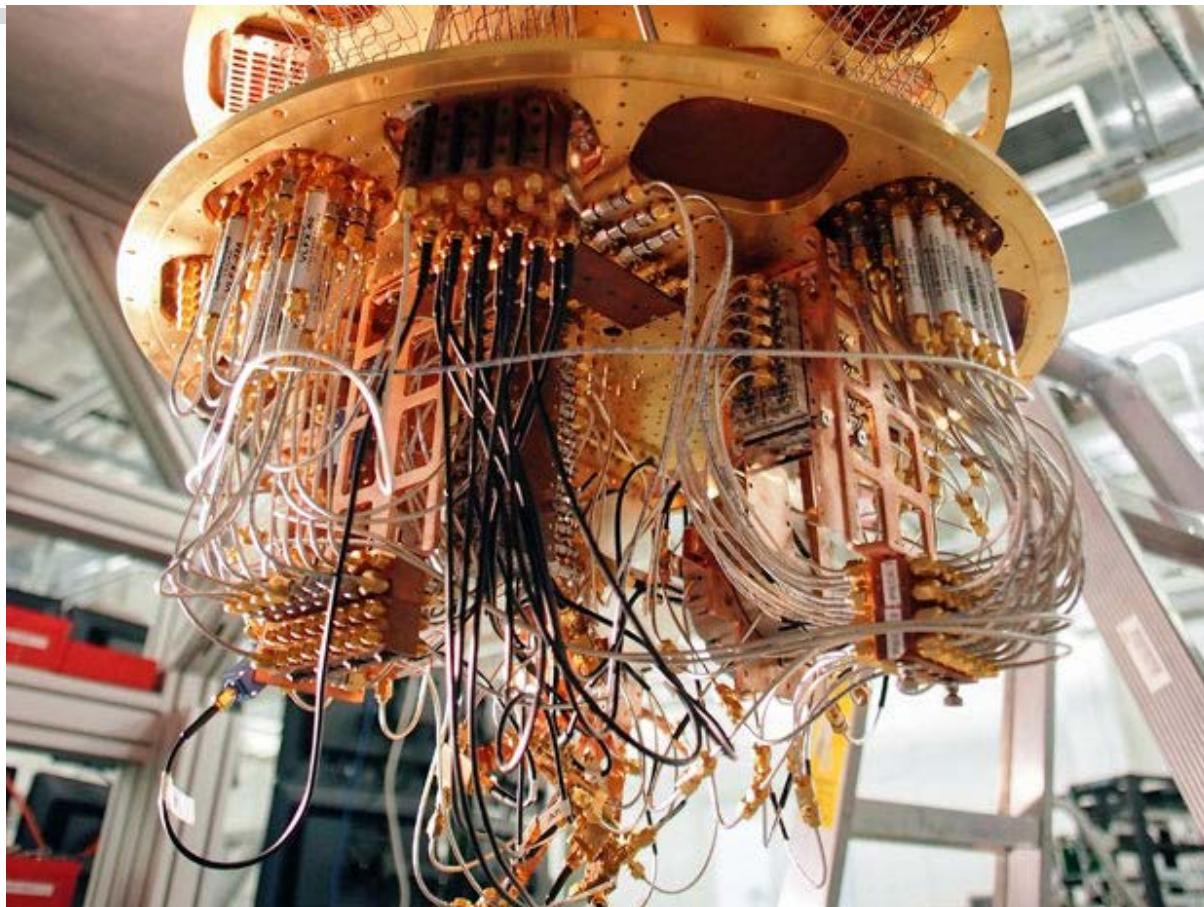
FIGURE 1 | D-Wave's recent generation Chimera connectivity graph G . Vertices indicate spin-valued variables represented by programmable qubits (h/s), and edges indicate programmable couplers (J_{ij} 's). G is a lattice of $K_{4,4}$ unit cells where missing qubits are the result of fabrication defects.

- Washington world-largest QC: 1152 qubits, 933 operational
- Quantum operation confirmed for 8-qubit register
- Consistent with either quantum or classical operation

State of the art for superconducting qubits

- Superconducting qubits are currently the most advanced technology for building scalable quantum circuits
- Gates and simple algorithms have been reported with up to 19 qubits (charge, flux, phase, and transmon qubits)
- Quantum gates take time 10-50 ns
- Coherence times T_1 , T_2 are currently in the range 10-50 μ s (2D) and up to several 100 μ s (3D)
- D-Wave is marketing superconducting quantum annealer (simulator)
- Google, IBM, Microsoft enter the race towards building a scalable quantum computer

КВАНТОВЫЙ ПРОЦЕССОР GOOGLE

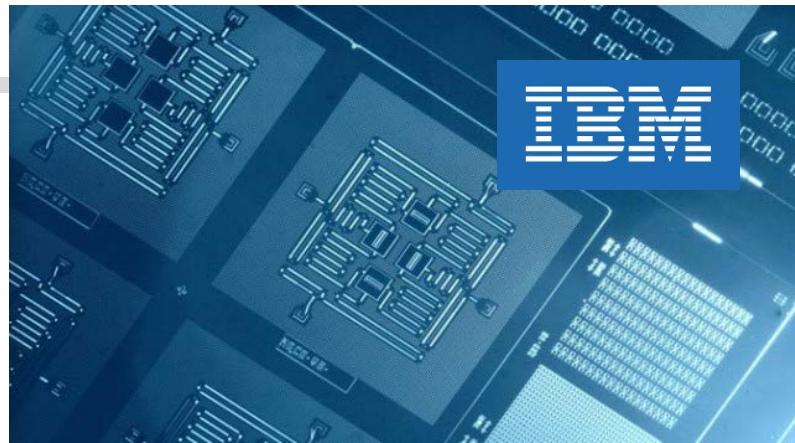


9 кубитов
Google
(2015)

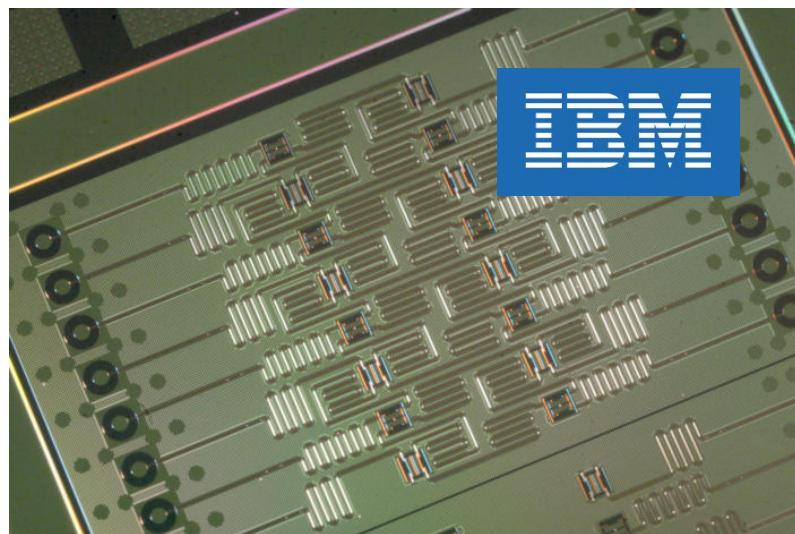
Google
John
Martinis

Google планировал к концу 2017 года продемонстрировать квантовый процессор на 49 кубитах. Цель - продемонстрировать „quantum supremacy”.

КВАНТОВЫЙ ПРОЦЕССОР IBM



5-кубитный процессор IBM (2015)



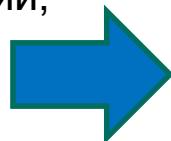
16-кубитный процессор IBM (2017)

СВЕРХПРОВОДЯЩИЕ КУБИТЫ В РОССИИ

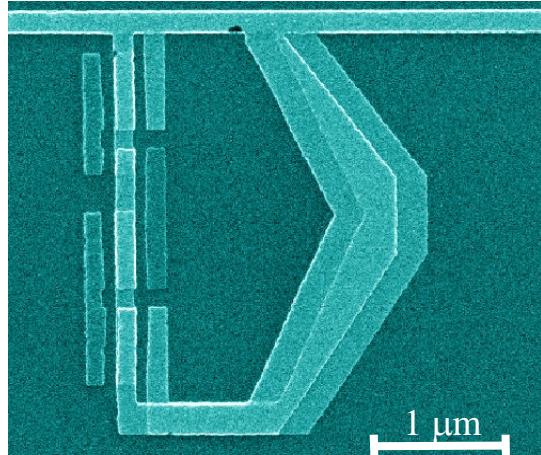
Начало: 2013-2015 гг.

июль 2013 г.

Первый кубит,
измеренный в России,
МИСиС и РКЦ

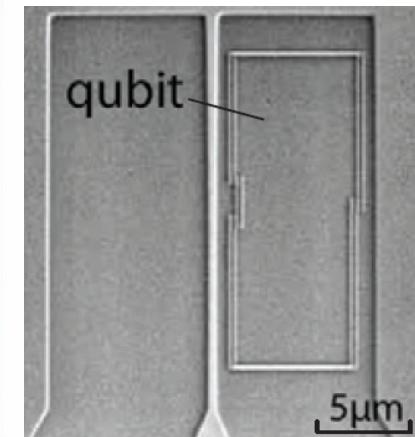
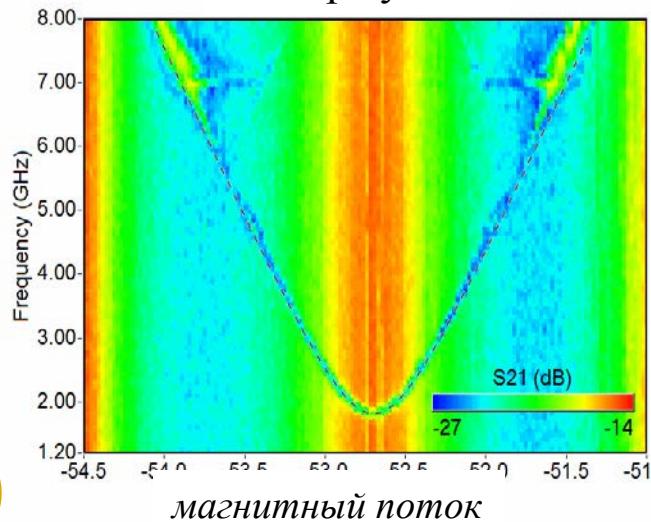


кубит изготовлен в
Германии



частота

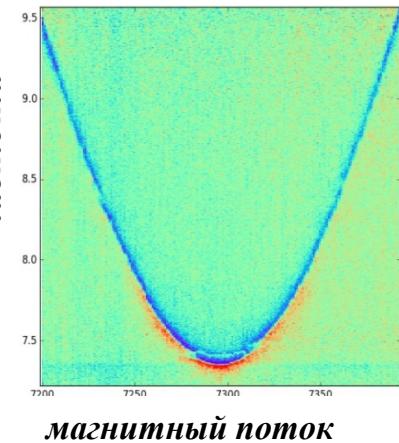
спектр кубита



май 2015 г.
первый кубит,
изготовленный
в России



частота



Практические задачи для квантовых компьютеров

Универсальные квантовые компьютеры

факторизация (алгоритм Шора)

поиск в базе данных (алгоритм Гровера и др.)

квантовая химия

расчет новых материалов

...

Адиабатические компьютеры, устройства квантового отжига

задачи оптимизации

машинаное обучение (combinatorial optimization problems)

искусственный интеллект

...

Аналоговые квантовые компьютеры (симуляторы)

моделирование квантовых систем

фотосинтез и лекарства

квантовый "Лего"

...